$\begin{array}{c} \text{INTEGRABILITY OF} \\ \text{CLASSICAL AFFINE W-ALGEBRAS} \end{array}$

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To the memory of Ernest Borisovich Vinberg

Abstract. We prove that all classical affine W-algebras $\mathcal{W}(\mathfrak{g}, f)$, where \mathfrak{g} is a simple Lie algebra and f is its non-zero nilpotent element, admit an integrable hierarchy of bi-Hamiltonian PDEs, except possibly for one nilpotent conjugacy class in G_2 , one in F_4 , and five in E_8 .

Contents

1. Introduction

2. Cyclic and quasi-cyclic elements. Integrable triples

- 2.1. Setup
- 2.2. Classification of cyclic elements
- 2.3. Quasi-cyclic elements
- 2.4. Integrable triples

3. Poisson vertex algebras, Hamiltonian equations and integrability

- 3.1. PVA and Hamiltonian PDE
- 3.2. Bi-Poisson vertex algebras and Lenard-Magri scheme of integrability

4. Classical affine W-algebras and generalized Drinfeld–Sokolov hierarchies

- 4.1. Definition of classical affine W-algebras
- 4.2. Integrable triples and generalized Drinfeld–Sokolov hierarchies

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1. Introduction

In order to define a Hamiltonian ODE one needs a Poisson algebra, that is a commutative associative algebra \mathcal{P} , endowed with a Poisson bracket $\mathcal{P} \otimes \mathcal{P} \to \mathcal{P}$, $a \otimes b \mapsto \{a,b\}$, and an element $h \in \mathcal{P}$, called a Hamiltonian function. Then the ODE

$$\frac{du}{dt} = \{h, u\}, \quad u = u(t) \in \mathcal{P},$$

is called Hamiltonian. Recall that, by definition of a Poisson algebra, the bracket $\{\cdot,\cdot\}$ should satisfy the Lie algebra axioms, and, for every $f \in \mathcal{P}$, the map $u \mapsto \{f,u\}$ should be a derivation of the associative product on \mathcal{P} (Leibniz rule).

In a similar fashion, in order to define a Hamiltonian PDE one needs a *Poisson* vertex algebra (abbreviated PVA), which is a differential algebra, that is a commutative associative algebra \mathcal{V} with a derivation ∂ , endowed with a PVA λ -bracket $\mathcal{V} \otimes \mathcal{V} \to \mathcal{V}$, $a \otimes b \mapsto \{a_{\lambda}b\}$, and an element $\int h \in \mathcal{V}/\partial \mathcal{V}$, called a *local Hamiltonian* functional. (One denotes by \int the canonical map $\mathcal{V} \to \mathcal{V}/\partial \mathcal{V}$ since it is the universal map satisfying integration by parts: $\int (\partial f)g = -\int f\partial g$.) A key property of a PVA \mathcal{V} is that the formula

$$\{\int f, \int g\} = \int \{f_{\lambda}g\}|_{\lambda=0} \tag{1.1}$$

produces a well-defined Lie algebra bracket on $V/\partial V$. The PDE

$$\frac{du}{dt} = \{ \int h, u \} := \{ h_{\lambda} u \} \big|_{\lambda = 0}, \quad u \in \mathcal{V}, \tag{1.2}$$

is called Hamiltonian. In (1.2) u = u(x,t) should be viewed as a function in the space variable x and the time t, $\partial = \partial/\partial x$ is the partial derivative with respect to x, while d/dt defines the time flow of the system. Recall that, by definition, the PVA λ -bracket should satisfy the Lie conformal algebra axioms, see axioms (i)–(iii) from Section 3.1, similar to the Lie algebra axioms, and the Leibniz rules (iv) and (iv') from Section 3.1. In particular, due to the first sesquilinearity axiom (i), the RHS of equation (1.2) is well defined (i.e., it does not depend on the choice of the representative of the coset $\int h$).

Recall that if \mathcal{V} carries two compatible PVA λ -brackets (i.e., any their linear combination is a PVA λ -bracket), such that an evolution PDE can be written in the form (1.2) for both of them, this evolution equation is called bi-Hamiltonian.

The notion of a PVA appears naturally in the theory of vertex algebras as their quasiclassical limit, cf. [DSK06], in the same way as a Poisson algebra appears naturally as a quasiclassical limit of a family of associative algebras.

However, the theory of Hamiltonian PDEs was started 15 years before the advent of the vertex algebra theory, in the work of Faddeev and Zakharov [ZF71], who attribute the construction to C. S. Gardner, which, in algebraic terms, is as follows. Let $\mathcal{V} = \mathbb{F}[u_i^{(n)} \mid i \in I = \{1, \dots, \ell\}, n \in \mathbb{Z}_{\geq 0}]$ be the algebra of differential polynomials in ℓ differential variables, with the derivation ∂ defined by $\partial u_i^{(n)} = u_i^{(n+1)}$. Given $\int h \in \mathcal{V}/\partial \mathcal{V}$, one defines a Hamiltonian PDE

$$\frac{du}{dt} = H(\partial) \frac{\delta}{\delta u} \int h, \tag{1.3}$$

where $u = (u_i)_{i \in I}$ is an ℓ -column vector of dependent variables, $H(\partial)$ is an $\ell \times \ell$ -matrix differential operator with coefficients in \mathcal{V} , called the *Poisson structure* of \mathcal{V} , and $\frac{\delta}{\delta u} \int h = \left(\frac{\delta}{\delta u_i} \int h\right)_{i \in I}$ is the ℓ -column vector of variational derivatives

$$\frac{\delta}{\delta u_i} \int h = \sum_{n \in \mathbb{Z}_{>0}} (-\partial)^n \frac{\partial h}{\partial u_i^{(n)}}, \quad i \in I.$$

(A Hamiltonian ODE can be written in local coordinates in a similar form $\frac{du}{dt} = H \frac{\partial}{\partial u} h$, where H is the transpose of the matrix of Poisson brackets of coordinate functions u_i , and $\frac{\partial}{\partial u}$ is the gradient.)

The basic assumption on the Poisson structure $H(\partial)$ is that formula

$$\{\int f, \int g\} = \int \frac{\delta \int g}{\delta u} H(\partial) \frac{\delta \int f}{\delta u}$$
 (1.4)

defines a Lie algebra structure on $\mathcal{V}/\partial\mathcal{V}$. A simple observation is that (1.3) coincides with (1.2), and (1.4) coincides with (1.1), if we let $H(\partial) = (\{u_{j\partial}u_i\}_{\rightarrow})_{i,j\in I}$, where the arrow means that ∂ should be moved to the right, see [BDSK09]. Thus PVAs provide a coordinate free approach to the theory of Hamiltonian PDEs.

The simplest example of a Hamiltonian PDE is the celebrated KdV equation

$$\frac{du}{dt} = 3uu' + cu''', \quad c \in \mathbb{F},\tag{1.5}$$

the only one studied in [ZF71]. It is Hamiltonian for the algebra of differential polynomials in one variable $\mathcal{V} = \mathbb{F}[u, u', u'', \dots]$, the local Hamiltonian functional $\int h = \frac{1}{2} \int u^2$, and the λ -bracket, defined on the generator u by

$$\{u_{\lambda}u\} = (\partial + 2\lambda)u + c\lambda^3,\tag{1.6}$$

and uniquely extended to $\mathcal V$ by the PVA axioms, which corresponds to the Poisson structure

$$H(\partial) = u' + 2u\partial + c\partial^3. \tag{1.7}$$

In fact, equation (1.5) is Hamiltonian also for another choice of the local Hamiltonian functional and the Poisson structure: $\int h_1 = \frac{1}{2} \int (u^3 + cuu'')$, $H_1(\partial) = \partial$, which makes the KdV equation a bi-Hamiltonian PDE.

For an arbitrary evolution PDE $\frac{du}{dt} = P$, $P \in \mathcal{V}$, an *integral of motion* is a local functional $\int f \in \mathcal{V}/\partial \mathcal{V}$ which is conserved by virtue of this equation. In particular, for the Hamiltonian equation (1.2) this is equivalent to the property that $\int h$ and $\int f$ are in *involution* (see (1.1)):

$$\{\int h, \int f\} = 0.$$

A bi-Hamiltonian PDE is called *integrable* if $\int h$ is contained in an infinite-dimensional abelian subalgebra of $V/\partial V$ with the Lie algebra bracket (1.1) for both PVA

 λ -brackets on \mathcal{V} . Given such a subalgebra of $\mathcal{V}/\partial\mathcal{V}$, choosing a basis $\{\int h_j\}_{j\in J}$, we obtain an integrable hierarchy of compatible bi-Hamiltonian PDEs:

$$\frac{du}{dt_j} = \{h_{j\lambda}u\}|_{\lambda=0}, \quad j \in J.$$
(1.8)

It is easy to see from the axioms of a PVA that "these flows commute", meaning that $\frac{d}{dt_i}\frac{d}{dt_j}=\frac{d}{dt_j}\frac{d}{dt_i}$.

We call a PVA \mathcal{V} with two compatible λ -brackets *integrable* if $\mathcal{V}/\partial\mathcal{V}$ with the two Lie algebra brackets (1.1) contains an infinite-dimensional abelian subalgebra with respect to both of them.

It was the paper [GGKM67] of Gardner, Green, Kruskal and Miura, where they proved that the KdV equation (1.5) admits infinitely many linearly independent integrals of motion, which initiated the whole theory of integrable PDEs. It follows from [ZF71] that these integrals are in involution with respect to the bracket (1.4), where $H(\partial)$ is given by (1.7). This leads to the whole KdV hierarchy as in (1.8).

In the paper [GD76], Gelfand and Dickey, using the idea of Lax operator [Lax68], constructed, for each integer $N \geq 2$, an integrable hierarchy of PDEs, called the Nth KdV hierarchy, and in [GD78] they showed that these hierarchies are Hamiltonian. (The corresponding bracket was conjectured earlier by Adler [Adl79].) Their 2nd KdV hierarchy coincides with the classical KdV hierarchy.

In the paper [Mag78] Magri proposed a simple algorithm, called nowadays the Lenard–Magri scheme, which allows one to prove that integrals of motion of a Hamiltonian PDE (1.3) are in involution, provided that the same equation can be written using a different Poisson structure $H_1(\partial)$ in place of $H(\partial)$, and a different local funtional $\int h_1$ in place of $\int h$. In such a case one obtains a bi-Hamiltonian hierarchy of PDEs. See [BDSK09] for details.

In their seminal paper [DS85] Drinfeld and Sokolov constructed the Lie algebra of local functionals obtained in [GD78] for the Nth KdV hierarchy via classical Hamiltonian reduction for the affine Kac–Moody algebra $\widehat{\mathfrak{sl}}_N$. This Drinfeld–Sokolov reduction was developed by them for an affine Kac–Moody algebra $\widehat{\mathfrak{g}}$, attached to an arbitrary simple Lie algebra \mathfrak{g} , which led to the construction of an integrable bi-Hamiltonian hierarchy of PDEs, called the Drinfeld–Sokolov hierarchy, attached to an arbitrary affine Kac–Moody algebra.

The Drinfeld–Sokolov reduction is based on a principal nilpotent element f of the simple Lie algebra \mathfrak{g} . It was extended in [FORTW92] to a certain class of nilpotent elements f of a simple Lie algebra \mathfrak{g} . The construction of a Drinfeld–Sokolov hierarchy in full generality turned out much more difficult, however.

The construction of the original Drinfeld–Sokolov hierarchy in [DS85] is based on Kostant's theorem that cyclic elements $f + E \in \mathfrak{g}$ (see Definition 2.2(a)), attached to a principal nilpotent element f, are semisimple. In a series of papers in the early 90's this construction was extended to other nilpotent elements admitting a semisimple cyclic element (see [DSKV13] for references).

The theory of integrable Hamiltonian hierarchies of PDEs has been naturally related to the theory of PVAs in [BDSK09]. An important class of PVAs, called the *classical affine W-algebras* and denoted by $W(\mathfrak{g}, f)$, where \mathfrak{g} is a simple Lie

algebra and f is its nilpotent element, was considered in [DSK06] as the quasiclassical limit of quantum affine W-algebras. It was then shown in [DSKV13] that the PVA $\mathcal{W}(\mathfrak{g}, f)$ can be obtained by a classical Hamiltonian reduction in the framework of PVA theory, analogous to the Drinfeld–Sokolov reduction. In the same paper it was proved that the Lie algebras of local functionals, constructed in [DS85] and [FORTW92], coincide with the Lie algebras $\mathcal{W}(\mathfrak{g}, f)/\partial \mathcal{W}(\mathfrak{g}, f)$ with the bracket (1.1).

Furthermore, the Drinfeld–Sokolov hierarchies and their generalizations have been constructed in [DSKV13] for the W-algebras $\mathcal{W}(\mathfrak{g}, f)$, using PVA techniques, provided that $f \in \mathfrak{g}$ is a nilpotent element of semisimple type. The latter means that there exists a semisimple cyclic element f + E attached to f. This establishes integrability of the PVA $\mathcal{W}(\mathfrak{g}, f)$ for f of semisimple type.

Unfortunately, the classification of cyclic elements in simple Lie algebras, obtained in [EKV13], shows that there are very few semisimple type nilpotent elements in classical Lie algebras, and only about half of the nilpotent elements in exceptional simple Lie algebras are of semisimple type.

The first basic idea of the present paper is that the Drinfeld–Sokolov method for constructing integrals of motion in [DS85], extended to the case of nilpotents f of semisimple type in [DSKV13], generalizes, after a simple modification, to all nilpotents f admitting a non-nilpotent cyclic element f + E (see Theorems 2.11 and 4.6).

The only nilpotent elements f which are left out from the above generalization are those of *nilpotent type*, i.e., such that all the cyclic elements f+E are nilpotent. According to [EKV13], there are altogether 15 conjugacy classes of nilpotent elements of nilpotent type in all exceptional Lie algebras, and, among classical Lie algebras, they exist only in \mathfrak{so}_n with $n \geq 7$, and correspond to partitions $(p_1 > p_2 = p_1 - 1 \geq \ldots)$, where p_1 is odd.

In order to treat the nilpotent elements f of nilpotent type, we use the idea of [DSKV13], another version of [FGMS95], that the Drinfeld–Sokolov method works also for those f which admit a semisimple quasi-cyclic element f+E (see Definition 2.2 (b)). In the present paper we show that this is also the case when f admits a non-nilpotent quasi-cyclic element (Theorems 2.14 and 4.6).

This establishes integrability of classical affine W-algebras $\mathcal{W}(\mathfrak{g}, f)$ for all classical simple Lie algebras \mathfrak{g} and all their nilpotent elements f, and for all the exceptional simple Lie algebras \mathfrak{g} and their nilpotent elements, except, possibly, the following types (in the notation of [CMG93]): $4A_1$, $2A_2 + 2A_1$, $2A_3$, $A_4 + A_3$ and A_7 in E_8 ; $\widetilde{A}_2 + A_1$ in F_4 ; \widetilde{A}_1 in G_2 .

The contents of the paper are as follows.

In Section 2 we define the notion of cyclic and quasi-cyclic elements (see Definitions 2.2(a) and (b)) and discuss their properties. We show that non-nilpotent cyclic and quasi-cyclic elements give rise to integrable triples in \mathfrak{g} , see Definition 2.10. In Section 3 we recall the notions of a PVA and integrable Hamiltonian PDE, and discuss the Lenard–Magri scheme of integrability. In Section 4 we recall the construction of classical affine W-algebras and show that any integrable triple gives rise to an integrable generalized Drinfeld–Sokolov hierarchy (Theorem 4.6). This theorem implies integrability of all classical affine W-algebras, associated to

classical Lie algebras, and all classical W-algebras $\mathcal{W}(\mathfrak{g}, f)$, associated to exceptional Lie algebras \mathfrak{g} , except for the seven nilpotents f mentioned above.

Throughout the paper the base field \mathbb{F} is an algebraically closed field of characteristic zero.

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2. Cyclic and quasi-cyclic elements. Integrable triples

2.1. Setup

Let \mathfrak{g} be a reductive finite-dimensional Lie algebra and let $f \in \mathfrak{g}$ be a non-zero nilpotent element. Recall that by the Jacobson–Morozov Theorem [CMG93, Thm. 3.3.1], any non-zero nilpotent element f is part of an \mathfrak{sl}_2 -triple $\mathfrak{s} = \{e, h, f\}$ in \mathfrak{g} , and by Kostant's Theorem [CMG93, Thm. 3.4.10] all \mathfrak{sl}_2 -triples containing f are conjugate by the centralizer of f in G, the adjoint group for \mathfrak{g} . It follows that all the constructions of this paper depend only on the G-orbit of f, and not on the chosen \mathfrak{sl}_2 -triple.

We have the ad(h/2)-eigenspace decomposition

$$\mathfrak{g} = \bigoplus_{k \in \frac{1}{2}\mathbb{Z}} \mathfrak{g}_k, \qquad \mathfrak{g}_k = \left\{ a \in \mathfrak{g} \mid [h, a] = 2ka \right\}.$$
 (2.1)

The largest $d \in \frac{1}{2}\mathbb{Z}$ such that $\mathfrak{g}_d \neq 0$ is called the *depth* of f. Note that $\dim \mathfrak{g}_k = \dim \mathfrak{g}_{-k}$ by \mathfrak{sl}_2 -representation theory, and $d \geq 1$ since $f \in \mathfrak{g}_{-1}$. For $j \in \frac{1}{2}\mathbb{Z}$, we shall use the notation $\mathfrak{g}_{>j} := \bigoplus_{k>j} \mathfrak{g}_k$, and similarly for $\mathfrak{g}_{\geq j}$, $\mathfrak{g}_{< j}$ and $\mathfrak{g}_{\leq j}$.

Remark 2.1. The depth d of f is easy to compute by knowing the Dynkin characteristic of f, defined as follows. Choose a Cartan subalgebra $\mathfrak h$ of $\mathfrak g$ contained in $\mathfrak g_0$ and choose a set of simple roots α_1,\ldots,α_r of $\mathfrak g$ such that $\alpha_i(h)\geq 0$, for all $i=1,\ldots,r$. If $\mathfrak g$ is simple and θ is the corresponding highest root, then $d=\frac{1}{2}\theta(h)$. More concretely, if $\theta=\sum_{i=1}^r a_i\alpha_i$, then

$$d = \frac{1}{2} \sum_{i=1}^{r} a_i \alpha_i(h).$$

In the general case of a reductive Lie algebra \mathfrak{g} the depth of f is equal to the maximum of the depths over all simple components of \mathfrak{g} . Recall from [CMG93, §3.5] that $\frac{1}{2}\alpha_i(h)$ can have only the values 0, 1/2, and 1; the collection of these numbers is called the Dynkin characteristic of f. Traditionally, the Dynkin characteristic is the collection of integers $\alpha_i(h)$, however, in the theory of W-algebras it is more natural to consider the halves of these integers.

Let $(\cdot|\cdot)$ be a non-degenerate symmetric invariant bilinear form on \mathfrak{g} . The subspace $\mathfrak{g}_{1/2}$ carries a skew-symmetric bilinear form ω , defined by

$$\omega(a,b) = (f|[a,b]), \quad a,b \in \mathfrak{g}_{1/2}.$$
 (2.2)

It is non-degenerate since ad $f: \mathfrak{g}_{1/2} \to \mathfrak{g}_{-1/2}$ is an isomorphism (by \mathfrak{sl}_2 -representation theory).

Definition 2.2.

(a) An element of g of the form

$$f + E$$
, where $E \in \mathfrak{g}_d \setminus \{0\}$, (2.3)

is called a *cyclic* element attached to the nilpotent element f.

(b) An element of g of the form

$$f + E$$
, where $E \in \mathfrak{g}_{d-\frac{1}{2}} \setminus \{0\}$, (2.4)

is called a *quasi-cyclic* element attached to the nilpotent element f, if the centralizer of E in $\mathfrak{g}_{1/2}$ is coisotropic with respect to the skew-symmetric form ω (i.e., its orthocomplement in $\mathfrak{g}_{1/2}$ is isotropic).

The following lemma is obvious.

Lemma 2.3.

- (a) The subspace \mathfrak{g}_d lies in the center of the subalgebra $\mathfrak{g}_{>0}$.
- (b) Let $E \in \mathfrak{g}_{d-\frac{1}{2}}$ and let \mathfrak{l}^{\perp} be the centralizer of E in $\mathfrak{g}_{1/2}$. Then E lies in the center of the subalgebra $\mathfrak{n} := \mathfrak{l}^{\perp} \oplus \mathfrak{g}_{\geq 1}$.

2.2. Classification of cyclic elements

The classification for reductive \mathfrak{g} easily reduces to the case when \mathfrak{g} is simple [EKV13], which we will assume in this subsection. We shall often use the well-known fact that an element of a reductive subalgebra in a reductive Lie algebra \mathfrak{g} is semisimple (resp. nilpotent) if and only if it is semisimple (resp. nilpotent) in \mathfrak{g} .

Definition 2.4 ([EKV13]). The nilpotent element f is called of *semisimple type* if there exists a cyclic element attached to f which is semisimple, and it is called of *nilpotent type* if all cyclic elements attached to f are nilpotent.

Let $\mathfrak{z}(\mathfrak{s})$ (resp. $Z(\mathfrak{s})$) be the centralizer of the \mathfrak{sl}_2 -triple \mathfrak{s} in \mathfrak{g} (resp. in the adjoint group G). More generally, for a subalgebra \mathfrak{q} of \mathfrak{g} we denote by $\mathfrak{z}(\mathfrak{q})$ its centralizer in \mathfrak{g} . Note that if a subalgebra \mathfrak{q} of \mathfrak{g} is normalized by \mathfrak{s} , then the ad x-grading of \mathfrak{g} induces that of \mathfrak{q} by letting $\mathfrak{q}_j = \mathfrak{q} \cap \mathfrak{g}_j$.

Theorem 2.5 ([EKV13]).

- (a) The nilpotent element f is of nilpotent type if and only if the depth of f is not an integer.
- (b) If the cyclic element f + E, $E \in \mathfrak{g}_d$, is semisimple, then the $Z(\mathfrak{s})$ -orbit of E in \mathfrak{g}_d is closed. If f is of semisimple type, then the set $\{E \in \mathfrak{g}_d \mid f + E \text{ is semisimple}\}$ contains a non-empty Zariski open subset.
- (c) If f is not of nilpotent type, then there exists a semisimple subalgebra $\mathfrak{q} \subset \mathfrak{g}$ normalized by \mathfrak{s} , such that $\mathfrak{q}_d \neq 0$, and an element $f^s \in \mathfrak{q}_{-1}$ of semisimple type in \mathfrak{q} , for which the element $f^n := f f^s$ lies in $\mathfrak{z}(\mathfrak{q})$. Consequently, by (b), there exists $E \in \mathfrak{q}_d$, such that $f^s + E$ is semisimple (in \mathfrak{q} , hence in \mathfrak{g}) and $[f + E, f^n] = 0$.

Proof. Claims (a) and (b) are proved in [EKV13, Thm.1.1 & Prop.2.2]. Claim (c) follows from discussions in [EKV13, Sect. 4] for classical Lie algebras, and [EKV13, Sect. 5] for exceptional Lie algebras. \Box

2.3. Quasi-cyclic elements

The classification of quasi-cyclic elements in reductive Lie algebras is discussed in [DSJKV20]. In this subsection we shall discuss only the quasi-cyclic elements in a simple Lie algebra \mathfrak{g} , associated to a nilpotent element f of nilpotent type.

Theorem 2.6 ([EKV13]).

- (a) There are no nilpotent elements of nilpotent type in \mathfrak{sl}_n and \mathfrak{sp}_n .
- (b) All nilpotent elements of nilpotent type in son correspond to partitions for which the largest part p₁ is odd and has multiplicity 1, and the next part p₂ equals p₁ − 1 and has even multiplicity.
- (c) All nilpotent elements f of nilpotent type in exceptional Lie algebras and their depths are listed in Table 1.

Proof. Part (a) and (b) follow from [EKV13, Sect. 4]. The list of Table 1 is [EKV13, Table 1.1]. \Box

Theorem 2.7 ([DSJKV20]).

- (a) For all nilpotent elements of nilpotent type in \mathfrak{so}_n there exists a non-nilpotent quasi-cyclic element.
- (b) There are no non-nilpotent quasi-cyclic elements for the following nilpotent elements of nilpotent type (see Table 1): $4A_1$, $2A_2 + 2A_1$, $2A_3$, $A_4 + A_3$ and A_7 in E_8 ; $\widetilde{A}_2 + A_1$ in F_4 ; \widetilde{A}_1 in G_2 .
- (c) For all other nilpotent elements of nilpotent type in exceptional simple Lie algebras there exists a non-nilpotent quasi-cyclic element (see Table 1).

Proof. (a) follows from Example 2.12 and Remark 2.15 below. (b) and (c) follow from the discussion preceding Table 1, Example 2.13 and Remark 2.15 below. \Box

Example 2.8. Let $\mathfrak{g} = G_2$ and let f be the nilpotent element denoted by \widetilde{A}_1 as in the last row in Table 1. This is a nilpotent element of nilpotent type, so all cyclic elements are nilpotent. We claim that there are no quasi-cyclic elements attached to f. Recall that the set of positive roots for \mathfrak{g} is $R_+ = \{\alpha, \beta, \alpha + \beta, \alpha + 2\beta, \alpha + 3\beta, 2\alpha + 3\beta\}$, where α and β are simple roots and β is a short root. Choose root vectors e_{γ} , $\gamma \in R = R_+ \cup (-R_+)$. Then, for the grading (2.1), we have: $\deg e_{\alpha} = 0$, $\deg e_{\beta} = 1/2$ (see Table 1), so that, for this grading,

$$\mathfrak{g}_{1/2} = \mathbb{F}e_{\beta} \oplus \mathbb{F}e_{\alpha+\beta}, \quad \mathfrak{g}_1 = \mathbb{F}e_{\alpha+2\beta}, \quad \mathfrak{g}_{3/2} = \mathbb{F}e_{\alpha+3\beta} \oplus \mathbb{F}e_{2\alpha+3\beta}.$$

The centralizer of $e = e_{\alpha+2\beta}$ in $\mathfrak{g}_{1/2}$ is zero, hence it is not coisotropic in $\mathfrak{g}_{1/2}$. Note that $\mathfrak{g}_{-1} = \mathbb{F}e_{-\alpha-2\beta}$, so that $f = e_{-\alpha-2\beta}$ is a short root vector.

Example 2.9. Let \mathfrak{g} be a simple Lie algebra, different from \mathfrak{sp}_n , and let $f = e_{-\theta}$, the lowest root vector. In this case, the depth is d = 1, and there exists a unique $E \in \mathfrak{g}_{1/2}$, up to a non-zero constant factor and action of $Z(\mathfrak{s})$, such that f + E is a semisimple quasi-cyclic element, see [DSKV14, Prop. 8.4].

In Table 1 we list the Dynkin characteristics of all nilpotent elements f of nilpotent type in exceptional simple Lie algebras \mathfrak{g} (the notation is the same as in [CMG93]). In the third column we list the depth d of f. In the fourth column we list the image in $\operatorname{End}\mathfrak{g}_{d-\frac{1}{2}}$ of the action of the Lie algebra $\mathfrak{z}(\mathfrak{s})$ on $\mathfrak{g}_{d-\frac{1}{2}}$, and in the fifth the rank of this action, defined as $\dim \left(\mathfrak{g}_{d-\frac{1}{2}} \ \middle | Z(\mathfrak{s})\right)$. Here for a linear reductive group G|V we use the standard notation $V \ \middle | G = \operatorname{Spec} \mathbb{F}[V]^G$. Finally, the last column says whether there exists a quasi-cyclic element attached to f which is semisimple or non-nilpotent. It follows from Table 1 that there are no quasi-cyclic elements attached to f if and only if $\dim \mathfrak{g}_{d-\frac{1}{3}} = 1$.

By $\mathfrak{st}_{\mathfrak{a}}$ we denote the standard representation of the Lie algebra \mathfrak{a} (it is 26-dimensional for $\mathfrak{a}=F_4$ and 7-dimensional for $\mathfrak{a}=G_2$); 1 stands for the trivial 1-dimensional representation; \oplus stands for the direct sum of linear Lie algebras: $(a_1 \oplus \mathfrak{a}_2)|(V_1 \oplus V_2)$.

Table 1. Quasi-cyclic elements attached to nilpotent elements of nilpotent type in exceptional simple Lie algebras

\mathfrak{g}	nilpotent f		d	$ \mathfrak{z}(\mathfrak{s}) \mathfrak{g}_{d-rac{1}{2}}$	rank	$f+E,E\in\mathfrak{g}_{d-\frac{1}{2}}$
$\overline{\mathrm{E}_{6}}$	$3A_1$	$0 - 0 - \frac{1}{2} - 0 - 0$	$\frac{3}{2}$	$\mathrm{ad}_{\mathfrak{sl}_3} \oplus 1$	3	∃ semisimple
	$2A_2 + A_1$	$\begin{array}{c} 0 \\ \frac{1}{2} - 0 - \frac{1}{2} - 0 - \frac{1}{2} \\ 0 \end{array}$	$\frac{5}{2}$	$\operatorname{st}_{\mathfrak{so}_3} \oplus 1$	2	\exists non-nilpotent
$\overline{\mathrm{E}_7}$	3A ₁	$0 - 0 - 0 - 0 - \frac{1}{2} - 0$	$\frac{3}{2}$	$\wedge^2 \mathrm{st}_{\mathfrak{sp}_6}$	3	\exists semisimple
	$4A_1$	$\frac{1}{2} - 0 - 0 - 0 - 0 - 0 - 0$	$\frac{3}{2}$	$\wedge^2\mathrm{st}_{\mathfrak{sp}_6}\oplus1$	4	\exists semisimple
	$2A_2 + A_1$	$0 - \frac{1}{2} - 0 - 0 - \frac{1}{2} - 0$	$\frac{5}{2}$	$\operatorname{st}_{\mathfrak{so}_3} \oplus \operatorname{st}_{\mathfrak{so}_3}$	2	\exists non-nilpotent
$\overline{\mathrm{E}_8}$	$3A_1$	$0 - \frac{1}{2} - 0 - 0 - 0 - 0 - 0$	$\frac{3}{2}$	$\mathrm{st}_{F_4}\oplus 1$	3	∃ semisimple
	$ $ 4A $_1$	0 - 0 - 0 - 0 - 0 - 0 - 0 - 0	$\frac{3}{2}$	$\wedge^2 \mathrm{st}_{\mathfrak{sp}_8}$	4	nilpotent only
	$2A_2 + A_1$	$0 - \frac{1}{2} - 0 - 0 - 0 - 0 - \frac{1}{2}$	$\frac{5}{2}$	$\operatorname{st}_{G_2} \oplus \operatorname{st}_{\mathfrak{so}_3}$	2	\exists non-nilpotent
	$2A_2 + 2A_1$	$0 - 0 - 0 - \frac{1}{2} - 0 - 0 - 0$	$\frac{5}{2}$	$\mathrm{ad}_{\mathfrak{so}_5}$	2	nilpotent only
	$2A_3$	$0 - 0 - 0 - \frac{1}{2} - 0 - 0 - \frac{1}{2}$	$\frac{7}{2}$	$\operatorname{st}_{\mathfrak{so}_5} \oplus 1$	2	nilpotent only
	$A_4 + A_3$	$0 - \frac{1}{2} - 0 - 0 - \frac{1}{2} - 0 - 0$	$\frac{9}{2}$	$\mathrm{st}_{\mathfrak{so}_3}$	1	nilpotent only
	A_7	$0 - \frac{1}{2} \cdot \frac{1}{2} - 0 - \frac{1}{2} - 0 - \frac{1}{2}$	$\frac{15}{2}$	1	1	never quasi-cyclic
$\overline{\mathrm{F}_4}$	$A_1 + \widetilde{A}_1$	$0 - \frac{1}{2} \Rightarrow 0 - 0$	$\frac{3}{2}$	$S^2\operatorname{st}_{\mathfrak{so}_3}$	3	∃ semisimple
	$A_2 + A_1$	$0 - \frac{1}{2} \Rightarrow 0 - \frac{1}{2}$	$\frac{5}{2}$	$\mathrm{st}_{\mathfrak{so}_3}$	1	nilpotent only
G_2	$ \widetilde{A}_1 $	$0 \Longrightarrow \frac{1}{2}$	$\frac{3}{2}$	1	1	never quasi-cyclic

The results of Table 1 are obtained as follows (see [DSJKV20] for details). Denote $\mathfrak{g}_{d-\frac{1}{2}}$ by \mathfrak{m} . First, we compute $\mathfrak{z}(\mathfrak{s})|\mathfrak{m}$ using the SLA package in the GAP computer algebra system. We note that, by Table 1, $Z(\mathfrak{s})|\mathfrak{m}$ is a polar representation with a Cartan subspace \mathfrak{m}_0 (see [DK85] for the definitions). In particular, any closed orbit of $Z(\mathfrak{s})$ in \mathfrak{m} intersects \mathfrak{m}_0 non-trivially [DK85]. It follows that if \mathfrak{m}_0 contains no elements E_0 , such that $f+E_0$ is quasi-cyclic, then \mathfrak{m} contains no elements E, such that f+E is non-nilpotent quasi-cyclic. Indeed, in the contrary case the orbit of minimal dimension in $Z(\mathfrak{s})E$ is closed and non-zero, hence there exists $E_0 \in \mathfrak{m}_0$, which lies in this orbit, such that $f+E_0$ is quasi-cyclic. This allows us to restrict consideration of quasi-cyclicity of f+E to $E \in \mathfrak{m}_0$. Then, one uses again SLA as well as Mathematica to study non-nilpotence or semisimplicity of quasi-cyclic elements f+E with $E \in \mathfrak{m}_0$ through minimal polynomials of their matrices in a faithful representation of a semisimple subalgebra containing them.

2.4. Integrable triples

Let \mathfrak{g} be a finite-dimensional Lie algebra, let $(\cdot | \cdot)$ be a non-degenerate symmetric invariant bilinear form on \mathfrak{g} , let $f \in \mathfrak{g}$ be a non-zero nilpotent element contained in an \mathfrak{sl}_2 -triple $\mathfrak{s} \subset \mathfrak{g}$, and consider the corresponding $\frac{1}{2}\mathbb{Z}$ -grading (2.1).

Definition 2.10. An integrable triple associated to f is (f_1, f_2, E) , where $f_1, f_2 \in \mathfrak{g}_{-1}$ and $E \in \mathfrak{g}_{\geq \frac{1}{2}}$ is a non-zero homogeneous element, such that the following three properties hold:

- (i) $f = f_1 + f_2$ and $[f_1, f_2] = 0$,
- (ii) $[E, \mathfrak{g}_{\geq 1}] = 0$ and the centralizer of E in $\mathfrak{g}_{\frac{1}{2}}$ is coisotropic with respect to the bilinear form (2.2).
- (iii) $f_1 + E$ is semisimple and $[f_2, E] = 0$.

In this case E is called an *integrable element* for f.

Note that for an integrable triple (f_1, f_2, E) the decomposition $f + E = (f_1 + E) + f_2$ is a Jordan decomposition of f + E, and that E is a central element of the subalgebra $\mathfrak{n} := \mathfrak{l}^{\perp} \oplus \mathfrak{g}_{>1}$,

Theorem 2.11. Let f be a nilpotent element of integer depth d of a reductive Lie algebra \mathfrak{g} .

- a) If f is of semisimple type, then there exists $E \in \mathfrak{g}_d$, such that (f, 0, E) is an integrable triple in \mathfrak{g} .
- b) If f^s , f^n and $E \in \mathfrak{g}_d$ are the elements, constructed in Theorem 2.5(c), then (f^s, f^n, E) is an integrable triple in \mathfrak{g} .

Proof. It follows immediately from Theorem 2.5(c). \Box

Example 2.12. Let $\mathfrak{g} = \mathfrak{so}_N$, $N \geq 7$, and let $f \in \mathfrak{g}$ be a nilpotent element of nilpotent type. Recall from [EKV13] that this nilpotent element corresponds to a partition $\underline{p} = (p_1^{r_1}, p_2^{r_2}, \dots, p_n^{r_n})$ of N (namely, $N = r_1 p_1 + \dots + r_n p_n$, where $p_1 > p_2 > \dots > p_n > 0$, $r_i \geq 1$, for $1 \leq i \leq n$), and $p_1 = p + 1$ occurs with multiplicity $r_1 = 1$, $p_2 = p$, and p is even (this implies that r_2 is also even). Then, there exists an integrable triple (f_1, f_2, E) associated to f, with $E \in \mathfrak{g}_{d-\frac{1}{2}}$, constructed as follows.

Let $V = \bigoplus_{\alpha \in I} \mathbb{F}e_{\alpha} \cong \mathbb{F}^N$ be the N-dimensional vector space with basis e_{α} , $\alpha \in I = \{(a,i,j)|1 \leq a \leq n, 1 \leq i \leq r_a, 1 \leq j \leq p_a\}$. Consider the following involution on the set I

$$(a, i, j)' = (a, r_a + 1 - i, p_a + 1 - j), \tag{2.5}$$

and also define

$$\epsilon_{(a,i,j)} = \begin{cases} (-1)^{(i-1)p_a+j}, & 1 \le i \le \lceil \frac{r_a}{2} \rceil, \\ -(-1)^{(i-1)p_a+j+r_a}, & \lceil \frac{r_a}{2} \rceil + 1 \le i \le r_a, \end{cases}$$
(2.6)

where $[\cdot]$ denotes the ceiling function.

Define a non-degenerate bilinear form on V on basis elements as follows:

$$\langle e_{\alpha}|e_{\beta}\rangle = -\delta_{\alpha,\beta'}\epsilon_{\alpha}, \qquad \alpha,\beta \in I.$$
 (2.7)

Let A^{\dagger} denote the adjoint of $A \in \text{End } V$ with respect to (2.7). Explicitly, in terms of elementary matrices, it is given by:

$$(E_{\alpha\beta})^{\dagger} = \epsilon_{\alpha} \epsilon_{\beta} E_{\beta'\alpha'}. \tag{2.8}$$

It follows from equations (2.5) and (2.6) that $\epsilon_{\alpha} = \epsilon_{\alpha'}$, for every $\alpha \in I$. Hence, the bilinear form (2.7) is symmetric and $\mathfrak{g} = \{A \in \text{End } V \mid A^{\dagger} = -A\} \simeq \mathfrak{so}_N$.

Consider the element

$$f = \sum_{(a,i,j)\in I, j\neq p_a} E_{(a,i,j+1),(a,i,j)} \in \mathfrak{g}.$$

It is a nilpotent element in the Jordan form corresponding to the partition \underline{p} . We include f in the \mathfrak{sl}_2 -triple $\{e, h, f\} \subset \mathfrak{g}$, where

$$h = \sum_{(a,i,j)\in I} (p_a + 1 - 2j) E_{(a,i,j),(a,i,j)},$$

$$e = \sum_{(a,i,j)\in I, j\neq p_a} j(p_a - j) E_{(a,i,j),(a,i,j+1)}.$$
(2.9)

For $\alpha, \beta \in I$, we define the following elements of \mathfrak{g} :

$$F_{\alpha\beta} = E_{\alpha\beta} - \epsilon_{\alpha}\epsilon_{\beta}E_{\beta'\alpha'} \ (= -F_{\alpha\beta}^{\dagger}).$$

The following commutation relations hold $(\alpha, \beta, \gamma, \eta \in I)$:

$$[F_{\alpha\beta}, F_{\gamma\eta}] = \delta_{\gamma\beta}F_{\alpha\eta} - \delta_{\eta\alpha}F_{\beta\gamma} - \epsilon_{\alpha}\epsilon_{\beta}\delta_{\alpha'\gamma}F_{\beta'\eta} + \epsilon_{\alpha}\epsilon_{\beta}\delta_{\eta\beta'}F_{\gamma\alpha'}. \tag{2.10}$$

Using (2.10) and the explicit form of $h \in \mathfrak{g}$ given in (2.9), we have that the depth of the decomposition (2.1) is d = (2p-1)/2 and

$$\begin{split} \mathfrak{g}_{1/2} &= \mathrm{Span}_{\mathbb{F}} \{ F_{(1,1,j),(2,i,j)} \mid 1 \leq i \leq r_2, 1 \leq j \leq p \} \\ &\oplus \mathrm{Span}_{\mathbb{F}} \left\{ F_{(a,i,i'),(b,j,j')} \mid a,b \geq 2, \frac{p_a - p_b}{2} - i' + j' = \frac{1}{2} \right\} = U \oplus W. \end{split}$$

Consider the element $E = F_{(1,1,1)(1,1,p)} \in \mathfrak{g}_{p-1} = \mathfrak{g}_{d-\frac{1}{2}}$. Clearly, $[E,\mathfrak{g}_{\geq 1}] = 0$. Moreover, by the commutation relations (2.10) we also have that [E,W] = 0. Let

$$A = \sum_{i=1}^{r_1} \sum_{j=1}^{p} a_{ij} F_{(1,1,j),(2,i,j)} \in U.$$

Using the commutation relations (2.10) we have

$$[E, A] = \sum_{i=1}^{r_2} a_{ip} F_{(1,1,p),(2,i,p)}.$$

This implies that the centralizer of E in $\mathfrak{g}_{1/2}$ is

$$\mathfrak{z}(E) \cap \mathfrak{g}_{1/2} = \operatorname{Span}_{\mathbb{F}} \{ F_{(1,1,j),(2,i,j)} \mid 1 \le i \le r_2, 1 \le j \le p-1 \} \oplus W.$$
 (2.11)

We claim that this is a coisotropic subspace of $\mathfrak{g}_{1/2}$, namely $\mathfrak{z}(E) \cap \mathfrak{g}_{1/2} = \mathfrak{l}^{\perp}$ for some isotropic subspace $\mathfrak{l} \subset \mathfrak{g}_{1/2}$. In fact, consider the following subspace of $\mathfrak{g}_{1/2}$:

$$\mathfrak{l} = \operatorname{Span}_{\mathbb{F}} \{ F_{(1,1,1),(2,i,1)} \mid 1 \le i \le r_2 \}.$$

Since, by the commutation relations (2.10) and the fact that $F_{(1,1,1),(1,1,p+1)} = 0$, we have that $[\mathfrak{l},\mathfrak{l}] = 0$, it follows that \mathfrak{l} is an isotropic subspace with respect to the skew-symmetric bilinear form given in (2.2) (we assume that the bilinear form on \mathfrak{g} in Section 2.1 is the trace form). Moreover, from (2.10) we also have that $[\mathfrak{l},W] = 0$. Hence, $W \subset \mathfrak{l}^{\perp}$. Using again the commutation relations (2.10) we get $(1 \leq i,j \leq r_2,1 \leq k \leq p)$

$$[F_{(1,1,1),(2,i,1)}, F_{(1,1,k),(2,i,k)}] = \epsilon_{(2,i,1)}\delta_{i+j,r_2+1}\delta_{k,p}F_{(1,1,k),(1,1,p+1)}.$$

It follows that

$$\omega(F_{(1,1,1),(2,i,1)},F_{(1,1,k),(2,j,k)}) = (f|[F_{(1,1,1),(2,i,1)},F_{(1,1,k),(2,j,k)}]) \neq 0$$

if and only if $i+j=r_2+1$ and k=p. Hence, $\mathfrak{z}(E)\cap\mathfrak{g}_{1/2}=\mathfrak{l}^{\perp}$ (cf. equation (2.11)) thus showing that f+E is a quasi-cyclic element attached to f.

Let $f_1 = \sum_{j=1}^p E_{(1,1,j+1),(1,1,j)} \in \mathfrak{g}$ and $f_2 = f - f_1$. It is immediate to check that $[f_1 + E, f_2] = 0$ and that f_2 lies in \mathfrak{g}_{-1} (hence, it is nilpotent). Moreover, the minimal polynomial of $f_1 + E$ is $p(\lambda) = \lambda(\lambda^p - 2)$. Hence, $f_1 + E$ is semisimple and (f_1, f_2, E) is an integrable triple associated to f.

Example 2.13. In simple Lie algebras of exceptional types there are three cases of nilpotent orbits of nilpotent type admitting non-nilpotent quasi-cyclic elements but no semisimple quasi-cyclic elements. Namely, these are nilpotent elements with label $2A_2 + A_1$ in E_6 , E_7 , and E_8 (see Table 1). For them we find integrable triples computationally, using the SLA package of the GAP system.

We choose the following representatives f:

$$\begin{array}{l} \text{for } \mathbf{E}_{6}, \, f_{00111} + f_{11100} + f_{11110} + f_{01111} + f_{01210}; \\ \text{for } \mathbf{E}_{7}, \, f_{011111} + f_{111111} + f_{011210} + f_{001221} + f_{112210}; \\ \text{for } \mathbf{E}_{8}, \, f_{0122221} + f_{1122211} + f_{1222210} + f_{0112321} + f_{1112321}. \end{array}$$

In all three cases, the zero weight space of the representation of $\mathfrak{z}(\mathfrak{s})$ on $\mathfrak{g}_{d-1/2}$ is a Cartan subspace, spanned by two root vectors b_1, b_2 , where

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\begin{array}{l} \text{for E}_6,\ b_1=e_{{}^{11211}}\ \text{and}\ b_2=e_{{}^{12221}};\\ \text{for E}_7,\ b_1=e_{{}^{122321}}\ \text{and}\ b_2=e_{{}^{123321}};\\ \text{for E}_8,\ b_1=e_{{}^{1234532}}\ \text{and}\ b_2=e_{{}^{1234542}}. \end{array}
```

In all three cases, an element $x_1b_1 + x_2b_2$ of this subspace has coisotropic centralizer in $\mathfrak{g}_{1/2}$ if and only if $x_2 = -x_1$. Take $E = c(b_1 - b_2)$, then in all three cases for any $c \neq 0$ the Jordan decomposition of f + E is $(f^s + E) + f^n$, where

for E₆,
$$f^n = f_{01210}$$
;
for E₇, $f^n = f_{001221}$;
for E₈, $f^n = f_{1222210}$.

Thus, since $f^n \in \mathfrak{g}_{-1}$, we have $f^s = f - f^n \in \mathfrak{g}_{-1}$. Hence, (f^s, f^n, E) is an integrable triple.

Additionally, let us remark that in all three of these cases f^s has label $2A_2$ and f^n has label A_1 . Note also that the subalgebra generated by f and E is the direct sum of an \mathfrak{sl}_3 and a 1-dimensional center spanned by f^n . Moreover f^s is a principal nilpotent in this \mathfrak{sl}_3 .

Theorem 2.14. Let f be a nilpotent element of depth d of nilpotent type in a simple Lie algebra \mathfrak{g} , such that there exists $E \in \mathfrak{g}_{d-\frac{1}{2}}$, for which f+E is a non-nilpotent quasi-cyclic element.

- (a) If f + E is semisimple, then (f, 0, E) is an integrable triple associated to f.
- (b) For f, not covered by (a), there exists an integrable triple.

Proof. Part (a) is clear. The cases, not covered by (a) are as follows. First, it is $\mathfrak{g} = \mathfrak{so}_N$, which is covered by Example 2.12. Second, it is $2A_2 + A_1$ in all algebras of type E, which is covered by Example 2.13. \square

Remark 2.15. Conversely, if (f_1, f_2, E) is an integrable triple for f, where $E \in \mathfrak{g}_d$ (resp. $E \in \mathfrak{g}_{d-\frac{1}{2}}$), then $f_1 + f_2 + E$ is a non-nilpotent cyclic (resp. quasi-cyclic) element. This follows from Definition 2.2.

Remark 2.16. It is clear from Example 2.8 that for $\mathfrak{g} = G_2$ and f a short root vector there are no integrable triples.

3. Poisson vertex algebras, Hamiltonian equations and integrability

3.1. PVA and Hamiltonian PDE

Recall (see, e.g., [BDSK09]) that a *Poisson vertex algebra* (abbreviated PVA) is a commutative associative algebra \mathcal{V} with a derivation ∂ , endowed with a λ -bracket $\mathcal{V} \otimes \mathcal{V} \to \mathcal{V}[\lambda]$, $a \otimes b \mapsto \{a_{\lambda}b\}$, satisfying the following axioms $(a, b, c \in \mathcal{V})$:

- (i) sesquilinearity: $\{\partial a_{\lambda}b\} = -\lambda \{a_{\lambda}b\}, \{a_{\lambda}\partial b\} = (\partial + \lambda)\{a_{\lambda}b\};$
- (ii) skew-symmetry: $\{a_{\lambda}b\} = -\{b_{-\lambda-\partial}a\}$ (where ∂ acts on the coefficients);
- (iii) Jacobi identity: $\{a_{\lambda}\{b_{\mu}c\}\} \{b_{\mu}\{a_{\lambda}c\}\} = \{\{a_{\lambda}b\}_{\lambda+\mu}c\};$
- (iv) left Leibniz rule: $\{a_{\lambda}bc\} = \{a_{\lambda}b\}c + b\{a_{\lambda}c\}.$

As a consequence of skew-symmetry and the left Leibniz rule, we also have

(iv') right Leibniz rule: $\{ab_{\lambda}c\} = \{a_{\lambda+\partial}c\}_{\to}b + \{b_{\lambda+\partial}c\}_{\to}a$, where \to means that ∂ is moved to the right.

Example 3.1. The most important example for this paper will be the affine PVA $\mathcal{V}(\mathfrak{g}, E)$, where \mathfrak{g} is a Lie algebra with a symmetric invariant bilinear form $(\cdot | \cdot)$ and $E \in \mathfrak{g}$. It is defined as the differential algebra $\mathcal{V}(\mathfrak{g}) = S(\mathbb{F}[\partial]\mathfrak{g})$, the algebra of differential polynomials over the vector space \mathfrak{g} , with the PVA λ -bracket given by

$$\{a_{\lambda}b\}_{z} = [a,b] + (a|b)\lambda + z(E|[a,b]), \quad a,b \in \mathfrak{g}, z \in \mathbb{F}, \tag{3.1}$$

and extended to \mathcal{V} by sesquilinearity axioms and the Leibniz rules.

Let \mathcal{V} be a PVA. We denote by $\int : \mathcal{V} \to \mathcal{V}/\partial \mathcal{V}$ the canonical quotient map of vector spaces. Recall that (see [BDSK09]) $\mathcal{V}/\partial \mathcal{V}$ carries a well-defined Lie algebra structure given by

$$\{ \int f, \int g \} = \int \{ f_{\lambda}g \} |_{\lambda=0},$$

and we have a representation of the Lie algebra $V/\partial V$ on V given by

$$\{\int f, g\} = \{f_{\lambda}g\}|_{\lambda=0}.$$

A Hamiltonian equation on $\mathcal V$ associated to a Hamiltonian functional $\int h \in \mathcal V/\partial \mathcal V$ is the evolution equation

$$\frac{du}{dt} = \{ \int h, u \}, \quad u \in \mathcal{V}. \tag{3.2}$$

The minimal requirement for integrability is to have an infinite collection of linearly independent integrals of motion in involution:

$$\int h_0 = \int h, \int h_1, \int h_2, \dots$$
 such that $\{\int h_m, \int h_n\} = 0$ for all $m, n \in \mathbb{Z}_{\geq 0}$.

In this case, we have the *integrable hierarchy* of Hamiltonian equations

$$\frac{du}{dt_n} = \{ \int h_n, u \}, \quad n \in \mathbb{Z}_{\geq 0}.$$
(3.3)

3.2. Bi-Poisson vertex algebras and Lenard–Magri scheme of integrability

Let $\{\cdot_{\lambda}\cdot\}_0$ and $\{\cdot_{\lambda}\cdot\}_{\infty}$ be two λ -brackets on the same differential algebra \mathcal{V} . We can consider the pencil of λ -brackets

$$\{\cdot_{\lambda}\cdot\}_z = \{\cdot_{\lambda}\cdot\}_0 + z\{\cdot_{\lambda}\cdot\}_{\infty} \quad , \qquad z \in \mathbb{F}. \tag{3.4}$$

As above, we say that \mathcal{V} is a bi-PVA if $\{\cdot_{\lambda}\cdot\}_z$ is a PVA λ -bracket on \mathcal{V} for every $z \in \mathbb{F}$. The affine PVA defined in Example 3.1 is in fact a bi-PVA.

Let \mathcal{V} be a bi-Poisson vertex algebra with λ -brackets $\{\cdot_{\lambda}\cdot\}_0$ and $\{\cdot_{\lambda}\cdot\}_{\infty}$. A bi-Hamiltonian equation is an evolution equation which can be written in Hamiltonian form with respect to both PVA λ -brackets and two Hamiltonian functionals $\int h_0$, $\int h_1 \in \mathcal{V}/\partial \mathcal{V}$:

$$\frac{du}{dt} = \{ \int h_0, u \}_0 = \{ \int h_1, u \}_{\infty}, \ u \in \mathcal{V}.$$

The most common way to prove integrability for a bi-Hamiltonian equation is to solve the so called *Lenard–Magri recurrence relation* (see [Mag78]):

$$\{\int h_n, \int u\}_0 = \{\int h_{n+1}, \int u\}_\infty, \quad n \in \mathbb{Z}_{\geq 0}, \ u \in \mathcal{V}.$$
 (3.5)

The Lenard–Magri recurrence relation (3.5) produces local functionals in involution. In order to prove this claim, one needs the following simple lemma (cf. [Mag78], [BDSK09]).

Lemma 3.2. Let \mathcal{U} be a vector space with two skew-commutative brackets $\{\cdot,\cdot\}_0$ and $\{\cdot,\cdot\}_{\infty}$ (not necessarily satisfying Jacobi identity). Let h_0,h_1,\ldots,h_N , where $N \geq 1$, be a sequence of elements of \mathcal{U} , satisfying the relation

$${h_n, u}_0 = {h_{n+1}, u}_{\infty}, \quad n = 0, \dots, N-1, u \in \mathcal{U}.$$
 (3.6)

Then

- a) $\{h_m, h_n\}_0 = 0 = \{h_m, h_n\}_{\infty}$, for all $m, n = 0, \dots, N$.
- b) If $N = \infty$, $\{h_0, \mathcal{U}\}_{\infty} = 0$ and $\{g_n\}_{n \in \mathbb{Z}_{\geq 0}}$ is another sequence satisfying (3.6), then $\{h_m, g_n\}_0 = 0 = \{h_m, g_n\}_{\infty}$, for every $m, n \in \mathbb{Z}_{\geq 0}$.

Proof. Due to skew-symmetry of the brackets, part a) holds for m = n. Without loss of generality we may assume that m > n and prove the result by induction on m - n. Then, by equation (3.6) for $u = h_m$ and skew-symmetry we have $\{h_m, h_n\}_0 = \{h_m, h_{n+1}\}_{\infty} = 0$ by the inductive assumption. Similarly, by equation (3.6) we have $\{h_m, h_n\}_{\infty} = \{h_{m-1}, h_n\}_0 = 0$. This proves part a).

We prove part b) by induction on $m \in \mathbb{Z}_{\geq 0}$. For m = 0 we have, by our assumption on h_0 , $\{h_0, g_n\}_{\infty} = 0$, for every $n \in \mathbb{Z}_{\geq 0}$. Furthermore, by our assumption on h_0 and skew-symmetry we have $\{h_0, g_m\}_0 = \{h_0, g_{m+1}\}_{\infty} = 0$. For $m \geq 1$ we have, by equation (3.6) and the inductive assumption, that $\{h_m, g_n\}_{\infty} = \{h_{m-1}, g_n\}_0 = 0$, for every $n \in \mathbb{Z}_{\geq 0}$. Similarly, using skew-symmetry and the Lenard-Magri recursion (3.6) for $\{g_n\}_{n \in \mathbb{Z}_{\geq 0}}$, we have $\{h_m, g_n\}_0 = \{h_m, g_{n+1}\}_{\infty} = 0$, thus concluding the proof of part b). \square

As a special case of Lemma 3.2, if \mathcal{V} is a bi-PVA and $\int h_0, \int h_1, \ldots \in \mathcal{V}/\partial \mathcal{V}$ satisfy the Lenard-Magri recurrence (3.5), then they are in involution:

$$\{\int h_m, \int h_n\}_0 = \{\int h_m, \int h_n\}_{\infty} = 0 \text{ for all } m, n \ge 0,$$

namely,

$$\mathcal{A} := \operatorname{Span} \{ \int h_n \}_{n=0}^{\infty} \subset \mathcal{V} / \partial \mathcal{V}$$

is an abelian subalgebra with respect to both Lie algebra brackets $\{\cdot,\cdot\}_0$ and $\{\cdot,\cdot\}_{\infty}$. In this way, we get the corresponding hierarchy of bi-Hamiltonian equations

$$\frac{du}{dt_n} = \{ \int h_n, u \}_0 = \{ \int h_{n+1}, u \}_{\infty}, \ n \in \mathbb{Z}_{\geq 0}, \ u \in \mathcal{V}.$$

If moreover

$$\{\int h_0, \int u\}_{\infty} = 0 \text{ for all } u \in \mathcal{V},$$
 (3.7)

and $\int g_0, \int g_1, \dots \in \mathcal{V}/\partial \mathcal{V}$ is any other sequence satisfying (3.5), then

$$\widetilde{\mathcal{A}} = \operatorname{Span}\{\int h_n, \int g_n\}_{n=0}^{\infty} \subset \mathcal{V}/\partial \mathcal{V}$$

is also an abelian subalgebra.

4. Classical affine W-algebras and generalized Drinfeld–Sokolov hierarchies

4.1. Definition of classical affine W-algebras

Let \mathfrak{g} be a reductive Lie algebra as in Section 2.1, $(\cdot | \cdot)$ a non-degenerate symmetric invariant bilinear form on \mathfrak{g} , f its non-zero nilpotent element, and consider the corresponding $\frac{1}{2}\mathbb{Z}$ -grading (2.1) and the skew-symmetric non-degenerate bilinear form ω on $\mathfrak{g}_{1/2}$ defined by (2.2). Fix an isotropic (with respect to ω) subspace $\mathfrak{l} \subset \mathfrak{g}_{1/2}$ and denote by $\mathfrak{l}^{\perp} = \{a \in \mathfrak{g}_{1/2} \mid \omega(a,b) = 0 \text{ for all } b \in \mathfrak{l}\} \subset \mathfrak{g}_{1/2}$ its orthogonal complement with respect to ω . Throughout the paper we consider the following nilpotent subalgebras of \mathfrak{g} :

$$\mathfrak{m} = \mathfrak{l} \oplus \mathfrak{g}_{>1} \subset \mathfrak{n} = \mathfrak{l}^{\perp} \oplus \mathfrak{g}_{>1},$$

where $\mathfrak{g}_{>1} = \bigoplus_{k>1} \mathfrak{g}_k$.

Fix an element $E \in \mathfrak{z}(\mathfrak{n})$ (the centralizer of \mathfrak{n} in \mathfrak{g}) and consider the affine bi-PVA $\mathcal{V}(\mathfrak{g}, E)$ from Example 3.1: it is the differential algebra $\mathcal{V}(\mathfrak{g}) = S(\mathbb{F}[\partial]\mathfrak{g})$, with its two PVA λ -brackets $\{\cdot_{\lambda}\cdot\}_0$ and $\{\cdot_{\lambda}\cdot\}_{\infty}$ (defined on $a, b \in \mathfrak{g}$ by $\{a_{\lambda}b\}_0 = [a, b] + (a|b)\lambda$ and $\{a_{\lambda}b\}_{\infty} = (E|[a, b])$, cf. equations (3.1) and (3.4)). Consider also the differential algebra ideal I, generated by the set

$$\langle m - (f|m) \rangle_{m \in \mathfrak{m}} \subset \mathcal{V}(\mathfrak{g}).$$

Note that, since $E \in \mathfrak{z}(\mathfrak{n})$, we have $\{a_{\lambda}w\}_{\infty} = 0$ for all $a \in \mathfrak{n}$ and $w \in \mathcal{V}(\mathfrak{g}, E)$. Consider the space

$$\widetilde{\mathcal{W}} = \{ w \in \mathcal{V}(\mathfrak{g}) \mid \{ a_{\lambda} w \}_0 \in I[\lambda], \text{ for every } a \in \mathfrak{n} \} \subset \mathcal{V}(\mathfrak{g}, E).$$

Lemma 4.1. $\widetilde{\mathcal{W}} \subset \mathcal{V}(\mathfrak{g}, E)$ is a bi-PVA subalgebra of $\mathcal{V}(\mathfrak{g}, E)$ and $I \subset \widetilde{\mathcal{W}}$ is a bi-PVA ideal.

Proof. Straightforward, see, e.g., [DSKV13, Sect. 3]. \square

Definition 4.2. The classical affine W-algebra associated to the triple (\mathfrak{g}, f, E) is the quotient

$$W(\mathfrak{g}, f, E) = \widetilde{W}/I, \tag{4.1}$$

which, by Lemma 4.1, has a natural structure of a bi-PVA, with the induced PVA λ -brackets $\{\cdot_{\lambda}\cdot\}_{0}$ and $\{\cdot_{\lambda}\cdot\}_{\infty}$.

Remark 4.3. Recall from [DSKV13] that the classical affine W-algebra depends only on the Lie algebra \mathfrak{g} and on the nilpotent orbit of f: for different choices of f in its nilpotent orbit, of the \mathfrak{sl}_2 -triple \mathfrak{s} containing f, and of the isotropic subspace $\mathfrak{l} \subset \mathfrak{g}_{1/2}$, we get isomorphic W-algebras.

Let $\mathfrak{p} \subset \mathfrak{g}$ be a subspace complementary to \mathfrak{m} in \mathfrak{g} : $\mathfrak{g} = \mathfrak{m} \oplus \mathfrak{p}$. We assume that \mathfrak{p} is compatible with the grading (2.1), so that $\mathfrak{g}_{\leq 0} \subset \mathfrak{p} \subset \mathfrak{g}_{\leq 1/2}$. Clearly, we can identify as differential algebras $\mathcal{V}(\mathfrak{g})/I \simeq \mathcal{V}(\mathfrak{p})$. Hence, the classical W-algebra can be viewed as a differential subalgebra of $\mathcal{V}(\mathfrak{p})$, the algebra of differential polynomials over \mathfrak{p} :

$$W(\mathfrak{g}, f, E) = \widetilde{W}/I \subset V(\mathfrak{g})/I \simeq V(\mathfrak{p}). \tag{4.2}$$

4.2. Integrable triples and generalized Drinfeld-Sokolov hierarchies

In [DSKV13] a generalized Drinfeld–Sokolov hierarchy was constructed, using the Lenard–Magri recurrence relation (3.5), under the assumption that the element $f + E \in \mathfrak{g}$ is semisimple. In this section we extend this result by constructing a generalized Drinfeld–Sokolov hierarchy for any integrable triple associated to f.

Let (f_1, f_2, E) be an integrable triple associated to f (cf. Definition 2.10), and let \mathfrak{l}^{\perp} be the centralizer of E in $\mathfrak{g}_{1/2}$, with \mathfrak{l} isotropic with respect to (2.2). Setting as above $\mathfrak{n} = \mathfrak{l}^{\perp} \oplus \mathfrak{g}_{>1}$, we have by definition that $[E, \mathfrak{n}] = 0$.

Let $\mathbb{K} = \mathbb{F}((z^{-1}))$ be the field of formal Laurent series in z^{-1} over \mathbb{F} . Consider the Lie algebra $\mathfrak{g}((z^{-1})) = \mathfrak{g} \otimes_{\mathbb{F}} \mathbb{K}$. Since, by Definition 2.10, $f_1 + E \in \mathfrak{g}$ is semisimple, $f_1 + zE \in \mathfrak{g}((z^{-1}))$ is also semisimple. Indeed, for a non-zero element $t \in \mathbb{K}$ we have the Lie algebra automorphism φ_t of $\mathfrak{g}((z^{-1}))$ acting as t^i on \mathfrak{g}_i . Hence, if $E \in \mathfrak{g}_k$, we have $t\varphi_t(f_1 + E) = f_1 + t^{k+1}E$. Then, working over the field extension of \mathbb{K} , containing $z^{1/(k+1)}$, we conclude that $f_1 + zE$ is semisimple in $\mathfrak{g}((z^{-1}))$. We thus have the direct sum decomposition

$$\mathfrak{g}((z^{-1})) = \mathfrak{h} \oplus \mathfrak{h}^{\perp}, \tag{4.3}$$

where

$$\mathfrak{h} := \operatorname{Ker} \operatorname{ad}(f_1 + zE) \text{ and } \mathfrak{h}^{\perp} := \operatorname{Im} \operatorname{ad}(f_1 + zE).$$
 (4.4)

The notation \mathfrak{h}^{\perp} relates to the fact that $\operatorname{Im} \operatorname{ad}(f_1 + zE)$ is the orthogonal complement of $\operatorname{Ker} \operatorname{ad}(f_1 + zE)$ with respect to the non-degenerate symmetric invariant bilinear form $(\cdot | \cdot)$ on $\mathfrak{g}((z^{-1}))$, extending the form $(\cdot | \cdot)$ on \mathfrak{g} by bilinearity. (But we will not use this fact.)

Lemma 4.4. For every element $A(z) \in \mathfrak{g}((z^{-1}))$, there exist unique $h(z) \in \mathfrak{h}$ and $U(z) \in \mathfrak{h}^{\perp}$ such that

$$h(z) + [f + zE, U(z)] = A(z).$$
 (4.5)

Proof. According to the direct sum decomposition (4.3), we can write, uniquely, A(z) = h(z) + B(z), where $h(z) \in \mathfrak{h}$ and $B(z) \in \mathfrak{h}^{\perp}$. Note that f + zE commutes with $f_1 + zE$, since $[f_2, f_1] = [f_2, E] = 0$ (cf. Definition 2.10). Hence $\mathfrak{h} = \operatorname{Ker} \operatorname{ad}(f_1 + zE)$ and $\mathfrak{h}^{\perp} = \operatorname{Im} \operatorname{ad}(f_1 + zE)$ are both $\operatorname{ad}(f + zE)$ -invariant. By definition, $f_1 + zE$ is the semisimple part of f + zE, and obviously $\operatorname{ad}(f_1 + zE)$ is invertible when restricted to \mathfrak{h}^{\perp} . Hence $\operatorname{ad}(f + zE)|_{\mathfrak{h}^{\perp}} : \mathfrak{h}^{\perp} \to \mathfrak{h}^{\perp}$ is invertible as well. Therefore, there exists a unique $U(z) \in \mathfrak{h}^{\perp}$ such that [f + zE, U(z)] = B(z). The claim follows. \square

If $E \in \mathfrak{g}_k$, we extend the $\frac{1}{2}\mathbb{Z}$ -grading (2.1) to $\mathfrak{g}((z^{-1}))$ by letting z have degree -k-1, so that f+zE and f_1+zE are homogeneous of degree -1. Then

$$\mathfrak{g}((z^{-1})) = \widehat{\bigoplus}_{i \in \frac{1}{2}\mathbb{Z}} \mathfrak{g}((z^{-1}))_i, \tag{4.6}$$

where $\mathfrak{g}((z^{-1}))_i \subset \mathfrak{g}((z^{-1}))$ is the space of homogeneous elements of degree i, and the direct sum is completed by allowing infinite series in positive degrees, cf. [DSKV13, Lem. 4.4(a)]. Note that, since $f_1 + zE$ is homogeneous, we have the corresponding decompositions of \mathfrak{h} and \mathfrak{h}^{\perp} :

$$\mathfrak{h} = \widehat{\bigoplus}_{i \in \frac{1}{2}\mathbb{Z}} \mathfrak{h}_i \text{ and } \mathfrak{h}^{\perp} = \widehat{\bigoplus}_{i \in \frac{1}{2}\mathbb{Z}} \mathfrak{h}_i^{\perp}.$$
 (4.7)

Recall that $\mathcal{V}(\mathfrak{p})$ is a commutative associative algebra with derivation ∂ . Consider the Lie algebra

$$\widetilde{\mathfrak{g}} = \mathbb{F}\partial \ltimes (\mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))),$$

where ∂ acts on the first factor. For every $U(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0} \subset \widetilde{\mathfrak{g}}$, we have a well-defined automorphism of the Lie algebra $\widetilde{\mathfrak{g}}$ given by $e^{\operatorname{ad} U(z)}$, cf. [DSKV13, Lem. 4.4(b)]. We extend the bilinear form $(\cdot | \cdot)$ on $\mathfrak{g}((z^{-1}))$ to a map

$$(\cdot|\cdot): \left(\mathcal{V}(\mathfrak{p})\otimes\mathfrak{g}((z^{-1}))\right)\times\left(\mathcal{V}(\mathfrak{p})\otimes\mathfrak{g}((z^{-1}))\right)\to\mathcal{V}(\mathfrak{p})((z^{-1})), \tag{4.8}$$

given by $(g \otimes a(z)|h \otimes b(z)) = gh(a(z)|b(z)).$

Using the bilinear form $(\cdot|\cdot)$ on \mathfrak{g} we get the isomorphism $\mathfrak{p}^* \simeq \mathfrak{m}^{\perp}$. Let $\{q_i\}_{i\in P}$ be a basis of \mathfrak{p} , and let $\{q^i\}_{i\in P}$ be the dual (with respect to $(\cdot|\cdot)$) basis of \mathfrak{m}^{\perp} , namely, such that $(q^j|q_i) = \delta_{ij}$. We denote

$$q = \sum_{i \in P} q_i \otimes q^i \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{m}^{\perp}. \tag{4.9}$$

The next result is a generalization of [DSKV13, Prop. 4.5] to the case of an integrable triple associated to the nilpotent element f.

Proposition 4.5. Let f be a non-zero nilpotent element of the Lie algebra \mathfrak{g} which admits an integrable triple (f_1, f_2, E) . Then there exist unique formal Laurent series $U(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>0}^{\perp}$ and $h(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>-1}$ such that

$$e^{\operatorname{ad} U(z)}(\partial + 1 \otimes (f + zE) + q) = \partial + 1 \otimes (f + zE) + h(z). \tag{4.10}$$

Moreover, an automorphism $e^{\operatorname{ad} U(z)}$, with $U(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0}$, solving (4.10) for some $h(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>-1}$, is defined uniquely up to multiplication on the left by automorphisms of the form $e^{\operatorname{ad} S(z)}$, $S(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>0}$.

Proof. Let us write $U(z) = \sum_{i \geq \frac{1}{2}} U_i(z)$, where $U_i(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_i^{\perp}$, $i \geq 1/2$, and $h(z) = \sum_{i \geq -\frac{1}{2}} h_i(z)$, where $h_i(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_i$, $i \geq -1/2$, and \mathfrak{h}_i , \mathfrak{h}_i^{\perp} are defined by (4.7). We determine $U_{i+1}(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{i+1}^{\perp}$ and $h_i(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_i$, inductively on $i \geq -1/2$, by equating the homogeneous components of degree i in each side of equation (4.10). This amounts to solving an equation in $h_i(z)$ and $U_{i+1}(z)$ of the form

$$h_i(z) + [1 \otimes (f + zE), U_{i+1}(z)] = A(z),$$
 (4.11)

where $A(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_i$ is a (differential polynomial) expression involving all the elements $U_{j+1}(z)$ and $h_j(z)$ for j < i (see the proof of [DSKV13, Prop.4.5(a)] for further details). By Lemma 4.4 there exists a unique solution $h_i(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_i$ and $U_{i+1}(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{i+1}^{\perp}$, thus proving the first part of the proposition.

Next, let $\widetilde{U}(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0}$, $\widetilde{h}(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>-1}$ be some other solution of (4.10): $e^{\operatorname{ad} \widetilde{U}(z)}(\partial + 1 \otimes (f + zE) + q) = \partial + 1 \otimes (f + zE) + \widetilde{h}(z)$. By the Baker-Campbell-Hausdorff formula, there exists $R(z) = \sum_{i>0}^{\infty} R_i(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0}$ such that $e^{\operatorname{ad} \widetilde{U}(z)} e^{-\operatorname{ad} U(z)} = e^{\operatorname{ad} R(z)}$. According to the direct sum decomposition (4.3), we can write, uniquely, R(z) = S(z) + T(z), where $S(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>0}$ and $T(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>0}^{\perp}$. To conclude the proof of the proposition we need to show that T(z) = 0. By construction, we have

$$\partial + 1 \otimes (f + zE) + \widetilde{h}(z) = e^{\operatorname{ad} R(z)} (\partial + 1 \otimes (f + zE) + h(z)). \tag{4.12}$$

Comparing the terms of degree -1/2 in both sides of the above equation, we get

$$[1 \otimes (f+zE), T_{1/2}(z)] = h_{-1/2}(z) - \widetilde{h}_{-1/2}(z) - [1 \otimes (f+zE), S_{1/2}(z)].$$

Since \mathfrak{h} and \mathfrak{h}^{\perp} are both $\mathrm{ad}(f+zE)$ -invariant, we have that $[1\otimes (f+zE), T_{1/2}(z)] \in (\mathcal{V}(\mathfrak{p})\otimes \mathfrak{h}_{-1/2})\cap (\mathcal{V}(\mathfrak{p})\otimes \mathfrak{h}_{-1/2}^{\perp})=0$. Hence $T_{1/2}(z)=0$, since $\mathrm{ad}(f+zE)|_{\mathfrak{h}^{\perp}}:\mathfrak{h}^{\perp}\to \mathfrak{h}^{\perp}$ is invertible. Let us assume, by induction, that $T_{j}(z)=0$ for all j< i. Comparing the terms of degree i-1 in both sides of equation (4.12), we easily get that, using the fact that \mathfrak{h} is a Lie subalgebra, $[1\otimes (f+zE),T_{i}(z)]\in (\mathcal{V}(\mathfrak{p})\otimes \mathfrak{h}_{i-1})\cap (\mathcal{V}(\mathfrak{p})\otimes \mathfrak{h}_{i-1})=0$, namely $T_{i}(z)=0$, as desired. \square

The main result of this section is the following theorem, which allows us to construct an integrable hierarchy of bi-Hamiltonian equations for classical affine W-algebras.

Theorem 4.6. Let \mathfrak{g} be a reductive Lie algebra with a non-degenerate symmetric invariant bilinear form $(\cdot | \cdot)$ extended to $\mathfrak{g}((z^{-1}))$ by bilinearity. Let f be a non-zero nilpotent element of \mathfrak{g} , and let (f_1, f_2, E) be an integrable triple for f and consider the decomposition (4.3) of $\mathfrak{g}((z^{-1}))$. Let $U(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0}$ and $h(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>-1}$ be a solution of equation (4.10). Let $\mathfrak{c}(\mathfrak{h})$ be the center of \mathfrak{h} , and let

$$\mathfrak{a} = \left\{ a \in \mathfrak{c}(\mathfrak{h}) \,\middle|\, [a, f_2] = 0 \right\} \subset \mathfrak{g}((z^{-1})). \tag{4.13}$$

For $a \in \mathfrak{a}$, let

$$\int g_a(z) = \int (1 \otimes a | h(z)) = \sum_{n \in \mathbb{Z}_{>0}} \int g_{a,n} z^{N-n}, \tag{4.14}$$

where N is the largest power of z appearing in $\int (1 \otimes a | h(z))$ with non-zero coefficient, and $(\cdot | \cdot)$ is defined in (4.8). Let $W = W(\mathfrak{g}, f, E) \subset V(\mathfrak{p})$ be the classical affine W-algebra with its compatible PVA structures $\{\cdot_{\lambda} \cdot\}_0$ and $\{\cdot_{\lambda} \cdot\}_{\infty}$ introduced in Definition 4.2. Then

$$\mathcal{A} = \operatorname{Span} \{ \int g_{a,n} \mid a \in \mathfrak{a}, n \in \mathbb{Z}_{>0} \}$$

is an infinite-dimensional abelian subalgebra of $W/\partial W \subset V(\mathfrak{p})/\partial V(\mathfrak{p})$ (with respect to both 0 and ∞ -Lie brackets), defining a hierarchy of bi-Hamiltonian equations

$$\frac{dw}{dt_{a,n}} = \{g_{a,n_{\lambda}}w\}_0\big|_{\lambda=0} = \{g_{a,n+1_{\lambda}}w\}_\infty\big|_{\lambda=0}, \quad w \in \mathcal{W}, a \in \mathfrak{a}, n \in \mathbb{Z}_{\geq 0}. \tag{4.15}$$

Moreover, if $a \in \mathfrak{a}$ is not central in $\mathfrak{g}((z^{-1}))$, then $\dim \operatorname{Span}\{\int g_{a,n}\}_{n \in \mathbb{Z}_{\geq 0}} = \infty$. Consequently, the hierarchy (4.15) is integrable.

Proof. First, we show that the results in [DSKV13, Sect. 4.5] remain valid if $\mathfrak{h} = \text{Ker ad}(f_1 + zE)$ and $a \in \mathfrak{a}$. Indeed, by the definition (4.13) of \mathfrak{a} , $[h(z), 1 \otimes a] = 0$, $[f_1 + zE, a] = 0$, $[f_2, a] = 0$, so that $[\partial + 1 \otimes (f + zE) + h(z), 1 \otimes a] = 0$. Applying the automorphism $e^{-\operatorname{ad} U(z)}$ and using equation (4.10), we thus get

$$[\partial + 1 \otimes (f + zE) + q, e^{-\operatorname{ad} U(z)}(1 \otimes a)] = 0,$$

which is the analogue of [DSKV13, Cor. 4.6]. The proofs of [DSKV13, Lem. 4.7, Lem. 4.8, Thm. 4.9] go through unchanged in the present setting. As a result, we have that the local functionals $\int g_{a,n} \in \mathcal{V}(\mathfrak{p})/\partial \mathcal{V}(\mathfrak{p})$, $n \in \mathbb{Z}_{\geq 0}$, defined by equation (4.14), solve the Lenard–Magri recurrence relation (3.5) and $\{\int g_{a,0}, \int w\}_{\infty} = 0$, for every $w \in \mathcal{W}$. Hence, by Lemma 3.2, all $\int g_{a,n}$ commute in the Lie algebra $\mathcal{W}/\partial \mathcal{W}$.

Next, by Proposition 4.5 and the proofs of [DSKV13, Lem. 4.10, Prop. 4.11], it follows that the Laurent series $\int g_a(z)$ defined by (4.14) is independent of the choice of $U(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{g}((z^{-1}))_{>0}$, $h(z) \in \mathcal{V}(\mathfrak{p}) \otimes \mathfrak{h}_{>-1}$, solving equation (4.10), and that its coefficients $\int g_{a,n}$, $n \in \mathbb{Z}_{\geq 0}$, lie in $\mathcal{W}/\partial\mathcal{W}$. Hence, by Lemma 3.2, we have that \mathcal{A} is an abelian subalgebra, and, by the discussion in Section 3.2, we get the hierarchy of bi-Hamiltonian equations (4.15).

Finally, we are left to show that if a is not in the center of $\mathfrak{g}((z^{-1}))$, then the local functionals $\int g_{a,n}$, $n \in \mathbb{Z}_{\geq 0}$, span an infinite-dimensional subspace. This follows verbatim from the results in [DSKV13, Sect. 4.7] noticing that $\mathrm{ad}(f+zE)$ and $\mathrm{ad}\,a$ commute. Since $f_1 + zE \in \mathfrak{a}$, we see that \mathcal{A} is infinite-dimensional, as claimed. \square

Remark 4.7. Theorem 4.6 holds in the more general setting of an arbitrary finitedimensional Lie algebra \mathfrak{g} with a non-degenerate symmetric invariant bilinear form $(\cdot|\cdot)$, an \mathfrak{sl}_2 -triple $\{e,h,f\}$ in \mathfrak{g} , and an integrable triple (f_1,f_2,E) for f.

References

- [Adl79] M. Adler, On a trace functional for formal pseudodifferential operators and the symplectic structure of the Korteweg-de Vries equation, Invent. Math. 50 (1979), 219–248.
- [BDSK09] A. Barakat, A. De Sole, V. G. Kac, Poisson vertex algebras in the theory of Hamiltonian equations, Japan. J. Math. 4 (2009), n. 2, 141–252.
- [CMG93] D. Collingwood, W. McGovern, Nilpotent Orbits in Semisimple Lie Algebra, Van Nostrand Reinhold Mathematics Series, Van Nostrand Reinhold Co., New York, 1993.
- $[\mathrm{DK85}]$ J. Dadok, V. G. Kac, Polar representations, J. Algebra **92** (1985), no. 2, 504–524.
- [DSJKV20] A. De Sole, M. Jibladze, V. G Kac, D. Valeri, Integrable triples in semisimple Lie algebras, arXiv: 2012.12913 (2020).
- [DSK06] A. De Sole, V. G. Kac, Finite vs. affine W-algebras, Japan. J. Math. 1 (2006), no.1, 137–261.
- [DSKV13] A. De Sole, V. G. Kac, D. Valeri, Classical W-algebras and generalized Drinfeld-Sokolov bi-Hamiltonian systems within the theory of Poisson vertex algebras, Comm. Math.Phys. **323** (2013), n. 2, 663–711.
- [DSKV14] A. De Sole, V. G. Kac, D. Valeri, Classical W-algebras and generalized Drinfeld-Sokolov hierarchies for minimal and short nilpotents, Comm. Math. Phys. 331 (2014), n. 2, 623-676. Erratum in Commun. Math. Phys. 333 (2015), n. 3, 1617-1619.
- [DS85] V. Drinfeld, V. Sokolov, Lie algebras and equations of KdV type, Soviet J. Math. 30 (1985), 1975–2036.
- [EKV13] A. G. Elashvili, V. G. Kac, E. B. Vinberg, Cyclic elements in semisimple Lie algebras, Transform. Groups 18 (2013), 97–130.
- [FORTW92] L. Fehér, L. O'Raifeartaigh, P. Ruelle, I. Tsutsui, A. Wipf, On Hamiltonian reductions of the Wess-Zumino-Novikov-Witten theories, Phys. Rep. 222 (1992), no.1, 1–64.
- [FGMS95] C. Fernández-Pousa, M. Gallas, L. Miramontes, J. Sánchez Guillén, W-algebras from soliton equations and Heisenberg subalgebras, Ann. Physics 243 (1995), no. 2, 372–419.
- [GGKM67] C. S. Gardner, J. M. Greene, Kruskal, R. M. Miura, Method for solving the Korteweg-de Vries equation, Phys. Rev. Lett. 19, no. 19, 1095–1097.
- [GD76] И. М. Гельфанд, Л. А. Дикий, Дробные степени операторов и гамильтоновы системы, Функц. анализ и его прил. 10 (1976), вып. 4, 13–29. Engl. transl.: I. M. Gelfand, L. A. Dickey, Fractional powers of operators and Hamiltonian systems, Funct. Anal. Appl. 10 (1976), 259–273.
- [GD78] I. M. Gelfand, L. A. Dickey, Family of Hamiltonian structures connected with integrable non-linear equations, preprint, IPM, Moscow (in Russian),

500

1978. English version in: Collected Papers of I. M. Gelfand, Vol. 1, Springer-Verlag, Berlin, 1987, pp. 625–646.

- [Lax68] P. D. Lax, Integrals of non-linear equations of evolution and solitary waves, Comm. Pure and Appl. Math., 21 (1968), no. 5, 467–490.
- [Mag78] F. Magri, A simple model of the integrable Hamiltonian equation, J. Math. Phys. 19 (1978), no. 5, 1156–1162.
- [ZF71] В. Е. Захаров, Л. Д. Фаддеев, Уравнение Кортевега-де Фриза—вполне интегрируемая гамильтонова система, Функц. анализ и его прил. 5 (1971), вып. 4, 18–27. Engl. transl.: V. E. Zakharov, L. D. Faddeev, Korteweg-de Vries equation: A completely integrable Hamiltonian system, Funct. Anal. Appl. 5 (1971), 280–287.

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