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# ON THE GLOBAL AND LOCAL SOLUTION OF THE MULTIDIMENSIONAL DARBOUX PROBLEM FOR SOME NONLINEAR WAVE EQUATIONS 

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Dedicated to the memory of Prof. Gaetano Fichera


#### Abstract

We consider a multidimensional analogue of the Darboux problem for wave equations with power nonlinearity. Depending on the spatial dimension of an equation, a power nonlinearity exponent and the sign in front of a nonlinear term, it is proved that the Darboux problem is globally solvable in some cases, but has no global solution in other cases though the local solvability of this problem remains in force.


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## 1. Statement of the Problem

Let us consider a nonlinear wave equation of the form

$$
\begin{equation*}
L u:=\frac{\partial^{2} u}{\partial t^{2}}-\Delta u+m u=f(u)+F, \tag{1}
\end{equation*}
$$

where $f$ and $F$ are given real functions, $f$ being nonlinear, $f(0)=0$, and $u$ is an unknown real function, $m=$ const $\geq 0, \Delta=\sum_{i=1}^{n} \frac{\partial^{2}}{\partial x_{i}^{2}}, n \geq 2$.

Denote by $D: t>|x|, x_{n}>0$, the half of the light cone of the future which is bounded by the part $S^{0}=D \cap\left\{x_{n}=0\right\}$ of the hyperplane $x_{n}=0$ and by the half $S: t=|x|, x_{n} \geq 0$, of the characteristic conoid $C: t=|x|$ of equation (1). Assume $D_{T}=\{(x, t) \in D: t<T\}, S_{T}^{0}=\left\{(x, t) \in S^{0}: t \leq T\right\}$, $S_{T}=\{(x, t) \in S: t \leq T\}, T>0$. When $T=\infty$, it is obvious that $D_{\infty}=D$, $S_{\infty}^{0}=S^{0}$ and $S_{\infty}=S$.

We will consider the problem on defining, in the domain $D_{T}$, a solution $u(x, t)$ of equation (1) by the boundary conditions

$$
\begin{equation*}
\left.u\right|_{S_{T}^{0}}=0,\left.\quad u\right|_{S_{T}}=g \tag{2}
\end{equation*}
$$

where $g$ is a given real function on $S_{T}$.
Problem (1),(2) is a multidimensional variant of the first Darboux problem for the nonlinear equation (1) when one part of the problem data support is a characteristic manifold, while the remaining part is a manifold of time type [1, Ch. III, § $\left.1.1^{0}\right]$.

Problems pertaining to the existence or nonexistence of a global solution of the Cauchy problem for nonlinear equations of form (1) with the boundary conditions $\left.u\right|_{t=0}=u_{0},\left.\frac{\partial u}{\partial t}\right|_{t=0}=u_{1}$ are considered in [2]-[17]. As for multidimensional variants of the first Darboux problem for linear hyperbolic equations of second order, they are well posed and their global solvability is proved in the corresponding function spaces [18]-[20].

In this paper, we discuss the concrete cases for the nonlinear function $f=$ $f(u)$, where problem (1),(2) is globally solvable in some cases, but has no global solution in other cases though the local solvability of this problem remain in force.

## 2. Global Solvability of Problem (1),(2) in the Case of Linearity of the Form $f(u)=-\lambda|u|^{p} u$

For $f(u)=-\lambda|u|^{p} u$, where $\lambda \neq 0$ and $p>0$ are given real numbers, equation (1) takes the form

$$
\begin{equation*}
L u:=\frac{\partial^{2} u}{\partial t^{2}}-\Delta u+m u=-\lambda|u|^{p} u+F . \tag{3}
\end{equation*}
$$

Note that equation (3) arises in relativistic quantum mechanics [21]-[24].
In this section, our consideration is limited to the case where the boundary conditions (2) are assumed homogeneous, i.e.

$$
\begin{equation*}
\left.u\right|_{S_{T}^{0}}=0,\left.\quad u\right|_{S_{T}}=0 . \tag{4}
\end{equation*}
$$

We assume that $\stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)=\left\{u \in W_{2}^{1}\left(D_{T}\right):\left.u\right|_{S_{T}^{0} \cup S_{T}}=0\right\}$, where $W_{2}^{1}\left(D_{T}\right)$ is the known Sobolev space with the norm

$$
\|\left. u\right|_{W_{2}^{1}\left(D_{T}\right)} ^{2}=\int_{D}\left[u^{2}+\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}\right] d x d t
$$

and the boundary condition $\left.u\right|_{S_{T}^{0} \cup S_{T}}=0$ should be understood in terms of trace theory [25, Ch. I, $\S \S 5,6]$.

Remark 1. The embedding operator $I: \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow L_{q}\left(D_{T}\right)$ is a linear continuous compact operator for $1<q<\frac{2(n+1)}{n-1}$, where $n>1$ [25, Ch.I, § 7]. Simultaneously, the Nemitski operator $K: L_{q}\left(D_{T}\right) \rightarrow L_{2}\left(D_{T}\right)$ acting by the formula $K u:=-\lambda|u|^{p} u$ is continuous and bounded if $q \geq 2(p+1)[26$, Ch. V, § 17.5], [27, Ch. III, $\S \S 12.10 ; 12.11]$. Therefore if $p<\frac{2}{n-1}$, i.e. $2(p+1)<$ $\frac{2(n+1)}{n-1}$, then there exists a number $q$ such that $1<2(p+1) \leq q<\frac{2(n+1)}{n-1}$ and hence the operator

$$
\begin{equation*}
K_{0}=K I: \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow L_{2}\left(D_{T}\right) \tag{5}
\end{equation*}
$$

is continuous and compact. Moreover, from $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ it follows that $u \in L_{p+1}\left(D_{T}\right)$. As has been mentioned above, it is assumed that here and everywhere $p>0$.

If $u \in C^{2}\left(\bar{D}_{T}\right)$ is a classical solution of problem (3),(4), then, after multiplying both parts of equation (3) by an arbitrary function $\varphi \in C^{2}\left(\bar{D}_{T}\right)$ that satisfies the condition $\left.\varphi\right|_{t=T}=0$ and applying integration by parts, we obtain

$$
\begin{align*}
\int_{S_{T}^{0} \cup S_{T}} \frac{\partial u}{\partial N} \varphi d s & -\int_{D_{T}} u_{t} \varphi_{t} d x d t+\int_{D_{T}} \nabla_{x} u \nabla_{x} \varphi d x d t+\int_{D_{T}} m u \varphi d x d t \\
& =-\lambda \int_{D_{T}}|u|^{p} u \varphi d x d t+\int_{D_{T}} F \varphi d x d t \tag{6}
\end{align*}
$$

where $\frac{\partial}{\partial N}=\nu_{0} \frac{\partial}{\partial t}-\sum_{i=1}^{n} \nu_{i} \frac{\partial}{\partial x_{i}}$ is a derivative with respect to the conormal, $\nu=$ $\left(\nu_{1}, \ldots, \nu_{n}, \nu_{0}\right)$ is the unit outward normal to $\partial D_{T}, \nabla_{x}=\left(\frac{\partial}{\partial x_{1}}, \ldots, \frac{\partial}{\partial x_{n}}\right)$. Since the hypersurface $S_{T}$ is the characteristic manifold on which the operator $\frac{\partial}{\partial N}$ is an internal differential operator, by (4) we have $\left.\frac{\partial}{\partial N}\right|_{S_{T}}=0$. Therefore, assuming additionally that the function $\left.\varphi\right|_{S_{T}^{0}}=0$, from equality (6) we obtain

$$
\begin{gather*}
-\int_{D_{T}} u_{t} \varphi_{t} d x d t+\int_{D_{T}} \nabla_{x} u \nabla_{x} \varphi d x d t+\int_{D_{T}} m u \varphi d x d t \\
=-\lambda \int_{D_{T}}|u|^{p} u \varphi d x d t+\int_{D_{T}} F \varphi d x d t \tag{7}
\end{gather*}
$$

Since, by virtue of Remark 1 , from $u \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ it follows that $|u|^{p} u \in$ $L_{2}\left(D_{T}\right)$, equality (7) can underlie the definition of a weak generalized solution of problem (3),(4) of the class $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$.

Definition 1. Let $F \in L_{2}\left(D_{T}\right)$ and $0<p<\frac{2}{n-1}$. A function $u \in$ $\stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is called a weak generalized solution of the nonlinear problem (3),(4) in the domain $D_{T}$ if the integral equality (7) is fulfilled for any function $\varphi \in W_{2}^{1}\left(D_{T}\right)$ such that $\left.\varphi\right|_{t=T}=0,\left.\varphi\right|_{S_{T}^{0}}=0$.

Assume that

$$
\stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup S_{T}\right)=\left\{u \in C^{2}\left(\bar{D}_{T}\right):\left.u\right|_{S_{T}^{0} \cup S_{T}}=0\right\} .
$$

Definition 2. Let $F \in L_{2}\left(D_{T}\right)$ and $0<p<\frac{2}{n-1}$. A function $u \in$ $\stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is called a strong generalized solution of the nonlinear problem (3),(4) in the domain $D_{T}$ if there exists a sequence of functions $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup\right.$ $S_{T}$ ) such that $u_{k} \rightarrow u$ in the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ and $\left[L u_{k}+\lambda\left|u_{k}\right|^{p} u_{k}\right] \rightarrow F$ in the space $L_{2}\left(D_{T}\right)$. Moreover, from Remark 1 it follows the sequence $\left\{\lambda\left|u_{k}\right|^{p} u_{k}\right\}$ converges to the function $\lambda|u|^{p} u$ in the space $L_{2}\left(D_{T}\right)$ as $u_{k} \rightarrow u$ in the space $\stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$.

Remark 2. One can easily verify that if $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is a strong generalized solution of problem (3),(4), then it is automatically a weak generalized solution of this problem.

Definition 3. Let $0<p<\frac{2}{n-1}, F \in L_{2, l o c}(D)$ and $F \in L_{2}\left(D_{T}\right)$ for any $T>0$. We say that problem (3),(4) is globally solvable if for any $T>0$ this problem has a strong generalized solution in the domain $D_{T}$ from the space $\stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$.

Lemma 1. Let $\lambda \geq 0,0<p<\frac{2}{n-1}$ and $F \in L_{2}\left(D_{T}\right)$. Then for any strong generalized solution $u \in \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ of problem (3),(4) in the domain $D_{T}$ the following a priori estimate is valid:

$$
\begin{equation*}
\|u\|_{\dot{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \leq \sqrt{\frac{e}{2}} T\|F\|_{L_{2}\left(D_{T}\right)} . \tag{8}
\end{equation*}
$$

Proof. Let $u \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ be a strong generalized solution of problem (3),(4). By virtue of Definition 2 there exists a sequence of functions $u_{k} \in$ $\stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup S_{T}\right)$ such that

$$
\begin{equation*}
\lim _{k \rightarrow \infty}\left\|u_{k}-u\right\|_{\dot{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}=0, \quad \lim _{k \rightarrow \infty}\left\|L u_{k}+\lambda\left|u_{k}\right|^{p} u+k-F\right\|_{L_{2}\left(D_{T}\right)}=0 \tag{9}
\end{equation*}
$$

Let us consider the function $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup S_{T}\right)$ as a solution of the problem

$$
\begin{gather*}
L u_{k}+\lambda\left|u_{k}\right|^{p} u_{k}=F_{k},  \tag{10}\\
\left.u_{k}\right|_{S_{T}^{0}}=0,\left.\quad u\right|_{S_{T}}=0 . \tag{11}
\end{gather*}
$$

Here

$$
\begin{equation*}
F_{k}=L u_{k}+\lambda\left|u_{k}\right|^{p} u_{k} \tag{12}
\end{equation*}
$$

Multiplying both parts of equation (10) by $\frac{\partial u_{k}}{\partial t}$ and integrating over the domain $D_{\tau}, 0<\tau \leq T$, we obtain

$$
\begin{gather*}
\frac{1}{2} \int_{D_{\tau}} \frac{\partial}{\partial t}\left(\frac{\partial u_{k}}{\partial t}\right)^{2} d x d t-\int_{D_{\tau}} \Delta u_{k} \frac{\partial u_{k}}{\partial t} d x d t+\frac{m}{2} \int_{D_{\tau}} \frac{\partial}{\partial t} u_{k}^{2} d x d t \\
\quad+\frac{\lambda}{p+2} \int_{D_{\tau}} \frac{\partial}{\partial t}\left|u_{k}\right|^{p+2} d x d t=\int_{D_{\tau}} F_{k} \frac{\partial u_{k}}{\partial t} d x d t \tag{13}
\end{gather*}
$$

Assume that $\Omega_{\tau}:=D_{T} \cap\{t=\tau\}, 0<\tau<T$. It is obvious that $\partial D_{\tau}=$ $S_{\tau}^{0} \cup S_{\tau} \cup \Omega_{\tau}$. Using (11), the equality $\left.\frac{\partial u}{\partial t}\right|_{S_{T}^{0}}=0$ and also the equalities $\left.\nu\right|_{\Omega_{\tau}}=$ $(0, \ldots, 0,1), \nu_{S_{T}^{0}}=(0, \ldots, 0,-1,0)$ and performing the integration by parts, we obtain

$$
\int_{D_{\tau}} \frac{\partial}{\partial t}\left(\frac{\partial u_{k}}{\partial t}\right)^{2} d x d t=\int_{\partial D_{\tau}}\left(\frac{\partial u_{k}}{\partial t}\right)^{2} \nu_{0} d s=\int_{\Omega_{\tau}}\left(\frac{\partial u_{k}}{\partial t}\right)^{2} d x+\int_{S_{\tau}}\left(\frac{\partial u_{k}}{\partial t}\right)^{2} \nu_{0} d s
$$

$$
\begin{aligned}
& \int_{D_{\tau}} \frac{\partial}{\partial t} u_{k}^{2} d x d t=\int_{\partial D_{\tau}} u_{k}^{2} \nu_{0} d s=\int_{\Omega_{\tau}} u_{k}^{2} d x \\
& \int_{D_{\tau}} \frac{\partial}{\partial t}\left|u_{k}\right|^{p+2} d x d t=\int_{\partial D_{\tau}}\left|u_{k}\right|^{p+2} \nu_{0} d s=\int_{\Omega_{\tau}}\left|u_{k}\right|^{p+2} d x \\
& \int_{D_{\tau}} \frac{\partial^{2} u_{k}}{\partial x_{i}^{2}} \frac{\partial u_{k}}{\partial t} d x d t=\int_{\partial D_{\tau}} \frac{\partial u_{k}}{\partial x_{i}} \frac{\partial u_{k}}{\partial t} \nu_{i} d s-\frac{1}{2} \int_{D_{\tau}} \frac{\partial}{\partial t}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2} d x d t \\
& \quad=\int_{\partial D_{\tau}} \frac{\partial u_{k}}{\partial x_{i}} \frac{\partial u_{k}}{\partial t} \nu_{i} d s-\frac{1}{2} \int_{D_{\tau}}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2} \nu_{0} d s \\
& \quad=\int_{S_{\tau}} \frac{\partial u_{k}}{\partial x_{i}} \frac{\partial u_{k}}{\partial t} \nu_{i} d s-\frac{1}{2} \int_{S_{\tau}}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2} \nu_{0} d s-\frac{1}{2} \int_{\Omega_{\tau}}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2} d x
\end{aligned}
$$

Hence by virtue of (13) it follows that

$$
\begin{align*}
\int_{D_{\tau}} F_{k} & \frac{\partial u_{k}}{\partial t} d x d t \\
& =\int_{S_{\tau}} \frac{1}{2 \nu_{0}}\left[\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}} \nu_{0}-\frac{\partial u_{k}}{\partial t} \nu_{i}\right)^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}\left(\nu_{0}^{2}-\sum_{j=1}^{n} \nu_{j}^{2}\right)\right] d s \\
& +\frac{1}{2} \int_{\Omega_{\tau}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x+\frac{\lambda}{p+2} \int_{\Omega_{\tau}}\left|u_{k}\right|^{p+2} d x . \tag{14}
\end{align*}
$$

Since $S_{\tau}$ is a characteristic manifold, we have

$$
\begin{equation*}
\left.\left(\nu_{0}^{2}-\sum_{j=1}^{n} \nu_{j}^{2}\right)\right|_{S_{\tau}}=0 \tag{15}
\end{equation*}
$$

Taking into account that the operator $\left(\nu_{0} \frac{\partial}{\partial x_{i}}-\nu_{i} \frac{\partial}{\partial t}\right), 1 \leq i \leq n$, is an internal differential operator on $S_{\tau}$, by virtue of (11) we have

$$
\begin{equation*}
\left.\left(\frac{\partial u_{k}}{\partial x_{i}} \nu_{0}-\frac{\partial u_{k}}{\partial t} \nu_{i}\right)\right|_{S_{\tau}}=0, \quad i=1, \ldots, n \tag{16}
\end{equation*}
$$

With regard for (15), (16), from (14) we have

$$
\int_{\Omega_{\tau}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x+\frac{2 \lambda}{p+2} \int_{\Omega_{\tau}}\left|u_{k}\right|^{p+2} d x=2 \int_{D_{\tau}} F_{k} \frac{\partial u_{k}}{\partial t} d x d t
$$

which, by virtue of $\lambda \geq 0$, implies in turn that

$$
\begin{equation*}
\int_{\Omega_{\tau}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x \leq 2 \int_{D_{\tau}} F_{k} \frac{\partial u_{k}}{\partial t} d x d t \tag{17}
\end{equation*}
$$

Using the notation

$$
w(\delta)=\int_{\Omega_{\delta}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x
$$

and taking into account that the inequality $2 F_{k} \frac{\partial u_{k}}{\partial t} \leq \varepsilon\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\frac{1}{\varepsilon} F_{k}^{2}$ holds for any $\varepsilon=$ const $>0$, from (17) we obtain

$$
\begin{equation*}
w(\delta) \leq \varepsilon \int_{0}^{\delta} w(\sigma) d \sigma+\frac{1}{\varepsilon}\left\|F_{k}\right\|_{L_{2}\left(D_{\delta}\right)}^{2}, \quad 0<\delta \leq T \tag{18}
\end{equation*}
$$

If we take into account that the value $\left\|F_{k}\right\|_{L_{2}\left(D_{\delta}\right)}^{2}$ as a function of $\delta$ is nondecreasing, then, by virtue Gronwall's lemma [28, Ch. I, § 2], from (18) it follows that

$$
\|w(\delta)\| \leq \frac{1}{\varepsilon}\left\|F_{k}\right\|_{L_{2}\left(D_{\delta}\right)}^{2} \exp \delta \varepsilon .
$$

Hence, since $\inf _{\varepsilon>0} \frac{\exp \delta \varepsilon}{\varepsilon}=e \delta$ for $\varepsilon=\frac{1}{\delta}$, we obtain

$$
\begin{equation*}
w(\delta) \leq e \delta\left\|F_{k}\right\|_{L_{2}\left(D_{\delta}\right)}^{2}, \quad 0<\delta \leq T \tag{19}
\end{equation*}
$$

which implies in turn that

$$
\begin{align*}
\left\|u_{k}\right\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}^{2} & =\int_{D_{T}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x d t \\
& =\int_{0}^{T} w(\delta) d \delta \leq \frac{e}{2} T^{2}\left\|F_{k}\right\|_{L_{2}\left(D_{T}\right)}^{2} \tag{20}
\end{align*}
$$

Here we have used the fact that in the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ one of the equivalent norms is given by the expression

$$
\left\{\int_{D_{T}}\left[m u_{k}^{2}+\left(\frac{\partial u_{k}}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u_{k}}{\partial x_{i}}\right)^{2}\right] d x d t\right\}^{\frac{1}{2}}
$$

independently of the assumption whether $m=0$ or $m>0$. Indeed, by a standard reasoning, the equalities $\left.u\right|_{S_{T}}=0$ and $u(x, t)=\int_{\psi(x)}^{t} \frac{\partial u(x, \tau)}{\partial t} d \tau,(x, t) \in$ $\bar{D}_{T}$, where $t-\psi(x)=0$ is an equation of the conical manifold $S_{T}$, imply the inequality [25, Ch. I, § 6]

$$
\int_{D_{T}} u^{2}(x, t) d x d t \leq T^{2} \int_{D_{T}}\left(\frac{\partial u}{\partial t}\right)^{2} d x d t
$$

Now, passing to the limit in inequality (20) as $k \rightarrow \infty$, we obtain (8), which proves the lemma.

Theorem 1. Let $\lambda>0,0<p<\frac{2}{n-1}, F \in L_{2, l o c}(D)$ and $F \in L_{2}\left(D_{T}\right)$ for any $T>0$. Then problem (3), (4) is globally solvable, i.e. for any $T>0$ this problem has a strong generalized solution $u \in \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ in the domain $D_{T}$.

Proof. Before proceeding to the discussion whether the nonlinear problem (3), (4) is solvable, we will consider the solvability for the linear case where it is assumed that in equation (3) the parameter $\lambda=0$, i.e. for the problem

$$
\left\{\begin{array}{l}
L u(x, t)=F(x, t), \quad(x, t) \in D_{T},  \tag{21}\\
\left.u\right|_{S_{T}^{0}}=0,\left.\quad u\right|_{S_{T}}=0
\end{array}\right.
$$

In that case, analogously to the above, we introduce, for $F \in L_{2}\left(D_{T}\right)$, the notion of a strong generalized solution $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ of problem (21) for which there exists a sequence of functions $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup S_{T}\right)$ such that $\lim _{k \rightarrow \infty}\left\|u_{k}-u\right\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}=0, \lim _{k \rightarrow \infty}\left\|L u_{k}-F\right\|_{L_{2}\left(D_{T}\right)}=0$. It should be noted here that by virtue of Lemma 1 for $\lambda=0$ the a priori estimate (8) holds for a strong generalized solution of problem (21), too.

Since the space $C_{0}^{\infty}\left(D_{T}\right)$ of finite, infinitely differentiable in $D_{T}$, functions is dense in $L_{2}\left(D_{T}\right)$, for given $F \in L_{2}\left(D_{T}\right)$ there exists a sequence of functions $F_{k} \in C_{0}^{\infty}\left(D_{T}\right)$ such that $\lim _{k \rightarrow \infty}\left\|F_{k}-F\right\|_{L_{2}\left(D_{T}\right)}=0$. If we continue the function $F_{k}$ in an odd manner with respect to the variable $x_{n}$ into the domain $D_{T}^{-}:=$ $\left\{(x, t) \in R^{n+1}: x_{n}<0,|x|<t<T\right\}$ and after that continue the resulting function by zero beyond the domain $D_{T} \cup D_{T}^{-}$and denote it by the previous symbol, then for fixed $k$ we obtain $F_{k} \in C^{\infty}\left(R_{+}^{n+1}\right)$ for which the support is $\operatorname{supp} F_{k} \subset D_{\infty} \cup D_{\infty}^{-}$, where $R_{+}^{n+1}=R^{n+1} \cap\{t \geq 0\}$. Denote by $u_{k}$ a solution of the Cauchy problem

$$
\begin{equation*}
L u_{k}=F_{k},\left.\quad u_{k}\right|_{t=0}=0,\left.\quad \frac{\partial u_{k}}{\partial t}\right|_{t=0}=0 \tag{22}
\end{equation*}
$$

which, as is known, exists, is unique and belongs to the space $C^{\infty}\left(R_{+}^{n+1}\right)[29$, Ch. V, §6]. Moreover, since supp $F_{k} \subset D_{\infty} \cup D_{\infty}^{-} \subset\left\{(x, t) \in R^{n+1}: t>|x|\right\}$ and $\left.u_{k}\right|_{t=0}=0,\left.\frac{\partial u_{k}}{\partial t}\right|_{t=0}=0$, by taking into account the geometry of the solution dependence domain of the linear wave equation $L u=F$ we have $\operatorname{supp} u_{k} \subset$ $\left\{(x, t) \in R^{n+1}: t>|x|\right\}$ and in particular $\left.u_{k}\right|_{S_{T}}=0$. On the other hand, the function $\widetilde{u}_{k}\left(x_{1}, \ldots, x_{n}, t\right)=-u_{k}\left(x_{1}, \ldots,-x_{n}, t\right)$ is also a solution of the Cauchy problem (22), since the function $F_{k}$ is odd with respect to the variable $x_{n}$. Hence by virtue of the uniqueness of a solution of the Cauchy problem we have $\widetilde{u}_{k}=u_{k}$, i.e. $u_{k}\left(x_{1}, \ldots,-x_{n}, t\right)=-u_{k}\left(x_{1}, \ldots, x_{n}, t\right)$ and thereby the function $u_{k}$ is also even with respect to the variable $x_{n}$. This in turn implies that $\left.u_{k}\right|_{x_{n}=0}=0$, which together with the condition $\left.u_{k}\right|_{S_{T}}=0$ implies that by preserving the previous notation for the restriction of the function $u_{k}$ to the domain $D_{T}$ we obtain $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0} \cup S_{T}\right)$. Furthermore, by virtue of (8) and
(22) the inequality

$$
\begin{equation*}
\left\|u_{k}-u_{l}\right\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \leq \sqrt{\frac{e}{2}} T\left\|F_{k}-F\right\|_{L_{2}\left(D_{T}\right)} \tag{23}
\end{equation*}
$$

is valid because the a priori estimate (8) holds for a strong generalized solution of the linear problem (21), too.

Since the sequence $\left\{F_{k}\right\}$ is fundamental in $L_{2}\left(D_{T}\right)$, by virtue of (23) the sequence $\left\{u_{k}\right\}$, too, is fundamental in the total space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$. Thus there exists a function $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ such that $\lim _{k \rightarrow \infty}\left\|u_{k}-u\right\|_{\dot{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}=0$ and since $L u_{k}=F_{k} \rightarrow F$ in the space $L_{2}\left(D_{T}\right)$, this function is, by definition, a strong generalized solution of problem (21). The uniqueness of this solution from the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ follows from the a priori estimate (8). Therefore for this solution $u$ of problem (21) we can write $u=L^{-1} F$, where $L^{-1}: L_{2}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is a linear continuous operator whose norm admits, by virtue of (8), an estimate

$$
\begin{equation*}
\left\|L^{-1}\right\|_{L_{2}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}} \leq \sqrt{\frac{e}{2}} T . \tag{24}
\end{equation*}
$$

Note that, by virtue of (24) and Remark 1, Definition 2 and Remark 2, for $F \in L_{2}\left(D_{T}\right), 0<p<\frac{2}{n-1}$, a function $u \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is a strong generalized solution of problem (3),(4) if and only if $u$ is a solution of the functional equation

$$
\begin{equation*}
u=L^{-1}\left(-\lambda|u|^{p} u+F\right) \tag{25}
\end{equation*}
$$

in the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$.
Rewrite equation (25) as follows:

$$
\begin{equation*}
u=A u:=L^{-1}\left(K_{0} u+F\right) \tag{26}
\end{equation*}
$$

where the operator $K_{0}: \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow L_{2}\left(D_{T}\right)$ from (5) is continuous and compact according to Remark 1. Therefore the operator $A: \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow$ $\stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is also continuous and compact by virtue of (24). At the same time, according to Lemma 1 , for any parameter $\mu \in[0,1]$ and for any solution of an equation with parameter $u=\mu A u$ we have the a priori estimate $\|u\|_{W_{2}^{1}\left(D_{T}, S_{T}\right)} \leq c\|F\|_{L_{2}\left(D_{T}\right)}$, where the positive constant $c$ does not depend on $u$, $\mu$ and $F$. Thus, by the Lere-Shauder theorem [30, Ch. VIII, § 35.5], equation (26) and therefore problem (3),(4), too, have at least one solution $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$. Theorem 1 is proved.
3. Nonexistence of Global Solvability of Problem (1),(2) in the Case of Nonlinearity of the Form $f(u)=\lambda|u|^{p+1}$
Below we will consider the case where the coefficients in problem (1),(2) are $m=0$ and $f(u)=\lambda|u|^{p+1}$ with $\lambda$ and $p$ being given positive numbers, i.e. the problem

$$
\begin{gather*}
\square u:=\frac{\partial^{2} u}{\partial t^{2}}-\Delta u=\lambda|u|^{p+1}+F  \tag{27}\\
\left.u\right|_{S_{T}^{0}}=0,\left.\quad u\right|_{S_{T}}=g \tag{28}
\end{gather*}
$$

in the domain $D_{T}, T>0$, where $g$ is a given real function on $S_{T}$ that by virtue of (28) satisfies the compatibility condition $\left.g\right|_{\partial S_{T}}=0$.

Remark 3. Assuming that $F \in L_{2}\left(D_{T}\right), g \in W_{2}^{1}\left(S_{T}\right)$ and $0<p<\frac{2}{n-1}$, analogously to Definitions 1 and 2 with regard to a weak and a strong generalized solution of problem (3),(4) in the domain $S_{T}$ and taking into account Remark 1, we introduce the notions of a weak and a strong generalized solution of problem (27), (28):
(i) a function $u \in W_{2}^{1}\left(D_{T}\right)$ is called a weak generalized solution of the nonlinear problem (27),(28) in the domain $D_{T}$ if for any function $\varphi \in W_{2}^{1}\left(D_{T}\right)$ such that $\left.\varphi\right|_{t=T}=0,\left.\varphi\right|_{S_{T}^{0}}=0$ the following integral equality is valid:

$$
\begin{align*}
-\int_{D_{T}} u_{t} \varphi_{t} d x d t+\int_{D_{T}} & \nabla_{x} u \nabla_{x} \varphi d x d t \\
& =\lambda \int_{D_{T}}|u|^{p+1} \varphi d x d t+\int_{D_{T}} F \varphi d x d t-\int_{S_{T}} \frac{\partial g}{\partial N} \varphi d s \tag{29}
\end{align*}
$$

where $\frac{\partial}{\partial N}=\nu_{0} \frac{\partial}{\partial t}-\sum_{i=1}^{n} \nu_{i} \frac{\partial}{\partial x_{i}}$ is an internal derivative with respect to the conormal which is an internal differential operator on $S_{T}$, since the conical manifold $S_{T}$ is characteristic, $\nu=\left(\nu_{1}, \ldots, \nu_{n}, \nu_{0}\right)$ is the outward unit normal to $\partial D_{T}$, $\nabla_{x}=\left(\frac{\partial}{\partial x_{1}}, \ldots, \frac{\partial}{\partial x_{n}}\right)$;
(ii) a function $u \in W_{2}^{1}\left(D_{T}\right)$ is called a strong generalized solution of the nonlinear problem (27),(28) in the domain $D_{T}$ if there exists a sequence of functions $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, S_{T}^{0}\right)=\left\{u \in C^{2}\left(\bar{D}_{T}\right):\left.u\right|_{S_{T}^{0}}=0\right\}$ such that $u_{k} \rightarrow u$ in the space $W_{2}^{1}\left(D_{T}\right),\left[\square u_{k}-\lambda\left|u_{k}\right|^{p+1}\right] \rightarrow F$ in the space $L_{2}\left(D_{T}\right)$ and $\left.u_{k}\right|_{S_{T}} \rightarrow g$ in the space $W_{2}^{1}\left(S_{T}\right)$.

Remark 4. In a standard manner [25, Ch. II, §5] one can prove that a weak generalized solution $u \in W_{2}^{1}\left(D_{T}\right)$ of problem (27),(28) satisfies the homogeneous boundary conditions (28) in the sense of trace theory.

It is obvious that a strong generalized solution $u \in W_{2}^{1}\left(D_{T}\right)$ of problem $(27),(28)$ is also a weak generalized solution of this problem.

Let us introduce into consideration a function $\varphi^{0}=\varphi^{0}(x, t)$ such that

$$
\begin{equation*}
\varphi^{0} \in C^{2}\left(\bar{D}_{T}\right),\left.\quad \varphi^{0}\right|_{D_{T=1}}>0,\left.\quad \varphi^{0}\right|_{x_{n}=0}=0,\left.\quad \varphi^{0}\right|_{t \geq 1}=0 \tag{30}
\end{equation*}
$$

and

$$
\begin{equation*}
\varkappa_{0}=\int_{D_{T=1}} \frac{\left|\square \varphi^{0}\right|^{\alpha^{\prime}}}{\left|\varphi^{0}\right| \alpha^{\prime}-1} d x d t<+\infty, \quad \alpha^{\prime}=1+\frac{1}{p} . \tag{31}
\end{equation*}
$$

It can be easily verified that for sufficiently large $k$ and $m$ we can take the function

$$
\varphi^{0}(x, t)= \begin{cases}x_{n}^{k}(1-t)^{m}, & (x, t) \in D_{T=1} \\ 0, & t \geq 1\end{cases}
$$

as a function $\varphi^{0}$ satisfying conditions (30) and (31).
If it is assumed $\varphi_{T}(x, t)=\varphi^{0}\left(\frac{x}{T}, \frac{t}{T}\right), T>0$, then by virtue of (30) we readily see that

$$
\begin{equation*}
\varphi_{T} \in C^{2}\left(\bar{D}_{T}\right),\left.\quad \varphi_{T}\right|_{D_{T}}>0,\left.\quad \varphi_{T}\right|_{x_{n}=0}=0,\left.\quad \varphi_{T}\right|_{t=T}=0 \tag{32}
\end{equation*}
$$

Assuming that the functions $F, g$ and $\varphi^{0}$ are fixed, we introduce into consideration the function of one variable $T$

$$
\begin{equation*}
\gamma(T)=\int_{D_{T}} F \varphi_{T} d x d t+\int_{S_{T}} g \frac{\partial \varphi_{T}}{\partial N} d s-\int_{S_{T}} \varphi_{T} \frac{\partial g}{\partial N} d s, \quad T>0 . \tag{33}
\end{equation*}
$$

We have the following theorem on the nonexistence of global solvability of problem (27),(28).

Theorem 2. Let $F \in L_{2, l o c}(D), g \in W_{2, l o c}^{1}(S)$ and $F \in L_{2}\left(D_{T}\right), g \in W_{2}^{1}\left(S_{T}\right)$ for any $T>0$. If $0<p<\frac{2}{n-1}$ and

$$
\begin{equation*}
\underline{\lim }_{T \rightarrow \infty} \gamma(T)>0 \tag{34}
\end{equation*}
$$

then there exists a positive number $T_{0}=T_{0}(F, g)$ such that, for $T>T_{0}$, problem (27), (28) cannot have a weak generalized solution $u \in W_{2}^{1}\left(D_{T}\right)$ in the domain $D_{T}$.

Proof. Let $u \in W_{2}^{1}\left(D_{T}\right)$ be a weak generalized solution of problem (27), (28) in the domain $D_{T}$, i.e. the integral equality (29) be fulfilled. By virtue of (32) we can take, in equality (29), the function $\varphi$ in the role of the test function $\varphi_{T}$. Integrating by parts the left-hand side of equality (29), we obtain

$$
\begin{align*}
& \int_{D_{T}} u \square \varphi_{T} d x d t=\lambda \int_{D_{T}}|u|^{p+1} \varphi_{T} d x d t \\
&+\int_{D_{T}} F \varphi_{T} d x d t+\int_{S_{T}} g \frac{\partial \varphi_{T}}{\partial N} d x-\int_{S_{T}} \varphi_{T} \frac{\partial g}{\partial N} d s \tag{35}
\end{align*}
$$

With (33) taken into account, equality (35) can be rewritten as

$$
\begin{equation*}
\lambda \int_{D_{T}}|u|^{p+1} \varphi_{T} d x d t=\int_{D_{T}} u \square \varphi_{T} d x d t-\gamma(T) \tag{36}
\end{equation*}
$$

If in the Young inequality with parameter $\varepsilon>0$ for $\alpha=p+1$

$$
a b \leq \frac{\varepsilon}{\alpha} a^{\alpha}+\frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} b^{\alpha^{\prime}}, \quad a, b \geq 0, \quad \alpha^{\prime}=\frac{\alpha}{\alpha-1}=1+\frac{1}{p}
$$

we take $a=|u| \varphi_{T}^{\frac{1}{\alpha}}, b=\frac{\left|\varphi_{T}\right|}{\varphi_{T}^{\frac{1}{T}}}$, then keeping in mind that $\frac{\alpha^{\prime}}{\alpha}=\alpha^{\prime}-1=\frac{1}{p}$, we obtain

$$
\begin{equation*}
\left|u \square \varphi_{T}\right|=|u| \varphi_{T}^{\frac{1}{\alpha}} \cdot \frac{\left|\square \varphi_{T}\right|}{\varphi_{T}^{\frac{1}{\alpha}}} \leq \frac{\varepsilon}{\alpha}|u|^{\alpha} \varphi_{T}+\frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \frac{\left|\square \varphi_{T}\right|^{\alpha^{\prime}}}{\varphi_{T}^{\alpha^{\alpha^{\prime}}-1}} \tag{37}
\end{equation*}
$$

By virtue of (37), from (36) we have

$$
\left(\lambda-\frac{\varepsilon}{\alpha}\right) \int_{D_{T}}|u|^{\alpha} \varphi_{T} d x d t \leq \frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{\alpha^{\prime}}}{\varphi_{T}^{\alpha^{\prime}-1}} d x d t-\gamma(T)
$$

whence for $\varepsilon<\lambda \alpha$ we obtain

$$
\begin{equation*}
\int_{D_{T}}|u|^{\alpha} \varphi_{T} d x d t \leq \frac{\alpha}{(\lambda \alpha-\varepsilon) \alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{\alpha^{\prime}}}{\varphi_{T}^{\alpha^{\prime}-1}} d x d t-\frac{\alpha}{\lambda \alpha-\varepsilon} \gamma(T) . \tag{38}
\end{equation*}
$$

Taking into account that $\alpha^{\prime}=\frac{\alpha}{\alpha-1}, \alpha=\frac{\alpha^{\prime}}{\alpha^{\prime}-1}$ and $\min _{0<\varepsilon<\lambda \alpha} \frac{\alpha}{(\lambda \alpha-\varepsilon) \alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}}=\frac{1}{\lambda^{\alpha^{\prime}}}$ for $\varepsilon=\lambda$, from (38) it follows that

$$
\begin{equation*}
\int_{D_{T}}|u|^{\alpha} \varphi_{T} d x d t \leq \frac{1}{\lambda^{\alpha^{\prime}}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{\alpha^{\prime}}}{\varphi_{T}^{\alpha^{\prime}-1}} d x d t-\frac{\alpha^{\prime}}{\lambda} \gamma(T) \tag{39}
\end{equation*}
$$

Since $\varphi_{T}(x, t)=\varphi^{0}\left(\frac{x}{T}, \frac{t}{T}\right)$, by virtue of (30), (31) it easy to verify after the substitution of the variables $t=T t^{\prime}, x=T x^{\prime}$ that

$$
\begin{equation*}
\int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{\alpha^{\prime}}}{\varphi_{T}^{\alpha^{\prime}-1}} d x d t=T^{n+1-2 \alpha^{\prime}} \int_{D_{T=1}} \frac{\left|\square \varphi^{0}\right|^{\alpha^{\prime}}}{\left(\varphi^{0}\right)^{\alpha^{\prime}-1}} d x^{\prime} d t^{\prime}=T^{n+1-2 \alpha^{\prime}} \varkappa_{0}<+\infty \tag{40}
\end{equation*}
$$

By (32) and (40), from inequality (39) we obtain

$$
\begin{equation*}
0 \leq \int_{D_{T}}|u|^{\alpha} \varphi_{T} d x d t \leq \frac{1}{\lambda^{\alpha^{\prime}}} T^{n+1-2 \alpha^{\prime}} \varkappa_{0}-\frac{\alpha^{\prime}}{\lambda} \gamma(T) \tag{41}
\end{equation*}
$$

If $p<\frac{2}{n-1}$, i.e. If $n+1-2 \alpha^{\prime}<0$, where $\alpha^{\prime}=1+\frac{1}{p}$, then by virtue of (31) we have $\lim _{T \rightarrow \infty} \frac{1}{\lambda^{\alpha^{\prime}}} T^{n+1-2 \alpha^{\prime}} \varkappa_{0}=0$. Hence by virtue of (34) there exists a positive number $T_{0}=T_{0}(F, g)$ such that, for $T>T_{0}$, the right-hand part of (41) is negative, while the left-hand part of this inequality is nonnegative. Thus if there exists a weak generalized solution $u \in W_{2}^{1}\left(D_{T}\right)$ of problem (27),(28) in the domain $D_{T}$, then necessarily $T \leq T_{0}$, which proves Theorem 2 .

Remark 5. We give some sufficient conditions imposed on the functions $F$ and $g$, which guarantee the fulfilment of condition (34):
(i) $F=$ const $>0, g=$ const;
(ii) $F \in L_{2, l o c}(D), g \in W_{2, l o c}^{1}(S)$ and $F \in L_{2}\left(D_{T}\right), g \in W_{2}^{1}\left(S_{T}\right)$ for any $T>0$, and also diam supp $g<+\infty$ and $F \geq 0, F(x, t) \geq c t^{-k}$ for $t \geq 1$, where $c=$ const $>0,0<k=$ const $<n+1$.

## 4. Local Solvability of Problem (1), (2) in the Case of

 Nonlinearity of the Form $f(u)=\lambda|u|^{p+1}$Remark 6. In proving Theorem 1, it was shown that the linear problem (21), which, for $m=0$, coincides with the corresponding linear problem (27), (28), has, for $\lambda=0$ and $g=0$, a unique solution $u=L^{-1} F$, where $L^{-1}: L_{2}\left(D_{T}\right) \rightarrow$ $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is a linear continuous operator whose norm admits estimate (24). It should also be noted that, analogously to Remark 1, for $0<p<\frac{2}{n-1}$ the operator

$$
\begin{equation*}
K_{1}: \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow L_{2}\left(D_{T}\right) \quad\left(K_{1} u=\lambda|u|^{p+1}\right) \tag{42}
\end{equation*}
$$

is continuous and compact. Thus for $g=0$ the nonlinear problem (27), (28) is equivalent to the functional equation

$$
\begin{equation*}
u=A u+u_{0} \tag{43}
\end{equation*}
$$

in the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$, where with (42) taken into account

$$
\begin{equation*}
A=L^{-1} K_{1}, \quad u_{0}=L^{-1} F \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \tag{44}
\end{equation*}
$$

Remark 7. Let $B(0, d):=\left\{u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right):\|u\|_{\dot{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \leq d\right\}$ be a closed (convex) ball in the Hilbert space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ with radius $d>0$ and center at a zero element. Since by virtue of Remark 6 the operator $A: \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ is continuous and compact for $0<p<$ $\frac{2}{n-1}$, by the Shauder principle in order to prove the solvability of equation (43) it is sufficient to show that the operator $A_{1}$ acting by the formula $A_{1} u=A u+u_{0}$ transverse the ball $B(0, d)$ into itself for some $d>0$ [30, Ch. VIII, § 35.3]. To this end, below we will derive the needed estimate for the value $\|A u\|_{\dot{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}$.

We further use the reasoning from [31]. If $u \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$, then we denote by $\widetilde{u}$ the function which continues in an even manner the function $u$ across the plane $t=T$. It is obvious that $\widetilde{u} \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}^{*}\right)$, where $D_{T}^{*}:|x|<t<$ $2 T-|x|, x_{n}>0$.

Using the inequality [32, Ch. X, § 1]

$$
\int_{\Omega}|v| d \Omega=(\operatorname{mes} \Omega)^{1-\frac{1}{q}}\|v\|_{q, \Omega}, \quad q \geq 1
$$

and taking into account the equalities

$$
\|\widetilde{u}\|_{L_{q}\left(D_{T}^{*}\right)}^{q}=2\|u\|_{L_{q}\left(D_{T}\right)}^{q}, \quad\|\widetilde{u}\|_{W_{2}^{1}}^{2}=2\|u\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}^{2}
$$

from the well known multiplicative inequality [25, Ch. I, § 7]

$$
\begin{aligned}
\|v\|_{q, \Omega} & \leq \beta\|\nabla v\|_{m, \Omega}^{\widetilde{\alpha}}\|v\|_{r, \Omega}^{1-\widetilde{\alpha}} \quad \forall v \in \stackrel{\circ}{W}_{2}^{1}(\Omega), \quad \Omega \subset R^{n+1}, \\
\widetilde{\alpha} & =\left(\frac{1}{r}-\frac{1}{q}\right)\left(\frac{1}{r}-\frac{1}{\widetilde{m}}\right)^{-1}, \quad \widetilde{m}=\frac{(n+1) m}{n+1-m}
\end{aligned}
$$

for $\Omega=D_{T}^{*} \subset R^{n+1}, v=\widetilde{u}, r=1, m=2$ and $1<q \leq \frac{2(n+1)}{n-1}$, where $\beta=$ const $>0$ does not depend on $v$ and $T$, we obtain the inequality

$$
\begin{equation*}
\|u\|_{L_{q}\left(D_{T}\right)} \leq c_{0}\left(\operatorname{mes} D_{T}\right)^{\frac{1}{q}+\frac{1}{n+1}-\frac{1}{2}}\|u\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \quad \forall u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \tag{45}
\end{equation*}
$$

where $c_{0}=$ const $>0$ does not depend on $u$.
Since mes $D_{T}=\frac{\omega_{n}}{2(n+1)} T^{n+1}$, where $\omega_{n}$ is the volume of the unit ball in $R^{n}$, for $q=2(p+1)$ inequality (45) implies

$$
\begin{align*}
\|u\|_{L_{2(p+1)}\left(D_{T}\right)} \leq c_{0} \widetilde{\ell}_{p, n} T^{(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}\|u\|_{\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \\
\forall u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \tag{46}
\end{align*}
$$

where $\widetilde{\ell}_{p, n}=\left(\frac{\omega_{n}}{2(n+1)}\right)^{\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}$.
For the value $\left\|K_{1} u\right\|_{L_{2}\left(D_{T}\right)}$, where $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ and the operator $K_{1}$ is given by equality (42), by virtue of (46) we obtain the estimate

$$
\begin{align*}
\left\|K_{1} u\right\|_{L_{2}\left(D_{T}\right)} & \leq \lambda\left[\int_{D_{T}}|u|^{2(p+1)} d x d t\right]^{\frac{1}{2}}=\lambda\|u\|_{L_{2(p+1)}\left(D_{T}\right)}^{2} \\
& \leq \lambda \ell_{p, n} T^{(p+1)(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}\|u\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}^{p+1} \tag{47}
\end{align*}
$$

where $\ell_{p, n}=\left[c_{0} \widetilde{\ell}_{p, n}\right]^{p+1}$.
Now, for $A u=L^{-1} K_{1} u$, from (24) and (47) follows the estimate

$$
\begin{align*}
\|A u\|_{\stackrel{\circ}{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}} \leq & \left\|L^{-1}\right\|_{L_{2}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}}\left\|K_{1} u\right\|_{L_{2}\left(D_{T}\right)} \\
\leq & \sqrt{\frac{e}{2}} \lambda \ell_{p, n} T^{1+(p+1)(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}\|u\|_{\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}^{p+1} \\
& \forall u \in \stackrel{\circ}{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) .} \tag{48}
\end{align*}
$$

Note that $\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}>0$ for $p<\frac{2}{n-1}$.
Consider the equation

$$
\begin{equation*}
a z^{p+1}+b=z \tag{49}
\end{equation*}
$$

with respect to the unknown $z$, where

$$
\begin{equation*}
a=\sqrt{\frac{e}{2}} \lambda \ell_{p, n} T^{1+(p+1)(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}, \quad b=\sqrt{\frac{e}{2}} T\|F\|_{L_{2}\left(D_{T}\right)} \tag{50}
\end{equation*}
$$

For $T>0$ it is obvious that $a>0$ and $b \geq 0$. A simple analysis analogous to that carried out for $p=2$ in [30, Ch. VIII, §35.4] shows that (i) in the case $b=0$ equation (49) has, along with the zero root $z_{1}=0$, the unique positive root $z_{2}=a^{-\frac{1}{p}}$; (ii) if $b>0$, then for $0<b<b_{0}$, where

$$
\begin{equation*}
b_{0}=\left[(p+1)^{-\frac{1}{p}}-(p+1)^{-\frac{p+1}{p}}\right] a^{-\frac{1}{p}}, \tag{51}
\end{equation*}
$$

equation (49) has two positive roots $z_{1}$ and $z_{2}, 0<z_{1}<z_{2}$. For $b=b_{0}$ these roots coincide and we have one positive root $z_{1}=z_{2}=z_{0}=[(p+1) a]^{-\frac{1}{p}}$; (iii) for $b>b_{0}$ equation (49) has no nonnegative roots.

Note that for $0<b<b_{0}$ we have the inequalities $z_{1}<z_{0}=[(p+1) a]^{-\frac{1}{p}}<z_{2}$.
By virtue of (50) and (51) the condition $b \leq b_{0}$ is equivalent to the condition
$\sqrt{\frac{e}{2}} T\|F\|_{L_{2}\left(D_{T}\right)}$

$$
\leq\left[\sqrt{\frac{e}{2}} \lambda \ell_{p, n} T^{1+(p+1)(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)}\right]^{-\frac{1}{p}}\left[(p+1)^{-\frac{1}{p}}-(p+1)^{-\frac{p+1}{p}}\right]
$$

or to

$$
\begin{equation*}
\|F\|_{L_{2}\left(D_{T}\right)} \leq \gamma_{n, \lambda, p} T^{-\alpha_{n}}, \quad \alpha_{n}>0 \tag{52}
\end{equation*}
$$

where

$$
\begin{aligned}
\gamma_{n, \lambda, p} & =\left[(p+1)^{-\frac{1}{p}}-(p+1)^{-\frac{p+1}{p}}\right]\left(\lambda \ell_{p, n}\right)^{-\frac{1}{p}} \exp \left[-\frac{1}{2}\left(1+\frac{1}{p}\right)\right], \\
\alpha_{n} & =1+\frac{1}{p}\left[1+(p+1)(n+1)\left(\frac{1}{2(p+1)}+\frac{1}{n+1}-\frac{1}{2}\right)\right]
\end{aligned}
$$

By the absolute continuity of the Lebesgue integral we have $\lim _{T \rightarrow 0}\|F\|_{L_{2}\left(D_{T}\right)}=$ 0 . Since at the same time time $\lim _{T \rightarrow 0} T^{-\alpha_{n}}=+\infty$, there exists a number $T_{1}=$ $T_{1}(F), 0<T_{1}<+\infty$, such that inequality (52) is fulfilled for

$$
\begin{equation*}
0<T \leq T_{1}(F) \tag{53}
\end{equation*}
$$

Now we will show that if condition (53) is fulfilled, then the operator $A_{1}$ : $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right) \rightarrow \stackrel{\circ}{W}_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ acting by the formula $A_{1} u=A u+u_{0}$ transfers the ball $B\left(0, z_{2}\right)$ from Remark 7 into itself, where $z_{2}$ is the largest positive root of equation (49). Indeed, if $u \in B\left(0, z_{2}\right)$, then by virtue of (48)(50) we have

$$
\left\|A_{1} u\right\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)} \leq a\|u\|_{W_{2}^{1}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)}^{p+1}+b \leq a z_{2}^{p+1}+b=z_{2} .
$$

Therefore the following theorem is valid according to Remarks 6 and 7.

Theorem 3. Let $F \in L_{2, l o c}(D), g=0,0<p<\frac{2}{n-1}$ and condition (53) be fulfilled for the value $T$. Then problem (27), (28) has at least one strong generalized solution $u \in \stackrel{\circ}{W_{2}^{1}}\left(D_{T}, S_{T}^{0} \cup S_{T}\right)$ in the domain $D_{T}$.

## References

1. A. V. Bitsadze, Some classes of partial differential equations. (Russian) Nauka, Moscow, 1981.
2. K. Jörgens, Das Anfangswertproblem im Grossen für eine Klasse nichtlinearer Wellengleichungen. Math. Z. 77(1961), 295-308.
3. H. A. Levine, Instability and nonexistence of global solutions to nonlinear wave equations of the form $P u_{t t}=-A u+\mathcal{F}(u)$. Trans. Amer. Math. Soc. 192(1974), 1-21.
4. F. John, Blow-up of solutions of nonlinear wave equations in three space dimensions. Manuscripta Math. 28(1979), No. 1-3, 235-268.
5. F. John, Blow-up for quasilinear wave equations in three space dimensions. Comm. Pure Appl. Math. 34(1981), No. 1, 29-51.
6. F. John and S. Klainerman, Almost global existence to nonlinear wave equations in three space dimensions. Comm. Pure Appl. Math. 37(1984), No. 4, 443-455.
7. T. Kato, Blow-up of solutions of some nonlinear hyperbolic equations. Comm. Pure Appl. Math. 33(1980), No. 4, 501-505.
8. J. Ginibre, A. Soffer, and G. Velo, The global Cauchy problem for the critical nonlinear wave equation. J. Funct. Anal. 110(1992), No. 1, 96-130.
9. W. A. Strauss, Nonlinear scattering theory at low energy. J. Funct. Anal. 41(1981), No. 1, 110-133.
10. V. Georgiev, H. Lindblad, and C. D. Sogge, Weighted Strichartz estimates and global existence for semilinear wave equations. Amer. J. Math. 119(1997), No. 6, 12911319.
11. T. C. Sideris, Nonexistence of global solutions to semilinear wave equations in high dimensions. J. Differential Equations 52(1984), No. 3, 378-406.
12. L. Hörmander, Lectures on nonlinear hyperbolic differential equations. Mathématiques \& Applications (Berlin) [Mathematics \& Applications], 26. Springer-Verlag, Berlin, 1997.
13. F. Merle and H. ZaAg, Determination of the blow-up rate for a critical semilinear wave equation. Math. Ann. 331(2005), No. 2, 395-416.
14. È. Mitidieri and S. I. Pokhozaev, A priori estimates and the absence of solutions of nonlinear partial differential equations and inequalities. (Russian) Tr. Mat. Inst. Steklova 234(2001), 1-384; English transl.: Proc. Steklov Inst. Math. 2001, No. 3 (234), 1-362.
15. E. Belchev, M. Kepka, and Z. Zhou, Finite-time blow-up of solutions to semilinear wave equations. Special issue dedicated to the memory of I. E. Segal. J. Funct. Anal. 190(2002), No. 1, 233-254.
16. G. Todorova and E. Vitillaro, Blow-up for nonlinear dissipative wave equations in $\mathbb{R}^{n}$. J. Math. Anal. Appl. 303(2005), No. 1, 242-257.
17. M. Keel, H. F. Smith, and C. D. Sogge, Almost global existence for quasilinear wave equations in three space dimensions. J. Amer. Math. Soc. 17(2004), No. 1, 109-153 (electronic).
18. A. V. Bitsadze, Equations of mixed type in three-dimensional regions. (Russian) Dokl. Akad. Nauk SSSR 143(1962), 1017-1019.
19. V. N. Vragov, Boundary value problems for nonclassical equations of mathematical physics. (Russian) Novosibirsk State University, Novosibirsk, 1983.
20. T. Sh. Kal'menov, Multidimensional regular boundary value problems for the wave equation. (Russian) Izv. Akad. Nauk Kazakh. SSR Ser. Fiz.-Mat. 1982, No. 3, 18-25.
21. L. I. Schiff, Nonlinear meson theory of nuclear forces. I. Neutral scalar mesons with point-contact repulsion. Phys. Rev., II. Ser. 84(1951), 1-9.
22. I. E. Segal, The global Cauchy problem for a relativistic scalar field with power interaction. Bull. Soc. Math. France 91(1963), 129-135.
23. J.-L. Lions, Quelques méthodes de résolution de problèmes aux limites non linèaires. Dunod, Gauthier-Villars, Paris, 1969.
24. M. Reed and B. Simon, Methods of modern mathematical physics. II. Fourier analysis, self-adjointness. Academic Press [Harcourt Brace Jovanovich, Publishers], New YorkLondon, 1975.
25. O. A. Ladyzhenskaya, Boundary value problems of mathematical physics. (Russian) Nauka, Moscow, 1973.
26. M. A. Krasnosel'skĭ̌, P. P. Zabreiko, E. I. Pustyl'nic, and P. E. Sobolevskĭ́, Integral operators in spaces of summable functions. (Russian) Nauka, Moscow, 1966.
27. S. FučIk and A. Kufner, Nonlinear differential equations. Studies in Applied Mechanics, 2. Elsevier Scientific Publishing Co., Amsterdam-New York, 1980; Russian transl.: Nauka, Moscow, 1988.
28. D. Henry, Geometric theory of semilinear parabolic equations. Lecture Notes in Mathematics, 840. Springer-Verlag, Berlin-New York, 1981; Russian transl.: Mir, Moscow, 1985.
29. L. HÖRMANDER, Linear partial differential operators. Die Grundlehren der mathematischen Wissenschaften, Bd. 116. Academic Press, Inc., Publishers, New York; SpringerVerlag, Berlin-Göttingen-Heidelberg, 1963; Russian transl.: Mir, Moscow, 1965.
30. V. A. Trenogin, Functional analysis. (Russian) Nauka, Moscow, 1993.
31. S. S. Kharibegashvili, On the nonexistence of global solutions of the characteristic Cauchy problem for a nonlinear wave equation in a conical domain. (Russian) Differ. Uravn. 42(2006), No. 2, 261-271, 288.
32. B. Z. Vulikh, A short course of the theory of functions of a real variable. (Russian) Nauka, Moscow, 1973.
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