Memoirs on Differential Equations and Mathematical Physics $$\mathrm{Volume}\ 54,\ 2011,\ 83-98$$

Irena Rachůnková

ASYMPTOTIC PROPERTIES OF HOMOCLINIC SOLUTIONS OF SOME SINGULAR NONLINEAR DIFFERENTIAL EQUATION

Dedicated to the memory of Professor Temuri Chanturia

Abstract. We investigate an asymptotic behaviour of homoclinic solutions of the singular differential equation (p(t)u')' = p(t)f(u). Here f is Lipschitz continuous on \mathbb{R} and has at least two zeros 0 and L > 0. The function p is continuous on $[0, \infty)$, has a positive continuous derivative on $(0, \infty)$ and p(0) = 0.

2010 Mathematics Subject Classification. 34D05, 34A12, 34B40.

Key words and phrases. Singular ordinary differential equation of the second order, time singularities, asymptotic formula, homoclinic solutions.

რეზიუმე. გამოკვლეულია (p(t)u')' = p(t)f(u) სინგულირული დიფერენციალური განტოლების ჰომოკლინიკური ამონახსნების ასიმპტოტური თვისებები. აქ $f: \mathbb{R} \to \mathbb{R}$ ლიპშიცურად უწყვეტი, ხოლო $p: [0, \infty) \to [0, \infty)$ უწყვეტი და $(0, \infty)$ შუალედში წარმოებადი ფუნქციაა. ამახთან f(0) = f(L) = 0, სადაც L > 0, p(0) = 0 და p'(t) > 0, როცა t > 0.

1. INTRODUCTION

We investigate the differential equation

$$(p(t)u')' = p(t)f(u), \quad t \in (0,\infty),$$
 (1)

and throughout the paper it will be assumed that f satisfies

$$f \in Lip_{loc}(\mathbb{R}), \quad \exists L \in (0,\infty): f(L) = 0,$$
 (2)

$$\exists L_0 \in [-\infty, 0): \quad xf(x) < 0, \ x \in (L_0, 0) \cup (0, L),$$
^x
(3)

$$\exists \bar{B} \in (L_0, 0) : F(\bar{B}) = F(L), \text{ where } F(x) = -\int_0^{\infty} f(z) \, \mathrm{d}z, \ x \in \mathbb{R}, \quad (4)$$

and p fulfils

$$p \in C[0,\infty) \cap C^1(0,\infty), \quad p(0) = 0,$$
 (5)

$$p'(t) > 0, \ t \in (0,\infty), \quad \lim_{t \to \infty} \frac{p'(t)}{p(t)} = 0.$$
 (6)

Due to p(0) = 0, equation (1) has a singularity at t = 0.

Definition 1. A function $u \in C^1[0,\infty) \cap C^2(0,\infty)$ which satisfies equation (1) for all $t \in (0,\infty)$ is called a *solution* of equation (1).

Consider a solution u of equation (1). Since $u \in C^1[0,\infty)$, we have $u(0), u'(0) \in \mathbb{R}$, and the assumption p(0) = 0 yields p(0)u'(0) = 0. We can find that M > 0 and $\delta > 0$ such that $|f(u(t))| \leq M$ for $t \in (0, \delta)$. Integrating equation (1) and using the fact that p is increasing, we get

$$|u'(t)| = \left|\frac{1}{p(t)}\int_0^t p(s)f(u(s))\,\mathrm{d}s\right| \le \frac{M}{p(t)}\int_0^t p(s)\,\mathrm{d}s \le Mt, \quad t \in (0,\delta).$$

Consequently, the condition u'(0) = 0 is necessary for each solution u of equation (1). Therefore the set of all solutions of equation (1) forms a one-parameter system of functions u satisfying $u(0) = A, A \in \mathbb{R}$.

Definition 2. Let u be a solution of equation (1) and let L be of (2) and (3). Denote $u_{\sup} = \sup\{u(t) : t \in [0,\infty)\}$. If $u_{\sup} = L$ ($u_{\sup} < L$ or $u_{\sup} > L$), then u is called a *homoclinic* solution (a *damped* solution or an *escape* solution).

The existence and properties of these three types of solutions have been investigated in [19]–[23]. In particular, we have proved that if $u(0) \in (0, L)$, than u is a damped solution ([22], Theorem 2.3). Clearly, for u(0) = 0 and u(0) = L, equation (1) has a unique solution $u \equiv 0$ and $u \equiv L$, respectively.

In this paper we focus our attention on homoclinic solutions. According to the above considerations, such solutions have to satisfy the initial conditions

$$u(0) = B, \quad u'(0) = 0, \quad B < 0.$$
 (7)

Note that if we extend the function p(t) in equation (1) from the half-line onto \mathbb{R} (as an even function), then a homoclinic solution of (1) has the same limit L as $t \to -\infty$ and $t \to \infty$. This is a motivation for Definition 2.

We have proved in [21], Lemma 3.5, that a solution u of equation (1) is homoclinic if and only if u is strictly increasing and $\lim_{t\to\infty} u(t) = L$. If such homoclinic solution exists, then many important physical properties of corresponding models (see below) can be obtained. In particular, equation (1) is a generalization of the equation

$$u'' + \frac{k-1}{t}u' = f(u), \quad t \in (0,\infty),$$
(8)

and we can find in [16] that equation (8) with k > 1 and special forms of f arise in many areas, for example, in the study of phase transitions of Van der Waals fluids [3], [10], [24], in the population genetics, where it serves as a model for the spatial distribution of the genetic composition of a population [8], [9], in the homogeneous nucleation theory [1], in relativistic cosmology for description of particles which can be treated as domains in the universe [18], in the nonlinear field theory, in particular, when describing bubbles generated by scalar fields of the Higgs type in the Minkowski spaces [7]. Numerical simulations of solutions of (8), where f is a polynomial with three zeros, have been presented in [6], [14], [17]. Close problems on the existence of positive solutions are investigated in [2], [4], [5].

The main result of the present paper is contained in Section 3, Theorem 12, where we deduce an asymptotic formula for homoclinic solutions of equation (1). Note that many important results dealing with asymptotic properties of various types of differential equations can be found in the monograph by I. Kiguradze and T. Chanturia [12].

2. The Existence of Homoclinic Solutions

Here we cite theorems on the existence of homoclinic solutions. Remind that assumptions (2)–(6) are common for all these theorems. For a given B < 0, we denote the solution of problem (1), (7) by u_B .

Theorem 3. Assume that problem (1), (7) has an escape solution and let \overline{B} be of (4). Then there exists $B^* < \overline{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Proof. Theorem 2.3 in [22] shows that for any $B \in [\overline{B}, 0)$ there exists a unique solution u_B of problem (1), (7) and u_B is damped. Thus, if we denote by \mathcal{M}_d a set of all B < 0 such that u_B is a damped solution of problem (1), (7), then we obtain $\mathcal{M}_d \neq \emptyset$. Moreover, \mathcal{M}_d is open in $(-\infty, 0)$, due to Theorem 14 in [19]. Further, denote by \mathcal{M}_e a set of all B < 0 such that u_B is an escape solution of problem (1), (7). By our assumption, we have $\mathcal{M}_e \neq \emptyset$ and, by Theorem 20 in [19], the set \mathcal{M}_e is open in $(-\infty, 0)$, as well. Therefore, the set $\mathcal{M}_h = (-\infty, 0) \setminus (\mathcal{M}_d \cup \mathcal{M}_e)$ is non-empty. Let us choose $B^* \in \mathcal{M}_h$. Then $B^* < \overline{B}$, and by virtue of Definition 2, the supremum of the solution u_{B^*} on $(0, \infty)$ cannot be less than L and cannot be greater than L. Consequently, this supremum is equal to L, and u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Theorem 4. Assume that L_0 of (3) satisfies

$$L_0 \in (-\infty, 0), \quad f(L_0) = 0.$$
 (9)

Then there exists $B^* \in (L_0, \overline{B})$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Proof. Define

$$\tilde{f}(x) = \begin{cases} f(x) & \text{for} \quad x \le L, \\ 0 & \text{for} \quad x > L, \end{cases}$$

and consider the auxiliary equation

$$(p(t)u')' = p(t)\hat{f}(u), \quad t \in (0,\infty).$$
 (10)

By Theorem 10 and Lemma 9 in [20], there exists $B \in (L_0, \overline{B})$ such that u_B is an escape solution of problem (10), (7). If we modify the proof of Theorem 3 working on $(L_0, 0)$ instead of on $(-\infty, 0)$, we get a homoclinic solution u_{B^*} of problem (10), (7) having its starting value B^* in (L_0, \overline{B}) . Since u_{B^*} is increasing on $(0, \infty)$ (see e.g., Lemma 3.5 in [21]), we have

$$B^* \le u_{B^*}(t) < L, \quad t \in [0, \infty),$$
(11)

and u_{B^*} is likewise a solution of equation (1).

Theorem 4 assumes that f has the negative finite zero L_0 . The following two theorems concern the case that $L_0 = -\infty$ and f is positive on $(-\infty, 0)$. Then a behavior of f near $-\infty$ plays an important role. Equations with fhaving sublinear or linear behavior near $-\infty$ are discussed in the following theorem.

Theorem 5. Assume that f(x) > 0 for $x \in (-\infty, 0)$ and

$$0 \le \limsup_{x \to -\infty} \frac{f(x)}{|x|} < \infty.$$
(12)

Then there exists $B^* < \overline{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Proof. In the linear case, that is if we assume

$$0 < \limsup_{x \to -\infty} \frac{f(x)}{|x|} < \infty,$$

the assertion follows from Theorem 5.1 in [21]. Consider the sublinear case in which we work with the condition

$$\limsup_{x \to -\infty} \frac{f(x)}{|x|} = 0.$$

Assumption f > 0 on $(-\infty, 0)$ provides us with

$$\lim_{x \to -\infty} \frac{f(x)}{|x|} = 0$$

and Theorem 19 in [19] guarantees the existence of $B < \overline{B}$ such that u_B is an escape solution of problem (10), (7). Theorem 3 and estimate (11) yield $B^* < \overline{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Theorem 6. Assume that f(x) > 0 for $x \in (-\infty, 0)$ and there exists $k \ge 2$ such that

$$\lim_{t \to 0+} \frac{p'(t)}{t^{k-2}} \in (0,\infty).$$
(13)

Further, let $r \in (1, \frac{k+2}{k-2})$ be such that f fulfils

x

$$\lim_{n \to -\infty} \frac{f(x)}{|x|^r} \in (0, \infty).$$
(14)

Then there exists $B^* < \overline{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Proof. Theorem 2.10 in [23] guarantees the existence of $B < \bar{B}$ such that u_B is an escape solution of problem (10), (7). Theorem 3 and estimate (11) yield $B^* < \bar{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Theorem 6 discusses a superlinear behavior of f near $-\infty$. Note that if k = 2, we can take arbitrary $r \in (0, \infty)$. The last existence theorem imposes an additional assumption on p only.

Theorem 7. Assume that *p* satisfies

$$\int_{0}^{1} \frac{\mathrm{d}s}{p(s)} < \infty. \tag{15}$$

Then there exists $B^* < \overline{B}$ such that u_{B^*} is a homoclinic solution of problem (1), (7) with $B = B^*$.

Proof. Using Theorem 18 in [19] instead of Theorem 2.10 in [23], we argue just as in the proof of Theorem 6. \Box

In the next section, the use will be made of the generalized Matell's theorem which can be found as Theorem 6.5 in the monograph by I. Kiguradze [11]. For our purpose we provide its following special case.

Consider an interval $J \subset \mathbb{R}$. We write AC(J) for the set of functions, absolutely continuous on J, and $AC_{loc}(J)$ for the set of functions belonging to AC(I) for each compact interval $I \subset J$. Choose T > 0 and a function

matrix $A(t) = (a_{i,j}(t))_{i,j \leq 2}$ which is defined on (T, ∞) . Denote by $\lambda(t)$ and $\mu(t)$ the eigenvalues of $A(t), t \in (T, \infty)$. Further, suppose that

$$\lambda = \lim_{t \to \infty} \lambda(t) \quad \text{and} \quad \mu = \lim_{t \to \infty} \mu(t)$$

are different eigenvalues of the matrix $A = \lim_{t \to \infty} A(t)$ and let **l** and **m** be eigenvectors of A corresponding to λ and μ , respectively.

Theorem 8 ([11]). Assume that

$$a_{i,j} \in AC_{loc}(T,\infty), \quad \left| \int_{T}^{\infty} a'_{i,j}(t) \,\mathrm{d}t \right| < \infty, \quad i,j = 1, 2,$$
 (16)

and there exists $c_0 > 0$ such that

$$\int_{s}^{t} \operatorname{Re}(\lambda(\tau) - \mu(\tau)) \, \mathrm{d}\tau \le c_0, \quad T \le s < t,$$
(17)

or

$$\int_{T}^{\infty} \operatorname{Re}(\lambda(\tau) - \mu(\tau)) \, \mathrm{d}\tau = \infty,$$

$$\int_{s}^{t} \operatorname{Re}(\lambda(\tau) - \mu(\tau)) \, \mathrm{d}\tau \ge -c_0, \quad T \le s < t.$$
(18)

Then the differential system

$$\mathbf{x}'(t) = A(t)\mathbf{x}(t) \tag{19}$$

has a fundamental system of solutions $\mathbf{x}(t)$, $\mathbf{y}(t)$ such that

$$\lim_{t \to \infty} \mathbf{x}(t) \mathrm{e}^{-\int_T^t \lambda(\tau) \,\mathrm{d}\tau} = \mathbf{l}, \quad \lim_{t \to \infty} \mathbf{y}(t) \mathrm{e}^{-\int_T^t \mu(\tau) \,\mathrm{d}\tau} = \mathbf{m}.$$
 (20)

3. Asymptotic Behavior of Homoclinic Solutions

In this section we assume that $B < \overline{B}$ is such that the corresponding solution u of the initial problem (1), (7) is homoclinic. Hence u fulfils

$$u(0) = B, \quad u'(0) = 0, \quad u'(t) > 0, \ t \in (0, \infty), \quad \lim_{t \to \infty} u(t) = L.$$
 (21)

Moreover, due to (1),

$$u''(t) + \frac{p'(t)}{p(t)}u'(t) = f(u(t)), \quad t > 0,$$
(22)

and, by multiplication and integration over [0, t],

$$\frac{u'^2(t)}{2} + \int_0^t \frac{p'(s)}{p(s)} u'^2(s) \,\mathrm{d}s = F(u(0)) - F(u(t)), \quad t > 0.$$
(23)

Irena Rachůnková

Therefore

$$0 \le \lim_{t \to \infty} \int_0^t \frac{p'(s)}{p(s)} u'^2(s) \,\mathrm{d}s \le F(B) - F(L) < \infty,$$

and hence there exists

$$\lim_{t \to \infty} \int_0^t \frac{p'(s)}{p(s)} u'^2(s) \,\mathrm{d}s.$$

Consequently, according to (23), $\lim_{t\to\infty} u'^2(t)$ exists, as well. Since u is bounded on $[0,\infty)$, we get

$$\lim_{t \to \infty} u^{\prime 2}(t) = \lim_{t \to \infty} u^{\prime}(t) = 0.$$
(24)

In order to derive an asymptotic formula for u we have to characterize a behavior of p in ∞ and that of f near L more precisely. In particular, we put

$$h(x) := \frac{f(x)}{x - L}, \quad x < L$$

and work with the following assumptions:

$$\exists c, \eta > 0: \quad h \in C^1[L - \eta, L], \quad \lim_{x \to L^-} h(x) = h(L) = c, \tag{25}$$

$$p' \in AC_{loc}(0,\infty), \quad \exists n \ge 2: \quad \lim_{t \to \infty} \frac{p'(t)}{t^{n-2}} \in (0,\infty).$$
 (26)

For the sake of simplicity we transform L to the origin by the substitution

$$z(t) = L - u(t), \quad t \in [0, \infty),$$
 (27)

and put

$$g(y) = -f(L-y), \quad y > 0.$$
 (28)

Then the function z given by (27) is a solution of the equation

$$(p(t)z')' = p(t)g(z), \quad t \in (0,\infty),$$
(29)

satisfies

$$z(0) = L + |B|, \quad z'(0) = 0, \quad z'(t) < 0, \ t \in (0, \infty),$$
(30)

$$\lim_{t \to \infty} z(t) = 0, \quad \lim_{t \to \infty} z'(t) = 0.$$
(31)

Lemma 9. Assume the above condition (25) holds and let z be given by (27). Then there exists T > 0 such that

$$|z'(t)| > \sqrt{\frac{c}{2}} z(t), \quad t \ge T.$$
 (32)

Proof. According to (29), the function z fulfils the following equation:

$$z''(t) = -\frac{p'(t)}{p(t)}z'(t) + g(z(t)), \quad t \in (0,\infty).$$
(33)

Define the Lyapunov function \boldsymbol{V} by

$$V(t) = \frac{z'^{2}(t)}{2} + G(z(t)), \qquad (34)$$

where

$$G(x) = -\int_{0}^{x} g(s) \,\mathrm{d}s.$$

Owing to (3), (4) and $B < \overline{B}$, the function G fulfils

$$G(L+|B|) = -\int_{0}^{L+|B|} g(s) \,\mathrm{d}s = \int_{B}^{L} f(s) \,\mathrm{d}s = F(B) - F(L) > 0.$$

Thus V(0) = G(L + |B|) > 0. Further, using (33), we have

$$V'(t) = z'(t)z''(t) - g(z(t))z'(t) = -\frac{p'(t)}{p(t)}z'^{2}(t) < 0, \quad t > 0.$$

Hence V is decreasing on $(0, \infty)$ and, by (31), (34), we get $\lim_{t \to \infty} V(t) = 0$. Consequently, V(t) > 0 for $t \in [0, \infty)$ which implies that

$$\frac{z'^2(t)}{2} > -G(z(t)), \quad t > 0.$$
(35)

Let y = L - x. Then, using (25) and (28), we deduce

$$-\lim_{y \to 0+} \frac{G(y)}{y^2} = \lim_{y \to 0+} \frac{g(y)}{2y} = \frac{1}{2} \lim_{x \to L^-} \frac{f(x)}{x - L} = \frac{c}{2}.$$

Hence by virtue of (31), there exists T > 0 such that

$$-\frac{G(z(t))}{z^2(t)} > \frac{c}{4}, \quad t \ge T$$

This, together with (35), results in

$$\frac{z'^2(t)}{2} > \frac{c}{4}z^2(t), \quad t \ge T.$$

Consequently, we get (32).

Lemma 10. Assume that the condition (25) holds and let z and g be given by (27) and (28), respectively. Then

$$\int_{1}^{\infty} \left| \frac{g(z(\tau))}{z(\tau)} - c \right| \, \mathrm{d}\tau < \infty.$$
(36)

 $Irena\ Rachůnková$

Proof. Let us put

$$\tilde{h}(y) = \frac{g(y)}{y}, \quad y > 0.$$
(37)

By (25) and (28), we have

$$h(L-y) = \tilde{h}(y), \ y > 0, \quad \tilde{h} \in C^1[0,\eta], \quad \lim_{y \to 0^+} \tilde{h}(y) = \tilde{h}(0) = c,$$
 (38)

and there exists $M_0 \in (0, \infty)$ such that

$$\left|\frac{\mathrm{d}\tilde{h}(y)}{\mathrm{d}y}\right| \le M_0, \quad y \in [0,\eta].$$

The Mean Value Theorem guarantees the existence of $\theta \in (0, 1)$ such that

$$\tilde{h}(y) = c + y \frac{\mathrm{d}\tilde{h}(\theta y)}{\mathrm{d}y}, \quad y \in (0, \eta].$$

By (31), there exists $T \ge 1$ such that $0 < z(t) \le \eta$ for $t \ge T$ and hence, according to (37),

$$\left|\frac{g(z(t))}{z(t)} - c\right| \le M_0 z(t), \quad t \ge T.$$
(39)

Using (2), (28) and z > 0 on [1, T], we can find $M_1 \in (0, \infty)$ such that

$$\int_{1}^{T} \left| \frac{g(z(\tau))}{z(\tau)} - c \right| \, \mathrm{d}\tau \le M_1,$$

and, without loss of generality, we may assume that T is chosen in such a way that (32) is valid, as well. Therefore, using (32) and (39), we get

$$\int_{1}^{t} \left| \frac{g(z(\tau))}{z(\tau)} - c \right| d\tau \le M_{1} + M_{0} \int_{T}^{t} z(\tau) d\tau <$$

$$< M_{1} + \sqrt{\frac{2}{c}} M_{0} \int_{T}^{t} |z'(\tau)| d\tau = M_{1} - \sqrt{\frac{2}{c}} M_{0} \int_{T}^{t} z'(\tau) d\tau =$$

$$= M_{1} + \sqrt{2c} M_{0} (z(T) - z(t)), \quad t \ge T.$$

Letting $t \to \infty$ and using (31), we obtain (36).

Lemma 11. Assume that the condition (26) holds. Then

$$\int_{1}^{\infty} \left(\frac{p'(\tau)}{p(\tau)}\right)^2 \,\mathrm{d}\tau < \infty. \tag{40}$$

Proof. The condition (26) implies that there exists $c_0 \in (0, \infty)$ such that

$$\lim_{t \to \infty} \frac{p'(t)}{t^{n-2}} = c_0, \quad \lim_{t \to \infty} \frac{p(t)}{t^{n-1}} = \frac{c_0}{n-1}.$$

Therefore

$$\lim_{t \to \infty} t^2 \left(\frac{p'(t)}{p(t)}\right)^2 = (n-1)^2.$$

Hence we can find $T \ge 1$ such that

$$\left(\frac{p'(t)}{p(t)}\right)^2 < \frac{n^2}{t^2}, \quad t \ge T, \tag{41}$$

and due to (5) and (6), we can find $M_3 \in (0, \infty)$ such that

$$\int_{1}^{T} \left(\frac{p'(\tau)}{p(\tau)}\right)^2 \, \mathrm{d}\tau \le M_3.$$

Consequently,

$$\int_{1}^{t} \left(\frac{p'(\tau)}{p(\tau)}\right)^{2} d\tau < M_{3} + n^{2} \int_{T}^{t} \frac{d\tau}{\tau^{2}} = n^{2} \left(\frac{1}{T} - \frac{1}{t}\right), \quad t \ge T.$$

Letting $t \to \infty$, we get (40).

The main result on the asymptotic behavior of homoclinic solutions is contained in the following theorem.

Theorem 12. Assume that (25) and (26) hold. Let $B < \overline{B}$ be such that the corresponding solution u of the initial problem (1), (7) is homoclinic. Then u fulfils the equation

$$\lim_{t \to \infty} (L - u(t)) e^{\sqrt{c}t} \sqrt{p(t)} \in (0, \infty).$$
(42)

Remark 13. A similar asymptotic formula for positive solutions of equation (8), where k > 1 and $f(x) = x - |x|^r \operatorname{sign} x, r > 1$, has been derived in [13], Theorem 6.1.

Proof. Step 1. Construction of an auxiliary linear differential system. Consider the function z given by (27). According to (29), z satisfies

$$z'' + \frac{p'(t)}{p(t)}z' = \frac{g(z(t))}{z(t)}z(t), \quad t \in (0,\infty).$$
(43)

Having z at hand, we introduce the linear differential equation

$$v'' + \frac{p'(t)}{p(t)}v' = \frac{g(z(t))}{z(t)}v,$$
(44)

and the corresponding linear differential system

$$x'_1 = x_2, \quad x'_2 = \frac{g(z(t))}{z(t)}x_1 - \frac{p'(t)}{p(t)}x_2.$$
 (45)

 \square

Denote

$$A(t) = (a_{i,j}(t))_{i,j \le 2} = \begin{pmatrix} 0 & 1\\ \frac{g(z(t))}{z(t)} & -\frac{p'(t)}{p(t)} \end{pmatrix}, \quad A = \begin{pmatrix} 0 & 1\\ c & 0 \end{pmatrix}.$$

By (6), (31), (37) and (38),

$$A = \lim_{t \to \infty} A(t).$$

Eigenvalues of A are the numbers $\lambda = \sqrt{c}$ and $\mu = -\sqrt{c}$, eigenvectors of A are $\mathbf{l} = (1, \sqrt{c})$ and $\mathbf{m} = (1, -\sqrt{c})$, respectively. Denote

$$D(t) = \left(\frac{p'(t)}{2p(t)}\right)^2 + \frac{g(z(t))}{z(t)}, \quad t \in (0,\infty).$$
(46)

Then eigenvalues of A(t) have the form

$$\lambda(t) = -\frac{p'(t)}{2p(t)} + \sqrt{D(t)}, \quad \mu(t) = -\frac{p'(t)}{2p(t)} - \sqrt{D(t)}, \quad t \in (0, \infty).$$
(47)

We can see that

$$\lim_{t\to\infty}\lambda(t)=\lambda,\quad \lim_{t\to\infty}\mu(t)=\mu.$$

Step 2. Verification of the Assumptions of Theorem 8. Due to (31) and (38), we can find $T \ge 1$ such that

$$0 < z(t) \le \eta, \quad D(t) > 0, \quad t \in (T, \infty).$$
 (48)

Therefore, by (37) and (38),

$$a_{21}(t) = \frac{g(z(t))}{z(t)} \in AC_{loc}(T, \infty),$$

and hence

$$\left| \int_{T}^{\infty} \left(\frac{g(z(t))}{z(t)} \right)' \, \mathrm{d}t \right| = \left| \lim_{t \to \infty} \frac{g(z(t))}{z(t)} - \frac{g(z(T))}{z(T)} \right| = \left| c - \frac{g(z(T))}{z(T)} \right| < \infty.$$

Further, by (26), $a_{22}(t) = -\frac{p'(t)}{p(t)} \in AC_{loc}(T, \infty)$. Hence due to (6),

$$\left| \int_{T}^{\infty} \left(\frac{p'(t)}{p(t)} \right)' \, \mathrm{d}t \right| = \left| \lim_{t \to \infty} \frac{p'(t)}{p(t)} - \frac{p'(T)}{p(T)} \right| = \frac{p'(T)}{p(T)} < \infty$$

Since $a_{11}(t) \equiv 0$ and $a_{12}(t) \equiv 1$, it is not difficult to see that (16) is satisfied. Using (47), we get $\operatorname{Re}(\lambda(t) - \mu(t)) = 2\sqrt{D(t)} > 0$ for $t \in (T, \infty)$. Since $\lim_{t \to \infty} \sqrt{D(t)} = \sqrt{c} > 0$, we have

$$\int_{T}^{\infty} \operatorname{Re}(\lambda(\tau) - \mu(\tau)) \, \mathrm{d}\tau = \infty, \quad \int_{s}^{t} \operatorname{Re}(\lambda(\tau) - \mu(\tau)) \, \mathrm{d}\tau > 0, \quad T \le s < t.$$

Consequently, (18) is valid.

Step 3. Application of Theorem 8. By Theorem 8, there exists a fundamental system $\mathbf{x}(t) = (x_1(t), x_2(t)), \mathbf{y}(t) = (y_1(t), y_2(t))$ of solutions of (45) such that (20) is valid. Hence

$$\lim_{t \to \infty} x_1(t) e^{-\int_T^t \lambda(\tau) \, \mathrm{d}\tau} = 1, \quad \lim_{t \to \infty} y_1(t) e^{-\int_T^t \mu(\tau) \, \mathrm{d}\tau} = 1.$$
(49)

Using (47), for $t \ge T$ we get

$$\begin{split} \exp\left(-\int_{T}^{t}\lambda(\tau)\,\mathrm{d}\tau\right) &= \exp\left(\int_{T}^{t}\left(\frac{p'(\tau)}{2p(\tau)} - \sqrt{D(\tau)}\right)\mathrm{d}\tau\right) = \\ &= \exp\left(\frac{1}{2}\ln\frac{p(t)}{p(T)}\right)\exp\left(-\int_{T}^{t}\sqrt{D(\tau)}\,\mathrm{d}\tau\right) = \\ &= \sqrt{\frac{p(t)}{p(T)}}\exp\left(-\int_{T}^{t}\sqrt{D(\tau)}\,\mathrm{d}\tau\right), \end{split}$$

and

$$\exp\left(-\int_{T}^{t}\mu(\tau)\,\mathrm{d}\tau\right) = \exp\left(\int_{T}^{t}\left(\frac{p'(\tau)}{2p(\tau)} + \sqrt{D(\tau)}\right)\,\mathrm{d}\tau\right) =$$
$$= \exp\left(\frac{1}{2}\ln\frac{p(t)}{p(T)}\right)\exp\left(\int_{T}^{t}\sqrt{D(\tau)}\,\mathrm{d}\tau\right) =$$
$$= \sqrt{\frac{p(t)}{p(T)}}\exp\left(\int_{T}^{t}\sqrt{D(\tau)}\,\mathrm{d}\tau\right).$$

Further,

$$\int_{T}^{t} \sqrt{D(\tau)} \,\mathrm{d}\tau = E_0(t) + \sqrt{c}(t-T),$$

where

$$E_0(t) = \int_T^t \frac{D(\tau) - c}{\sqrt{D(\tau)} + \sqrt{c}} \,\mathrm{d}\tau, \quad t \ge T.$$
(50)

Hence

$$\exp\left(-\int_{T}^{t}\lambda(\tau)\,\mathrm{d}\tau\right) = \sqrt{\frac{p(t)}{p(T)}}\mathrm{e}^{-E_{0}(t)}\mathrm{e}^{-\sqrt{c}(t-T)}, \quad t \ge T, \qquad (51)$$

$$\exp\left(-\int_{T}^{t}\mu(\tau)\,\mathrm{d}\tau\right) = \sqrt{\frac{p(t)}{p(T)}}\mathrm{e}^{E_{0}(t)}\mathrm{e}^{\sqrt{c}(t-T)}, \quad t \ge T.$$
(52)

Using (36), (40) and (46), we can find $K_0 \in (0, \infty)$ such that for $t \ge T$,

$$\int_{T}^{t} \left| \frac{D(\tau) - c}{\sqrt{D(\tau)} + \sqrt{c}} \right| d\tau \leq \\ \leq \frac{1}{\sqrt{c}} \left(\int_{T}^{t} \left(\frac{p'(\tau)}{2p(\tau)} \right)^{2} d\tau + \int_{T}^{t} \left| \frac{g(z(\tau))}{z(\tau)} - c \right| d\tau \right) \leq K_{0}.$$

Consequently, due to (50),

$$\lim_{t \to \infty} E_0(t) = E_0 \in \mathbb{R}$$

Therefore (49), (51) and (52) imply

$$1 = \lim_{t \to \infty} x_1(t) \sqrt{\frac{p(t)}{p(T)}} e^{-E_0} e^{-\sqrt{c}(t-T)},$$

$$1 = \lim_{t \to \infty} y_1(t) \sqrt{\frac{p(t)}{p(T)}} e^{E_0} e^{\sqrt{c}(t-T)}.$$

Since by (26),

$$\lim_{t \to \infty} \sqrt{p(t)} e^{-\sqrt{c}t} = \lim_{t \to \infty} \sqrt{\frac{p(t)}{t^{n-1}}} t^{(n-1)/2} e^{-\sqrt{c}t} = 0,$$
$$\lim_{t \to \infty} \sqrt{p(t)} e^{\sqrt{c}t} = \infty,$$

we obtain

$$\lim_{t \to \infty} x_1(t) = \infty, \quad \lim_{t \to \infty} y_1(t) = 0.$$
(53)

Step 4. Asymptotic Formula. According to (43), z is likewise a solution of (44). Therefore there are $c_1, c_2 \in \mathbb{R}$ such that $z(t) = c_1 x_1(t) + c_2 y_1(t)$, $t \in (0, \infty)$. Having in mind (30), (31), (49) and (53), we get $c_1 = 0$, $c_2 y_1(t) > 0$ on $(0, \infty)$, and $c_2 \in (0, \infty)$. Consequently, $z(t) = c_2 y_1(t)$ and

$$1 = \lim_{t \to \infty} \frac{1}{c_2} z(t) \sqrt{\frac{p(t)}{p(T)}} e^{E_0} e^{\sqrt{c}(t-T)},$$

which together with (27) yields (42).

Acknowledgements

This work was supported by the Council of Czech Government MSM 6198959214.

References

- 1. F. F. ABRAHAM, Homogeneous nucleation theory. Acad. Press, New York, 1974.
- H. BERESTYCKI, P.-L. LIONS, AND L. A. PELETIER, An ODE approach to the existence of positive solutions for semilinear problems in R^N. Indiana Univ. Math. J. 30 (1981), No. 1, 141–157.
- V. BONGIORNO, L. E. SCRIVEN, AND H. T. DAVIS, Molecular theory of fluid interfaces, J. Colloid and Interface Science 57 (1967), 462–475.
- D. BONHEURE, J. M. GOMES, AND L. SANCHEZ, Positive solutions of a second-order singular ordinary differential equation. *Nonlinear Anal.* 61 (2005), No. 8, 1383–1399.
- M. CONTI, L. MERIZZI, AND S. TERRACINI, Radial solutions of superlinear equations on R^N. I. A global variational approach. Arch. Ration. Mech. Anal. 153 (2000), No. 4, 291–316.
- F. DELL'ISOLA, H. GOUIN, AND G. ROTOLI, Nucleation of spherical shell-like interfaces by second gradient theory: Numerical simulations. *Eur. J. Mech.*, B 15 (1996), No. 4, 545–568.
- G. H. DERRICK, Comments on nonlinear wave equations as models for elementary particles. J. Mathematical Phys. 5 (1964), 1252–1254.
- P. C. FIFE, Mathematical aspects of reacting and diffusing systems. Lecture Notes in Biomathematics, 28. Springer-Verlag, Berlin-New York, 1979.
- R. A. FISHER, The wave of advance of advantegeous genes. J. Eugenics 7 (1937), No. 4, 355–369.
- H. GOUIN AND G. ROTOLI, An analytical approximation of density profile and surface tension of microscopic bubbles for van der Waals fluids. *Mech. Res. Commun.* 24 (1997), No. 3, 255–260.
- 11. I. KIGURADZE, Some singular boundary value problems for ordinary differential equations. (Russian) *Izdat. Tbilis. Univ., Tbilisi*, 1975.
- I. KIGURADZE AND T. CHANTURIA, Asymptotic properties of solutions of nonautonomous ordinary differential equations. Translated from the 1985 Russian original. Mathematics and its Applications (Soviet Series), 89. Kluwer Academic Publishers Group, Dordrecht, 1993.
- I. KIGURADZE AND B. SHEKHTER, Singular boundary value problems for secondorder ordinary differential equations. (Russian) Itogi Nauki i Tekhniki, Current problems in mathematics. Newest results, Vol. 30 (Russian), 105–201, 204, Akad. Nauk SSSR, Vsesoyuz. Inst. Nauchn. i Tekhn. Inform., Moscow, 1987; English transl.: J. Soviet Math. 43 (1988), No. 2, 2340–2417.
- G. KITZHOFER, O. KOCH, P. LIMA, AND E. WEINMÜLLER, Efficient numerical solution of the density profile equation in hydrodynamics. J. Sci. Comput. 32 (2007), No. 3, 411–424.
- O. KOCH, P. KOFLER, AND E. WEINMÜLLER, Initial value problems for systems of ordinary first and second order differential equations with a singularity of the first kind. Analysis (Munich) 21 (2001), No. 4, 373–389.
- N. B. KONYUKHOVA, P. M. LIMA, M. L. MORGADO, AND M. B. SOLOLVIEV, Bubbles and droplets in nonlinear physics models: analysis and numerical simulation of singular nonlinear boundary value problem. (Russian) *Zh. Vychisl. Mat. Mat. Fiz.* 48 (2008), No. 11, 2019–2023; English transl.: *Comput. Math. Math. Phys.* 48 (2008), No. 11, 2018–2058.
- P. M. LIMA, N. V. CHEMETOV, N. B. KONYUKHOVA, AND A. I. SUKOV, Analyticalnumerical investigation of bubble-type solutions of nonlinear singular problems. J. Comput. Appl. Math. 189 (2006), No. 1-2, 260–273.
- A. P. LINDE, Particle Physics and Inflationary Cosmology. Harwood Academic, Chur, Switzerland, 1990.
- I. RACHŮNKOVÁ AND J. TOMEČEK, Bubble-type solutions of nonlinear singular problems. Math. Comput. Modelling 51 (2010), No. 5-6, 658–669.

- I. RACHŮNKOVÁ AND J. TOMEČEK, Strictly increasing solutions of a nonlinear singular differential equation arising in hydrodynamics. *Nonlinear Anal.* 72 (2010), No. 3-4, 2114–2118.
- I. RACHŮNKOVÁ, J. TOMEČEK, Homoclinic solutions of singular nonautonomous second-order differential equations. *Bound. Value Probl.* 2009, Art. ID 959636, 21 pp.
- I. RACHŮNKOVÁ, L. RACHŮNEK, AND J. TOMEČEK, Existence of oscillatory solutions of singular nonlinear differential equations. *Abstr. Appl. Anal.* 2011, Art. ID 408525, 20 pp.
- I. RACHŮNKOVÁ AND J. TOMEČEK, Superlinear singular problems on the half line. Bound. Value Probl. 2010, Art. ID 429813, 18 pp.
- 24. J. D. VAN DER WAALS AND R. KOHNSTAMM, Lehrbuch der Thermodynamik. vol. 1, Barth Verlag, Leipzig, 1908.

(Received 13.09.2011)

Authors' address:

Department of Mathematics Faculty of Science, Palacký University 17 listopadu 12, 771 46 Olomouc Czech Republic E-mail: irena.rachunkova@upol.cz