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# ON ONE BOUNDARY VALUE PROBLEM FOR A NONLINEAR EQUATION WITH THE ITERATED WAVE OPERATOR IN THE PRINCIPAL PART 

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Dedicated to the memory of Professor J.-L. Lions


#### Abstract

One boundary value problem for a hyperbolic equation with power nonlinearity and the iterated wave operator in the principal part is considered in a conical domain. Depending on the index of nonlinearity and spatial dimensionality of the equation the question on the existence and uniqueness of a solution of a boundary value problem is investigated. The question as to the absence of a solution of this problem is also considered.


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## 1. Statement of the Problem

In the Euclidean space $R^{n+1}$ of independent variables $x_{1}, x_{2}, \ldots, x_{n}, t$ consider a nonlinear equation of the form

$$
\begin{equation*}
L_{\lambda} u:=\square^{2} u=\lambda f(u)+F, \tag{1.1}
\end{equation*}
$$

where $\lambda$ is a given real constant, $f: R \rightarrow R$ is a given continuous nonlinear function, $f(0)=0, F$ is a given, and $u$ is an unknown real function, $\square=$ $\frac{\partial^{2}}{\partial t^{2}}-\sum_{i=1}^{n} \frac{\partial^{2}}{\partial x_{i}^{2}}, n \geq 2$.

Let $D_{T}:|x|<t<T$ be the domain which is the intersection of the light cone of future $K_{O}^{+}: t>|x|$ with apex at the origin $O(0,0, \ldots, 0)$ and a half-space $H_{T}: t<T, T=$ const $>0$. Assume $S_{T}=\partial D_{T} \cap \partial K_{O}^{+}, S_{T}^{0}=\partial D_{T} \cap \partial H_{T}$. It is obvious that $S_{T}: t=|x|, 0 \leq t \leq T$, is a characteristic conical manifold for equation (1.1), $S_{T}^{0}:|x| \leq T, t=T$, and $\partial D_{T}=S_{T} \cup S_{T}^{0}$.

For equation (1.1) consider the boundary value problem on determination of its solution $u\left(x_{1}, \ldots, x_{n}, t\right)$ in the domain $D_{T}$ by the boundary conditions

$$
\begin{gather*}
\left.u\right|_{\partial D_{T}}=0  \tag{1.2}\\
\left.\frac{\partial u}{\partial t}\right|_{S_{T}^{0}}=0 . \tag{1.3}
\end{gather*}
$$

It should be noted that for nonlinear hyperbolic equations the question of the local or global solvability of the Cauchy problem with the initial conditions for $t=0$ is considered in vast literature [see, e.g., 1-20].

As is known, for a nonlinear wave equation of the form $\square u=\lambda f(u)+F$ the characteristic problem in the light cone of future $K_{O}^{+}: t>|x|$ with a boundary condition of the form $\left.u\right|_{\partial K_{O}^{+}}=g$ in the linear case, i.e., with $\lambda=0$ is well-posed and the global solvability takes place in appropriate function spaces [21-25], while in the nonlinear case, when the function $f(u)$ has exponential nature and $\lambda \neq 0$, this problem was considered in [26-28].

Assume $\stackrel{\circ}{C^{k}}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)=\left\{u \in C^{k}\left(\bar{D}_{T}\right):\left.u\right|_{\partial D_{T}}=0,\left.\frac{\partial u}{\partial t}\right|_{S_{T}^{0}}=0\right\}, k \geq 1$.
Let $u \in \stackrel{\circ}{C^{4}}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ be a classical solution of problem (1.1), (1.2), (1.3). Multiplying both parts of equation (1.1) by an arbitrary function $\phi \in$ ${ }^{\circ}{ }^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ and integrating the resulting equation by parts in the domain $D_{T}$ we obtain

$$
\begin{equation*}
\int_{D_{T}} \square u \square \phi d x d t=\lambda \int_{D_{T}} f(u) \phi d x d t+\int_{D_{T}} F \phi d x d t \tag{1.4}
\end{equation*}
$$

Here we have used the equality

$$
\int_{D_{T}} \square u \square \phi d x d t=\int_{\partial D_{T}} \frac{\partial \phi}{\partial N} \square u d s-\int_{\partial D_{T}} \phi \frac{\partial}{\partial N} \square u d s+\int_{D_{T}} \phi \square^{2} u d x d t
$$

and the fact that since $S_{T}=\partial D_{T} \cap \partial K_{O}^{+}$is a characteristic manifold, the derivative on the conormal $\frac{\partial}{\partial N}=\gamma_{n+1} \frac{\partial}{\partial t}-\sum_{i=1}^{n} \gamma_{i} \frac{\partial}{\partial x_{i}}$, where $\gamma=\left(\gamma_{1}, \ldots, \gamma_{n}, \gamma_{n+1}\right)$ is the unit vector of the outer normal to $\partial D_{T}$, is an inner differential operator on the characteristic manifold $S_{T}$ and thus if $v \in C^{1}\left(\bar{D}_{T}\right)$ and $\left.v\right|_{S_{T}}=0$, then $\left.\frac{\partial v}{\partial N}\right|_{S_{T}}=0$.

Let us introduce the Hilbert space ${ }_{W}^{\circ}{ }_{2, \square}^{1}\left(D_{T}\right)$ as the completion with respect to the norm

$$
\begin{equation*}
\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2}=\int_{D_{T}}\left[u^{2}+\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}+(\square u)^{2}\right] d x d t \tag{1.5}
\end{equation*}
$$

of the classical space $\stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$. It follows from (1.5) that if $u \in$ $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$, then $u \in \stackrel{\circ}{W}_{2}^{1}\left(D_{T}\right)$ and $\square u \in L_{2}\left(D_{T}\right)$. Here $W_{2}^{1}\left(D_{T}\right)$ is the known Sobolev space [29, p. 56] consisting of elements from $L_{2}\left(D_{T}\right)$, which have first order generalized derivatives in $L_{2}\left(D_{T}\right)$, and ${ }_{W}^{\circ}{ }_{2}^{1}\left(D_{T}\right)=\left\{u \in W_{2}^{1}\left(D_{T}\right)\right.$ : $\left.\left.u\right|_{\partial D_{T}}=0\right\}$, where the equality $\left.u\right|_{\partial D_{T}}=0$ should be understood in the sense of the trace theory [29, p. 70]. Moreover, since from (1.5) it follows that $u \in \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}\right)$ and $\square u \in L_{2}\left(D_{T}\right)$, then an arbitrary function $u$ from the space $\stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$ satisfies the homogeneous condition (1.3) in the sense of the trace theory [29, p. 130].

Let us assume that equality (1.4) underlies the determination of a generalized solution of problem (1.1), (1.2), (1.3).

Definition 1. Let $F \in L_{2}\left(D_{T}\right)$. We call a function $u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ a weak generalized solution of problem (1.1), (1.2), (1.3) if $f(u) \in L_{2}\left(D_{T}\right)$ and for any function $\phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ the integral equality (1.4) is valid, i.e.

$$
\begin{equation*}
\int_{D_{T}} \square u \square \phi d x d t=\lambda \int_{D_{T}} f(u) \phi d x d t+\int_{D_{T}} F \phi d x d t \quad \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{1.6}
\end{equation*}
$$

It is easy to verify that if a solution $u$ of problem (1.1), (1.2), (1.3) in the sense of Definition 1 belongs to the class $C^{4}\left(\bar{D}_{T}\right)$, then it will be a classical solution of this problem.
2. Solvability of Problem (1.1), (1.2), (1.3) in the Case of A Nonlinearity of the Form $f(u)=|u|^{\alpha} \operatorname{sgn} u$
Assume that in equation (1.1) the nonlinear function $f$ has the form

$$
\begin{equation*}
f(u)=|u|^{\alpha} \operatorname{sgn} u, \quad \alpha=\text { const }>0, \quad \alpha \neq 1 \tag{2.1}
\end{equation*}
$$

Then in accordance with (2.1) equation (1.1) and the integral equality (1.6) take the form

$$
\begin{equation*}
L_{\lambda} u:=\square^{2} u=\lambda|u|^{\alpha} \operatorname{sgn} u+F \tag{2.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{D_{T}} \square u \square \phi d x d t=\lambda \int_{D_{T}} \phi|u|^{\alpha} \operatorname{sgn} u d x d t+\int_{D_{T}} F \phi d x d t \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{2.3}
\end{equation*}
$$

Lemma 1. The inequality

$$
\begin{equation*}
\|u\|_{\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)} \leq c\|\square u\|_{L_{2}\left(D_{T}\right)} \forall u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{2.4}
\end{equation*}
$$

is valid, where the norm of the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ is given by equality (1.5) and the positive constant $c$ does not depend on $u$.

Proof. Let $\Omega_{\tau}:=\bar{D}_{T} \cap\{t=\tau\}, D_{\tau}=D_{T} \cap\{t<\tau\}, S_{\tau}=\left\{(x, t) \in \partial D_{\tau}\right.$ : $t=|x|\}, 0<\tau \leq T$ and $\gamma=\left(\gamma_{1}, \ldots, \gamma_{n}, \gamma_{n+1}\right)$ be the unit vector of the outer normal to $\partial D_{\tau}$.

Since the space $\stackrel{\circ}{C^{2}}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ is a dense subspace of the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ it is sufficient to prove inequality (2.4) for functions from the space $\stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$. For $u \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$, taking into account the equalities $\left.u\right|_{S_{\tau}}=0, \Omega_{\tau}=\partial D_{\tau} \cap\{t=\tau\}$ and $\left.\gamma\right|_{\Omega_{\tau}}=(0, \ldots, 0,1)$, it is easy to obtain
by integration by parts

$$
\begin{align*}
\int_{D_{\tau}} \frac{\partial^{2} u}{\partial t^{2}} \frac{\partial u}{\partial t} d x d t= & \frac{1}{2} \int_{D_{\tau}} \frac{\partial}{\partial t}\left(\frac{\partial u}{\partial t}\right)^{2} d x d t=\frac{1}{2} \int_{\partial D_{\tau}}\left(\frac{\partial u}{\partial t}\right)^{2} \gamma_{n+1} d s \\
= & \frac{1}{2} \int_{\Omega_{\tau}}\left(\frac{\partial u}{\partial t}\right)^{2} d x+\frac{1}{2} \int_{S_{\tau}}\left(\frac{\partial u}{\partial t}\right)^{2} \gamma_{n+1} d s, \quad \tau \leq T,  \tag{2.5}\\
\int_{D_{\tau}} \frac{\partial^{2} u}{\partial x_{i}^{2}} \frac{\partial u}{\partial t} d x d t= & \int_{\partial D_{\tau}} \frac{\partial u}{\partial x_{i}} \frac{\partial u}{\partial t} \gamma_{i} d s-\frac{1}{2} \int_{D_{\tau}} \frac{\partial}{\partial \tau}\left(\frac{\partial u}{\partial x_{i}}\right)^{2} d x d t \\
= & \int_{\partial D_{\tau}} \frac{\partial u}{\partial x_{i}} \frac{\partial u}{\partial t} \gamma_{i} d s-\frac{1}{2} \int_{\partial D_{\tau}}\left(\frac{\partial u}{\partial x_{i}}\right)^{2} \gamma_{n+1} d s \\
= & \int_{\partial D_{\tau}} \frac{\partial u}{\partial x_{i}} \frac{\partial u}{\partial t} \gamma_{i} d s--\frac{1}{2} \int_{S_{\tau}}\left(\frac{\partial u}{\partial x_{i}}\right)^{2} \gamma_{n+1} d s \\
& -\frac{1}{2} \int_{\Omega_{\tau}}\left(\frac{\partial u}{\partial x_{i}}\right)^{2} d s, \quad \tau \leq T . \tag{2.6}
\end{align*}
$$

It follows from (2.5) and (2.6) that

$$
\begin{align*}
\int_{D_{\tau}} \square u \frac{\partial u}{\partial t} d x d t= & \int_{S_{\tau}} \frac{1}{2 \gamma_{n+1}}\left[\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}} \gamma_{n+1}-\frac{\partial u}{\partial t} \gamma_{i}\right)^{2}\right. \\
& \left.+\left(\frac{\partial u}{\partial t}\right)^{2}\left(\gamma_{n+1}^{2}-\sum_{j=1}^{n} \gamma_{j}^{2}\right)\right] d s \\
& +\frac{1}{2} \int_{\Omega_{\tau}}\left[\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}\right] d x, \quad \tau \leq T . \tag{2.7}
\end{align*}
$$

Since $\left.u\right|_{S_{\tau}}=0$ and the operator $\left(\gamma_{n+1} \frac{\partial}{\partial x_{i}}-\gamma_{i} \frac{\partial}{\partial t}\right), 1 \leq i \leq n$, is an internal differential operator on $S_{\tau}$, we have the equalities

$$
\begin{equation*}
\left.\left(\frac{\partial u}{\partial x_{i}} \gamma_{n+1}-\frac{\partial u}{\partial t} \gamma_{i}\right)\right|_{S_{\tau}}=0, \quad i=1, \ldots, n . \tag{2.8}
\end{equation*}
$$

Therefore, taking into account that $\gamma_{n+1}^{2}-\sum_{j=1}^{n} \gamma_{j}^{2}=0$ on the characteristic manifold $S_{\tau}$, in view of (2.7) and (2.8) we have

$$
\begin{equation*}
\int_{\Omega_{\tau}}\left[\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}\right] d x=2 \int_{D_{\tau}} \square u \frac{\partial u}{\partial t} d x d t, \quad \tau \leq T . \tag{2.9}
\end{equation*}
$$

Assuming $w(\delta)=\int_{\Omega_{\delta}}\left[\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}\right] d x$, and using inequality $2 \square u \frac{\partial u}{\partial t} \leq$ $\varepsilon\left(\frac{\partial u}{\partial t}\right)^{2}+\frac{1}{\varepsilon}|\square u|^{2}$, which is valid for any $\varepsilon=$ const $>0$, from (2.9) we obtain

$$
\begin{equation*}
w(\delta) \leq \varepsilon \int_{0}^{\delta} w(\sigma) d \sigma+\frac{1}{\varepsilon}\|\square\|_{L_{2}\left(D_{\delta}\right)}^{2}, \quad 0<\delta \leq T \tag{2.10}
\end{equation*}
$$

From (2.10), taking into account that the value $\|\square\|_{L_{2}\left(D_{\delta}\right)}^{2}$ as a function of $\delta$ is non-decreasing, in view of Gronwall's lemma [30, p. 13] it follows that

$$
w(\delta) \leq \frac{1}{\varepsilon}\|\square\|_{L_{2}\left(D_{\delta}\right)}^{2} \exp \delta \varepsilon
$$

Hence, taking into account the fact that $\inf _{\varepsilon>0} \frac{1}{\varepsilon} \exp \delta \varepsilon=e \delta$ and it is reached at $\varepsilon=\frac{1}{\delta}$, we obtain

$$
\begin{equation*}
w(\delta) \leq e \delta\|\square\|_{L_{2}\left(D_{\delta}\right)}^{2}, \quad 0<\delta \leq T \tag{2.11}
\end{equation*}
$$

From (2.11), in turn, it follows that

$$
\begin{equation*}
\int_{D_{T}}\left[\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}\right] d x d t=\int_{0}^{T} w(\delta) d \delta \leq \frac{e}{2} T^{2}\|\square u\|_{L_{2}\left(D_{T}\right)}^{2} \tag{2.12}
\end{equation*}
$$

Using the equalities $\left.u\right|_{S_{T}}=0$ and $u(x, t)=\int_{|x|}^{t} \frac{\partial u(x, t)}{\partial t} d \tau,(x, t) \in \bar{D}_{T}$, which are valid for any function $u \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$, by a standard reasoning [29, p. 63] we easily obtain the inequality

$$
\begin{equation*}
\int_{D_{T}} u^{2}(x, t) d x d t \leq T^{2} \int_{D_{T}}\left(\frac{\partial u}{\partial t}\right)^{2} d x d t \tag{2.13}
\end{equation*}
$$

By virtue of (2.11) and (2.13) we have

$$
\begin{aligned}
\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2} & =\int_{D_{T}}\left[u^{2}+\left(\frac{\partial u}{\partial t}\right)^{2}+\sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}}\right)^{2}+(\square u)^{2}\right] d x d t \\
& \leq\left(1+\frac{e}{2} T^{2}+\frac{e}{2} T^{4}\right)\|\square\|_{L_{2}\left(D_{T}\right)}^{2}
\end{aligned}
$$

whence inequality (2.4) with the constant $c^{2}=1+\frac{e}{2} T^{2}+\frac{e}{2} T^{4}$ follows.
Lemma 2. Assume $F \in L_{2}\left(D_{T}\right), 0<\alpha<1$, and in the case $\alpha>1$ additionally require that $\lambda<0$. Then in the case with a nonlinearity of form (2.1) for a weak generalized solution $u \in \stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$ of problem (1.1), (1.2), (1.3) , i.e., problem (2.2), (1.2), (1.3) in the sense of the integral equality (2.3) with $|u|^{\alpha} \in L_{2}\left(D_{T}\right)$, we have an a priori estimate

$$
\begin{equation*}
\|u\|_{\dot{W}_{2, \square}^{1}\left(D_{T}\right)} \leq c_{1}\|F\|_{L_{2}\left(D_{T}\right)}+c_{2} \tag{2.14}
\end{equation*}
$$

with non-negative constants $c_{i}(T, \alpha, \lambda), i=1,2$, which do not depend on $u, F$ and $c_{1}>0$.

Proof. First, let $\alpha>1$ and $\lambda<0$. Assuming in equality (2.3) that $\phi=u \in$ $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ and taking into account (1.5), for any $\varepsilon>0$ we have

$$
\begin{align*}
\|\square u\|_{L_{2}\left(D_{T}\right)}^{2} & =\int_{D_{T}}(\square u)^{2} d x d t=\lambda \int_{D_{T}}|u|^{\alpha+1} d x d t+\int_{D_{T}} F u d x d t \\
& \leq \int_{D_{T}} F u d x d t \leq \frac{1}{4 \varepsilon} \int_{D_{T}} F^{2} d x d t+\varepsilon\|u\|_{L_{2}\left(D_{T}\right)}^{2} \\
& \leq \frac{1}{4 \varepsilon}\|F\|_{L_{2}\left(D_{T}\right)}^{2}+\varepsilon\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2} \tag{2.15}
\end{align*}
$$

Due to (2.4) and (2.15) we have

$$
\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2} \leq c^{2}\|\square u\|_{L_{2}\left(D_{T}\right)}^{2} \leq \frac{c^{2}}{4 \varepsilon}\|F\|_{L_{2}\left(D_{T}\right)}^{2}+c^{2} \varepsilon\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2},
$$

from which for $\varepsilon=\frac{1}{2 c^{2}}<\frac{1}{c^{2}}$ we obtain

$$
\begin{equation*}
\|u\|_{\dot{W}_{2, \square}^{1}\left(D_{T}\right)}^{2} \leq \frac{c^{2}}{4 \varepsilon\left(1-\varepsilon c^{2}\right)}\|F\|_{L_{2}\left(D_{T}\right)}^{2}=c^{4}\|F\|_{L_{2}\left(D_{T}\right)}^{2} . \tag{2.16}
\end{equation*}
$$

From (2.16) in the case $\alpha>1$ and $\lambda<0$ follows inequality (2.14) with $c_{1}=c^{2}$ and $c_{2}=0$.

Let now $0<\alpha<1$. Using the known inequality

$$
a b \leq \frac{\varepsilon a^{p}}{p}+\frac{b^{q}}{q \varepsilon^{q-1}}
$$

with a parameter $\varepsilon>0$ for $a=|u|^{\alpha+1}, b=1, p=\frac{2}{\alpha+1}>1, q=\frac{2}{1-\alpha}, \frac{1}{p}+\frac{1}{q}=1$, in the same way as for inequality (2.15) we have

$$
\begin{align*}
& \|\square u\|_{L_{2}\left(D_{T}\right)}^{2}=\int_{D_{T}}(\square u)^{2} d x d t=\lambda \int_{D_{T}}|u|^{\alpha+1} d x d t+\int_{D_{T}} F u d x d t \\
& \quad \leq|\lambda| \int_{D_{T}}\left[\varepsilon \frac{1+\alpha}{2}|u|^{2}+\frac{1-\alpha}{2 \varepsilon^{q-1}}\right] d x d t+\frac{1}{4 \varepsilon} \int_{D_{T}} F^{2} d x d t+\varepsilon \int_{D_{T}} u^{2} d x d t \\
& \quad=\frac{1}{4 \varepsilon}\|F\|_{L_{2}\left(D_{T}\right)}^{2}+\varepsilon\left(|\lambda| \frac{1+\alpha}{2}+1\right)\|u\|_{L_{2}\left(D_{T}\right)}^{2}+|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \operatorname{mes} D_{T} \tag{2.17}
\end{align*}
$$

In view of (1.5) and (2.4) it follows from (2.17) that

$$
\begin{aligned}
& \|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2} \leq c^{2}\|\square u\|_{L_{2}\left(D_{T}\right)}^{2} \leq \frac{c^{2}}{4 \varepsilon}\|F\|_{L_{2}\left(D_{T}\right)}^{2} \\
& \quad+\varepsilon c^{2}\left(|\lambda| \frac{1+\alpha}{2}+1\right)\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2}+c^{2}|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \operatorname{mes} D_{T}, \quad q=\frac{2}{1-\alpha},
\end{aligned}
$$

whence for $\varepsilon=\frac{1}{2} c^{-2}\left(|\lambda| \frac{1+\alpha}{2}+1\right)^{-1}$ we obtain

$$
\begin{align*}
\|u\|_{W_{2, \square}^{1}\left(D_{T}\right)}^{2} \leq & {\left[1-\varepsilon c^{2}\left(|\lambda| \frac{1+\alpha}{2}+1\right)\right]^{-1} } \\
& \times\left(\frac{c^{2}}{4 \varepsilon}\|F\|_{L_{2}\left(D_{T}\right)}^{2}+c^{2}|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \operatorname{mes} D_{T}\right) \\
= & c^{4}\left(|\lambda| \frac{1+\alpha}{2}+1\right)\|F\|_{L_{2}\left(D_{T}\right)}^{2}+2 c^{2}|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \operatorname{mes} D_{T} . \tag{2.18}
\end{align*}
$$

From (2.18), in the case $0<\alpha<1$, follows inequality (2.14) with $c_{1}=$ $c^{2}\left(|\lambda| \frac{1+\alpha}{2}+1\right)^{\frac{1}{2}}$ and $c_{2}=c\left(2|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \text { mes } D_{T}\right)^{\frac{1}{2}}$, where $q=\frac{1}{1-\alpha}$. Lemma 2 is completely proved.

Remark 1. From the proof of Lemma 2 it follows that in estimate (2.14) the constants $c_{1}$ and $c_{2}$ are equal:

1) $\alpha>1, \quad \lambda<0: \quad c_{1}=c^{2}, \quad c_{2}=0$;
2) $0<\alpha<1, \quad-\infty<\lambda<+\infty$ :

$$
\begin{equation*}
c_{1}=c^{2}\left(|\lambda| \frac{1+\alpha}{2}+1\right)^{\frac{1}{2}}, \quad c_{2}=c\left(2|\lambda| \frac{1-\alpha}{2 \varepsilon^{q-1}} \operatorname{mes} D_{T}\right)^{\frac{1}{2}} \tag{2.20}
\end{equation*}
$$

where the constant $c=\left(1+\frac{e}{2} T^{2}+\frac{e}{2} T^{4}\right)^{\frac{1}{2}}$ is taken from estimate (2.4) and $q=\frac{2}{1-\alpha}$.

Remark 2. Below we will consider the linear problem corresponding to (1.1), (1.2), (1.3), when $\lambda=0$. In that case, for $F \in L_{2}\left(D_{T}\right)$ we analogously to the above introduce the concept of a weak generalized solution $u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$, when the integral equality

$$
\begin{equation*}
(u, \phi)_{\square}:=\int_{D_{T}} \square u \square \phi d x d t=\int_{D_{T}} F \phi d x d t \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{2.21}
\end{equation*}
$$

holds.
Remark 3. In view of (1.5) and (2.4), taking into account that

$$
\begin{aligned}
\left|(\square u, \square \phi)_{L_{2}\left(D_{T}\right)}\right| & =\left|\int_{D_{T}} \square u \square \phi d x d t\right| \leq\|\square u\|_{L_{2}\left(D_{T}\right)}\|\square \phi\|_{L_{2}\left(D_{T}\right)} \\
& \leq\|\square u\|_{W_{2, \square}^{1}\left(D_{T}\right)}\|\square \phi\|_{\dot{W}_{2, \square}^{1}\left(D_{T}\right)},
\end{aligned}
$$

the bilinear form

$$
(u, \phi)_{\square}:=\int_{D_{T}} \square u \square \phi d x d t
$$

of (2.21) can be considered as a scalar product in the Hilbert space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$. Therefore, since for $F \in L_{2}\left(D_{T}\right)$

$$
\left|\int_{D_{T}} F \phi d x d t\right| \leq\|F\|_{L_{2}\left(D_{T}\right)}\|\phi\|_{L_{2}\left(D_{T}\right)} \leq\|F\|_{L_{2}\left(D_{T}\right)}\|\phi\|_{\dot{W}_{2, \square}^{1}\left(D_{T}\right)},
$$

due to the theorem of Riesz [31, p. 83] there is a unique function $u$ in the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ that satisfies equality (2.21) for any $\phi \in \stackrel{\circ}{W}_{2, \square}\left(D_{T}\right)$ and for whose norm an estimate

$$
\begin{equation*}
\|u\|_{\stackrel{W}{W}_{2, \square}^{1}\left(D_{T}\right)} \leq\|F\|_{L_{2}\left(D_{T}\right)} \tag{2.22}
\end{equation*}
$$

is valid. Thus, being introduced the notation $u=L_{0}^{-1} F$, we find that to the linear problem corresponding to (1.1), (1.2), (1.3), when $\lambda=0$, there corresponds a linear bounded operator

$$
L_{0}^{-1}: L_{2}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right),
$$

for whose norm, by virtue of (2.22), an estimate

$$
\begin{equation*}
\left\|L_{0}^{-1}\right\|_{L_{2}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)} \leq\|F\|_{L_{2}\left(D_{T}\right)} \tag{2.23}
\end{equation*}
$$

is true.
Taking into account Definition 1 and Remark 3, equality (2.3), which is equivalent to problem (2.2), (1.2), (1.3), can be rewritten in the form of an equivalent equation

$$
\begin{equation*}
u=L_{0}^{-1}\left[\lambda|u|^{\alpha} \operatorname{sgn} u+F\right] \tag{2.24}
\end{equation*}
$$

in the Hilbert space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$.
Remark 4. The embedding operator $I:{ }_{W}^{W}{ }_{2}^{1}\left(D_{T}\right) \rightarrow L_{q}\left(D_{T}\right)$ is a linear continuous compact operator for $1<q<\frac{2(n+1)}{n-1}$ when $n \geq 2$ [29, p. 81]. At the same time, the operator of Nemytskii $N: L_{q}\left(D_{T}\right) \rightarrow L_{2}\left(D_{T}\right)$, which acts according to the formula $N u=\lambda|u|^{\alpha} \operatorname{sgn} u, \alpha>1$, is continuous and bounded for $q \geq 2 \alpha$ [32, p. 349], [33, pp. 66, 67]. Thus, if $1<\alpha<\frac{n+1}{n-1}$, then there exists a number $q$ such that $1<2 \alpha \leq q<\frac{2(n+1)}{n-1}$ and hence the operator

$$
\begin{equation*}
N_{1}=N I: \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}\right) \rightarrow L_{2}\left(D_{T}\right) \tag{2.25}
\end{equation*}
$$

is a continuous and compact operator. In that case, since $u \in{ }_{W}^{\circ}{ }_{2}^{1}\left(D_{T}\right)$, it is clear that $f(u)=|u|^{\alpha} \operatorname{sgn} u \in L_{2}\left(D_{T}\right)$. Further, since in view of (1.5) the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ is continuously embedded into the space ${ }_{W}^{W}{ }_{2}^{1}\left(D_{T}\right)$, taking into account (2.25) the operator

$$
\begin{equation*}
N_{2}=N I I_{1}: \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \rightarrow L_{2}\left(D_{T}\right), \tag{2.26}
\end{equation*}
$$

where $I_{1}: \stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W}{ }_{2}^{1}\left(D_{T}\right)$ is the embedding operator, is continuous and compact for $1<\alpha<\frac{n+1}{n-1}$. For $0<\alpha<1$, operator (2.26) is also continuous
and compact, since according to the theorem of Rellich [29, p. 64] the space $\stackrel{\circ}{W}_{2}^{1}\left(D_{T}\right)$ is continuously and compactly embedded into $L_{2}\left(D_{T}\right)$, and the space $L_{2}\left(D_{T}\right)$, in turn, is continuously embedded into $L_{p}\left(D_{T}\right)$ for $0<p<2$.

Let us rewrite equation (2.24) in the form

$$
\begin{equation*}
u=A u:=L_{0}^{-1}\left(N_{2} u+F\right), \tag{2.27}
\end{equation*}
$$

where the operator $N_{2}: \stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right) \rightarrow L_{2}\left(D_{T}\right)$, for $0<\alpha<\frac{n+1}{n-1}, \alpha \neq 1$, is continuous and compact in view of Remark 4. Then taking into account (2.23) the operator $A: \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \rightarrow \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ in (2.27) is also continuous and compact. At the same, time according to the a priori estimate (2.14) of Lemma 2, in which the constants $c_{1}$ and $c_{2}$ are given by equalities (2.19) and (2.20), for any parameter $\tau \in[0,1]$ and for any solution $u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ of the equation $u=\tau A u$ with this parameter we have the a priori estimation (2.14) with constants $c_{1}>0$ and $c_{2} \geq 0$ not depending on $u, \tau$ and $F$. Therefore, according to the theorem of Lere-Schauder [34, p. 375], equation (2.27) and, consequently, problem (2.2), (1.2), (1.3) have at least one weak generalized solution $u$ in the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$.

Thus the following statement is valid.
Theorem 1. Let $0<\alpha<\frac{n+1}{n-1}, \alpha \neq 1, \lambda \neq 0$ and in the case $\alpha>1$ additionally require that $\lambda<0$. Then for any $F \in L_{2}\left(D_{T}\right)$ problem (2.2), (1.2), (1.3) has at least one weak generalized solution $u \in \stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$.
3. the Uniqueness of a Solution of Problem (1.1), (1.2), (1.3) in the Case of A Nonlinearity of the Form $f(u)=|u|^{\alpha} \operatorname{sgn} u$

Let $F \in L_{2}\left(D_{T}\right)$, and $u_{1}, u_{2}$ be two weak generalized solutions of problem (2.2), (1.2), (1.3) in the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$, i.e., according to (2.3) the following equalities

$$
\begin{equation*}
\int_{D_{T}} \square u_{i} \square \phi d x d t=\lambda \int_{D_{T}} \phi\left|u_{i}\right|^{\alpha} \operatorname{sgn} u_{i} d x d t+\int_{D_{T}} F \phi d x d t \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{3.1}
\end{equation*}
$$

are fulfilled and $\left|u_{i}\right|^{\alpha} \in L_{2}\left(D_{T}\right), i=1,2$.
For the difference $v=u_{2}-u_{1}$, from (3.1) it follows that

$$
\begin{align*}
& \int_{D_{T}} \square v \square \phi d x d t \\
& \quad=\lambda \int_{D_{T}} \phi\left(\left|u_{2}\right|^{\alpha} \operatorname{sgn} u_{2}-\left|u_{1}\right|^{\alpha} \operatorname{sgn} u_{1}\right) d x d t \quad \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) . \tag{3.2}
\end{align*}
$$

Assuming $\phi=v \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ in equality (3.2), we obtain

$$
\begin{equation*}
\int_{D_{T}}(\square v)^{2} d x d t=\lambda \int_{D_{T}}\left(\left|u_{2}\right|^{\alpha} \operatorname{sgn} u_{2}-\left|u_{1}\right|^{\alpha} \operatorname{sgn} u_{1}\right)\left(u_{2}-u_{1}\right) d x d t \tag{3.3}
\end{equation*}
$$

Note that for the finite values of $u_{1}$ and $u_{2}$ with $\alpha>0$ the inequality

$$
\begin{equation*}
\left(\left|u_{2}\right|^{\alpha} \operatorname{sgn} u_{2}-\left|u_{1}\right|^{\alpha} \operatorname{sgn} u_{1}\right)\left(u_{2}-u_{1}\right) \geq 0 \tag{3.4}
\end{equation*}
$$

holds.
From (3.3) and inequality (3.4), which is true for almost all points $(x, t) \in D_{T}$ with $u_{i} \in \stackrel{\circ}{W} \underset{2, \square}{1}\left(D_{T}\right), i=1,2$, when $\alpha>0$ and $\lambda<0$, it follows that

$$
\int_{D_{T}}(\square v)^{2} d x d t \leq 0
$$

whence, due to (2.4), we obtain $v=0$, i.e. $u_{1}=u_{2}$.
Thus the following statement is valid.
Theorem 2. Let $\alpha>0, \alpha \neq 1$ and $\lambda<0$. Then for any $F \in L_{2}\left(D_{T}\right)$, problem (2.2), (1.2), (1.3) cannot have more than one generalized solution in the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$.

In turn, Theorems 1 and 2 give rise to
Theorem 3. Let $0<\alpha<\frac{n+1}{n-1}, \alpha \neq 1$ and $\lambda<0$. Then for any $F \in$ $L_{2}\left(D_{T}\right)$, problem (2.2), (1.2), (1.3) has a unique weak generalized solution $u \in$ $\stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$.
4. The Absence of a Solution of Problem (1.1), (1.2), (1.3) in the Case of a Nonlinearity of the Form $f(u)=|u|^{\alpha}$
Assume now in equation (1.1) and therefore in the integral equality (1.3) that $f(u)=|u|^{\alpha}, \alpha>1$.

Theorem 4. Let $F^{0} \in L_{2}\left(D_{T}\right),\left\|F^{0}\right\|_{L_{2}\left(D_{T}\right)} \neq 0, F^{0} \geq 0$, and $F=\mu F^{0}, \mu=$ const $>0$. Then in the case $f(u)=|u|^{\alpha}, \alpha>1$, with $\lambda>0$ there exists a number $\mu_{0}=\mu_{0}\left(F^{0}, \lambda, \alpha\right)>0$ such that for $\mu>\mu_{0}$ problem (1.1), (1.2), (1.3) cannot have a weak generalized solution in the space $\stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$.

Proof. Let us assume that when the conditions of the theorem are satisfied the solution $u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ of problem (1.1), (1.2), (1.3) exists for any fixed $\mu>0$. Then equality (1.6) takes the form

$$
\begin{equation*}
\int_{D_{T}} \square u \square \phi d x d t=\lambda \int_{D_{T}}|u|^{\alpha} \phi d x d t+\mu \int_{D_{T}} F^{0} \phi d x d t \quad \forall \phi \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right) \tag{4.1}
\end{equation*}
$$

It is easy to verify that

$$
\begin{equation*}
\int_{D_{T}} \square u \square \phi d x d t=\int_{D_{T}} u \square^{2} \phi d x d t \quad \forall \phi \in \stackrel{\circ}{C}^{4}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right), \tag{4.2}
\end{equation*}
$$

where $\stackrel{\circ}{C}^{4}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)=\left\{u \in C^{4}\left(\bar{D}_{T}\right):\left.u\right|_{\partial D_{T}}=0,\left.\frac{\partial u}{\partial t}\right|_{S_{T}^{0}}=0\right\} \subset \stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$. Indeed, since $u \in \stackrel{\circ}{W}_{2, \square}^{1}\left(D_{T}\right)$ and the space $\stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ is dense in $\stackrel{\circ}{W}{ }_{2, \square}^{1}\left(D_{T}\right)$, there exists a sequence $u_{k} \in \stackrel{\circ}{C}^{2}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ such that

$$
\begin{equation*}
\lim _{k \rightarrow \infty}\left\|u_{k}-u\right\|_{\dot{W}_{2, \square}^{1}\left(D_{T}\right)}=0 \tag{4.3}
\end{equation*}
$$

Taking into account that

$$
\begin{equation*}
\int_{D_{T}} \square u_{k} \square \phi d x d t=\int_{\partial D_{T}} \frac{\partial u_{k}}{\partial N} \square \phi d s-\int_{\partial D_{T}} u_{k} \frac{\partial}{\partial N} \square \phi d s+\int_{D_{T}} u_{k} \square^{2} \phi d x d t, \tag{4.4}
\end{equation*}
$$

where the derivative on the conormal $\frac{\partial}{\partial N}=\gamma_{n+1} \frac{\partial}{\partial t}-\sum_{i=1}^{n} \gamma_{i} \frac{\partial}{\partial x_{i}}$ is an inner differential operator on the characteristic manifold $S_{T}$ and therefore $\left.\frac{\partial u_{k}}{\partial N}\right|_{S_{T}}=0$ since $\left.u_{k}\right|_{S_{T}}=0$, from (4.4), due to the fact that $\left.u_{k}\right|_{S_{T}^{0}}=\left.\frac{\partial u_{k}}{\partial t}\right|_{S_{T}^{0}}=0$, and $\partial D_{T}=S_{T} \cup S_{T}^{0}$, we obtain

$$
\begin{equation*}
\int_{D_{T}} \square u_{k} \square \phi d x d t=\int_{D_{T}} u_{k} \square^{2} \phi d x d t \tag{4.5}
\end{equation*}
$$

Passing in (4.5) to the limit as $k \rightarrow \infty$, in view of (1.5) and (4.3) we obtain (4.2).
Taking into account (4.2), we rewrite equality (4.1) as

$$
\begin{align*}
& \lambda \int_{D_{T}}|u|^{\alpha} \phi d x d t=\int_{D_{T}} u \square^{2} \phi d x d t \\
& \quad-\mu \int_{D_{T}} F^{0} \phi d x d t \quad \forall \phi \in \stackrel{\circ}{C}^{4}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right) \tag{4.6}
\end{align*}
$$

Below we use the method of test functions [12, p. 10-12]. Let us select a test function $\phi \in \stackrel{\circ}{C}^{4}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right)$ such that $\left.\phi\right|_{D_{T}}>0$. If in Young's inequality with a parameter $\varepsilon>0$

$$
a b \leq \frac{\varepsilon}{\alpha} a^{\alpha}+\frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} b^{\alpha^{\prime}} ; \quad a, b \geq 0, \quad \alpha^{\prime}=\frac{\alpha}{\alpha-1}
$$

we take $a=|u| \phi^{\frac{1}{\alpha}}, b=\frac{\left|\square^{2} \phi\right|}{\phi^{\frac{1}{\alpha}}}$, then due to the fact that $\frac{\alpha^{\prime}}{\alpha}=\alpha^{\prime}-1$ we have

$$
\begin{equation*}
\left|u \square^{2} \phi\right|=|u| \phi^{\frac{1}{\alpha}} \frac{\left|\square^{2} \phi\right|}{\phi^{\frac{1}{\alpha}}} \leq \frac{\varepsilon}{\alpha}|u|^{\alpha} \phi+\frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \frac{\left|\square^{2} \phi\right|^{\alpha^{\prime}}}{\phi^{\alpha^{\prime}-1}} . \tag{4.7}
\end{equation*}
$$

By virtue of (4.7) and (4.6) we obtain the inequality

$$
\left(\lambda-\frac{\varepsilon}{\alpha}\right) \int_{D_{T}}|u|^{\alpha} \phi d x d t \leq \frac{1}{\alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \int_{D_{T}} \frac{\left|\square^{2} \phi\right|^{\alpha^{\prime}}}{\phi^{\alpha^{\prime}-1}} d x d t-\mu \int_{D_{T}} F^{0} \phi d x d t
$$

which, for $\varepsilon<\lambda \alpha$, implies

$$
\begin{equation*}
\int_{D_{T}}|u|^{\alpha} \phi d x d t \leq \frac{\alpha}{(\lambda \alpha-\varepsilon) \alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}} \int_{D_{T}} \frac{\left|\square^{2} \phi\right|^{\alpha^{\prime}}}{\phi^{\alpha^{\prime}-1}} d x d t-\frac{\alpha \mu}{\lambda \alpha-\varepsilon} \int_{D_{T}} F^{0} \phi d x d t \tag{4.8}
\end{equation*}
$$

Taking into account the equalities $\alpha^{\prime}=\frac{\alpha}{\alpha-1}, \alpha=\frac{\alpha^{\prime}}{\alpha^{\prime}-1}$ and $\min _{0<\varepsilon<\lambda \alpha} \frac{\alpha}{(\lambda \alpha-\varepsilon) \alpha^{\prime} \varepsilon^{\alpha^{\prime}-1}}=$ $\frac{1}{\lambda^{\alpha^{\prime}}}$ which is obtained at $\varepsilon=\lambda$, it follows from (4.8) that

$$
\begin{equation*}
\int_{D_{T}}|u|^{\alpha} \phi d x d t \leq \frac{1}{\lambda^{\alpha^{\prime}}} \int_{D_{T}} \frac{\left|\square^{2} \phi\right|^{\alpha^{\prime}}}{\phi^{\alpha^{\prime}-1}} d x d t-\frac{\alpha^{\prime} \mu}{\lambda} \int_{D_{T}} F^{0} \phi d x d t . \tag{4.9}
\end{equation*}
$$

Note that the existence of a test function $\phi$ such that

$$
\begin{equation*}
\phi \in \stackrel{\circ}{C}^{4}\left(\bar{D}_{T}, \partial D_{T}, S_{T}^{0}\right),\left.\phi\right|_{D_{T}}>0, \kappa=\int_{D_{T}} \frac{\left|\square^{2} \phi\right|^{\alpha^{\prime}}}{\phi^{\alpha^{\prime}-1}} d x d t<+\infty \tag{4.10}
\end{equation*}
$$

is not difficult to show. Indeed, it is easy to verify that the function

$$
\phi(x, t)=\left[\left(t^{2}-|x|^{2}\right)\left((T-t)^{2}-|x|^{2}\right)\right]^{m}
$$

satisfies conditions (4.10) for a sufficiently large positive $m$.
Since, by assumption, $F^{0} \in L_{2}\left(D_{T}\right),\left\|F^{0}\right\|_{L_{2}\left(D_{T}\right)} \neq 0, F^{0} \geq 0$, and mes $D_{T}<$ $+\infty$, due to the fact that $\left.\phi\right|_{D_{T}}>0$ we have

$$
\begin{equation*}
0<\kappa_{1}=\int_{D_{T}} F^{0} \phi d x d t<+\infty \tag{4.11}
\end{equation*}
$$

Let us denote by $g(\mu)$ the right side of inequality (4.9) which is a linear function with respect to $\mu$, then in view of (4.10) and (4.11)

$$
\begin{equation*}
g(\mu)<0 \text { for } \mu>\mu_{0} \text { and } g(\mu)>0 \text { for } \mu<\mu_{0} \tag{4.12}
\end{equation*}
$$

where

$$
g(\mu)=\frac{\kappa_{0}}{\lambda^{\alpha^{\prime}}}-\frac{\alpha^{\prime} \mu}{\lambda} \kappa_{1}, \quad \mu_{0}=\frac{\lambda}{\alpha^{\prime} \lambda^{\alpha^{\prime}}} \frac{\kappa_{0}}{\kappa_{1}}>0 .
$$

According to (4.12) with $\mu>\mu_{0}$, the right side of inequality (4.9) is negative, while the left side of this inequality is non-negative. The obtained contradiction proves Theorem 4.

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