

# The Cauchy–Goursat Problem for Wave Equations with Nonlinear Dissipative Term

S. S. Kharibegashvili\* and O. M. Jokhadze\*\*

*Tbilisi State University, Tbilisi, Georgia*

Received February 28, 2011

**Abstract**—The Cauchy–Goursat problem for wave equations with nonlinear dissipative term is studied. The existence, uniqueness, and blow-up of global solutions of this problem are considered. The local solvability of this problem is also discussed.

**DOI:** 10.1134/S000143461311028X

**Keywords:** *wave equation, Cauchy–Goursat problem, Volterra-type integral equation, Green’s formula, dissipative term, blow-up of a solution.*

## 1. STATEMENT OF THE PROBLEM

In the plane of independent variables  $x$  and  $t$ , let us consider the following wave equation with nonlinear dissipative term [1, p. 57 (Russian transl.)], [2]

$$Lu := u_{tt} - u_{xx} + g(x, t, u)u_t = f(x, t), \quad (1.1)$$

where  $f$  and  $g$  are given functions, and  $u$  is an unknown real function.

Let  $D_T := \{(x, t) : 0 < x < t, 0 < t < T\}$  denote the triangular domain bounded by the characteristic closed interval  $\gamma_{1,T}: x = t, 0 \leq t \leq T$ , as well as by the closed intervals  $\gamma_{2,T}: x = 0, 0 \leq t \leq T$  and  $\gamma_{3,T}: t = T, 0 \leq x \leq T$ . For Eq. (1.1), consider the Cauchy–Goursat problem of finding the solution  $u(x, t)$  in the domain  $D_T$ , subject to the conditions [3, p. 284]

$$u_x|_{\gamma_{2,T}} = 0, \quad u|_{\gamma_{1,T}} = 0. \quad (1.2)$$

Note that, for hyperbolic-type nonlinear equations, the questions of the existence, uniqueness, and blow-up of global solutions of initial, mixed, nonlocal, and other problems were studied in numerous papers (see, for example, [4]–[18]). It is well known that, in the linear case, i.e., for  $g(x, t, u) = g(x, t)$ , problem (1.1), (1.2), is well posed and its global solvability was established in the corresponding function spaces (see, for example, [1], [19]–[23]).

In what follows, it will be shown that, under certain conditions on the nonlinear function  $g(x, t, u)$ , problem (1.1), (1.2) is locally solvable; we shall also obtain conditions for the global solvability whose violation, in general, may lead to the blow-up of the solution in finite time.

**Definition 1.1.** Let  $f \in C(\overline{D}_T)$ , and let  $g \in C(\overline{D}_T \times \mathbb{R})$ , where  $\mathbb{R} := (-\infty, +\infty)$ . A function  $u$  is called a *strong generalized solution of Problem (1.1), (1.2)* of class  $C^1$  in the domain  $D_T$  if  $u \in C^1(\overline{D}_T)$  and there exists a sequence of functions  $u_n \in \dot{C}^2(\overline{D}_T, \Gamma_T)$  such that  $u_n \rightarrow u$  and  $Lu_n \rightarrow f$  as  $n \rightarrow \infty$  in the spaces  $C^1(\overline{D}_T)$  and  $C(\overline{D}_T)$ , respectively, where

$$\dot{C}^2(\overline{D}_T, \Gamma_T) := \{v \in C^2(\overline{D}_T) : v_x|_{\gamma_{2,T}} = 0, v|_{\gamma_{1,T}} = 0\}, \quad \Gamma_T := \gamma_{1,T} \cup \gamma_{2,T}.$$

\*E-mail: kharibegashvili@yahoo.com

\*\*E-mail: ojokhadze@yahoo.com

**Remark 1.1.** Obviously, the classical solution of problem (1.1), (1.2) from the space  $u \in \mathring{C}^2(\overline{D}_T, \Gamma_T)$  is a strong generalized solution of this problem of class  $C^1$  in the domain  $D_T$ . In turn, if the strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$  belongs to the space  $C^2(\overline{D}_T)$ , then it will also be the classical solution of this problem.

**Definition 1.2.** Let  $f \in C(\overline{D}_\infty)$ , and let  $g \in C(\overline{D}_\infty \times \mathbb{R})$ . We say that problem (1.1), (1.2) is *globally solvable for the class  $C^1$*  if, for any finite  $T > 0$ , this problem has a strong generalized solution of class  $C^1$  in the domain  $D_T$ .

2. A PRIORI ESTIMATES OF THE SOLUTION OF PROBLEM (1.1), (1.2) FOR THE CLASSES  $C(\overline{D}_T)$ ,  $C^1(\overline{D}_T)$

**Lemma 2.1.** Let  $f \in C(\overline{D}_T)$ , let  $g \in C(\overline{D}_T \times \mathbb{R})$ , and let

$$g(x, t, s) \geq -M_T, \quad (x, t, s) \in \overline{D}_T \times \mathbb{R}, \quad M_T := \text{const} > 0. \tag{2.1}$$

Then, in the domain  $D_T$ , the strong generalized solution of problem (1.1), (1.2) of class  $C^1$  satisfies the a priori estimate

$$\|u\|_{C(\overline{D}_T)} \leq c_0 \|f\|_{C(\overline{D}_T)}, \tag{2.2}$$

where  $c_0 = c_0(T, M_T)$  is a positive constant independent of  $u$  and  $f$ .

**Proof.** Let  $u$  be a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ . Then, in view of Definition 1.1, there exists a sequence of functions  $u_n \in \mathring{C}^2(\overline{D}_T, \Gamma_T)$  such that

$$\lim_{n \rightarrow \infty} \|u_n - u\|_{C^1(\overline{D}_T)} = 0, \quad \lim_{n \rightarrow \infty} \|Lu_n - f\|_{C(\overline{D}_T)} = 0, \tag{2.3}$$

and, therefore, also

$$\lim_{n \rightarrow \infty} \|g(x, t, u_n)u_{nt} - g(x, t, u)u_t\|_{C(\overline{D}_T)} = 0. \tag{2.4}$$

Consider the function  $u_n \in \mathring{C}^2(\overline{D}_T, \Gamma_T)$  as the solution of the following problem:

$$Lu_n = f_n, \tag{2.5}$$

$$\left. \frac{\partial u_n}{\partial x} \right|_{\gamma_{2,T}} = 0, \quad u_n|_{\gamma_{1,T}} = 0. \tag{2.6}$$

Here

$$f_n := Lu_n. \tag{2.7}$$

Multiplying both sides of relation (2.5) by  $\partial u_n / \partial t$  and integrating the resulting equality over the domain  $D_\tau := \{(x, t) \in D_T : 0 < t < \tau\}$ ,  $0 < \tau \leq T$ , we obtain

$$\begin{aligned} & \frac{1}{2} \int_{D_\tau} \frac{\partial}{\partial t} \left( \frac{\partial u_n}{\partial t} \right)^2 dx dt - \int_{D_\tau} \frac{\partial^2 u_n}{\partial x^2} \frac{\partial u_n}{\partial t} dx dt + \int_{D_\tau} g(x, t, u_n) \left( \frac{\partial u_n}{\partial t} \right)^2 dx dt \\ & = \int_{D_\tau} f_n \frac{\partial u_n}{\partial t} dx dt. \end{aligned}$$

Set  $\Omega_\tau := \overline{D}_\infty \cap \{t = \tau\}$ ,  $0 < \tau \leq T$ . Then, in view of (2.6), applying Green's formula to the left-hand side of the last equality, we can write

$$\begin{aligned} \int_{D_\tau} f_n \frac{\partial u_n}{\partial t} dx dt &= \int_{\gamma_{1,\tau}} \frac{1}{2\nu_t} \left[ \left( \frac{\partial u_n}{\partial x} \nu_t - \frac{\partial u_n}{\partial t} \nu_x \right)^2 + \left( \frac{\partial u_n}{\partial t} \right)^2 (\nu_t^2 - \nu_x^2) \right] ds \\ &+ \frac{1}{2} \int_{\Omega_\tau} \left[ \left( \frac{\partial u_n}{\partial x} \right)^2 + \left( \frac{\partial u_n}{\partial t} \right)^2 \right] dx + \int_{D_\tau} g(x, t, u_n) \left( \frac{\partial u_n}{\partial t} \right)^2 dx dt, \end{aligned} \tag{2.8}$$

where  $\nu := (\nu_x, \nu_t)$  is the unit vector of the outer/exterior normal to  $\partial D_\tau$  and  $\gamma_{1,\tau} := \gamma_{1,T} \cap \{t \leq \tau\}$ . Since the operator  $\nu_t (\partial/\partial x) - \nu_x (\partial/\partial t)$  is the inner differential operator on  $\gamma_{1,T}$ , in view of the second condition from (2.6), it follows that

$$\left( \frac{\partial u_n}{\partial x} \nu_t - \frac{\partial u_n}{\partial t} \nu_x \right) \Big|_{\gamma_{1,\tau}} = 0. \tag{2.9}$$

Further, it is easy to see that

$$(\nu_t^2 - \nu_x^2)|_{\gamma_{1,\tau}} = 0. \tag{2.10}$$

Therefore, from (2.8)–(2.10) we obtain

$$w_n(\tau) := \int_{\Omega_\tau} \left[ \left( \frac{\partial u_n}{\partial x} \right)^2 + \left( \frac{\partial u_n}{\partial t} \right)^2 \right] dx \leq 2 \int_{D_\tau} f_n \frac{\partial u_n}{\partial t} dx dt + 2M_T \int_{D_\tau} \left( \frac{\partial u_n}{\partial t} \right)^2 dx dt. \tag{2.11}$$

Taking into account the inequality

$$2f_n \frac{\partial u_n}{\partial t} \leq \left( \frac{\partial u_n}{\partial t} \right)^2 + f_n^2$$

and using (2.11), we obtain

$$w_n(\tau) \leq (1 + 2M_T) \int_{D_\tau} \left( \frac{\partial u_n}{\partial t} \right)^2 dx dt + \int_{D_\tau} f_n^2 dx dt.$$

Hence, using the expression for the function  $w_n(\tau)$ , we can write

$$w_n(\tau) \leq m_T \int_0^\tau w_n(\sigma) d\sigma + \|f_n\|_{L_2(D_\tau)}^2,$$

where  $m_T := \max(1 + 2M_T, 1)$ . Hence, since  $\|f_n\|_{L_2(D_\tau)}^2$  is a nondecreasing function of  $\tau$ , by Gronwall’s lemma [24, p. 13 (Russian transl.)], we have

$$w_n(\tau) \leq \exp(m_T \tau) \|f_n\|_{L_2(D_\tau)}^2. \tag{2.12}$$

If  $(x, t) \in \overline{D}_T$ , then, in view of the second condition from (2.6), the following equality holds:

$$u_n(x, t) = u_n(x, t) - u_n(t, t) = \int_t^x \frac{\partial u_n(\sigma, t)}{\partial x} d\sigma;$$

hence, in view of (2.12) we have

$$\begin{aligned} |u_n(x, t)|^2 &\leq \int_x^t d\sigma \int_x^t \left[ \frac{\partial u_n(\sigma, t)}{\partial x} \right]^2 d\sigma \leq (t - x) \int_{\Omega_t} \left[ \frac{\partial u_n(\sigma, t)}{\partial x} \right]^2 d\sigma \\ &\leq (t - x) w_n(t) \leq t w_n(t) \leq T \exp(m_T T) \|f_n\|_{C(\overline{D}_T)}^2 \text{mes } D_T \\ &= 2^{-1} T^3 \exp(m_T T) \|f_n\|_{C(\overline{D}_T)}^2. \end{aligned} \tag{2.13}$$

Using (2.13), we obtain

$$\|u_n\|_{C(\overline{D}_T)} \leq T \sqrt{\frac{T}{2}} \exp\left(\frac{m_T T}{2}\right) \|f_n\|_{C(\overline{D}_T)}.$$

Passing to the limit as  $n \rightarrow \infty$  in this inequality and taking into account (2.3), (2.7), we can write

$$\|u\|_{C(\overline{D}_T)} \leq T \sqrt{\frac{T}{2}} \exp\left(\frac{m_T T}{2}\right) \|f\|_{C(\overline{D}_T)}. \tag{2.14}$$

Estimate (2.2) is proved. □

**Remark 2.1.** It follows from (2.14) that the constant  $c_0$  in estimate (2.2) can be taken to be

$$c_0 := T\sqrt{\frac{T}{2}} \exp\left(\frac{m_T T}{2}\right). \tag{2.15}$$

Below, using the classical method of characteristics and taking into account estimate (2.2), we obtain an a priori estimate in the space  $C^1(\overline{D}_T)$  for the strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ .

The following statement is valid.

**Lemma 2.2.** *Under the assumptions of Lemma 2.1, the strong generalized solution of problem (1.1), (1.2) of class  $C^1$  satisfies the following a priori estimate in the domain  $D_T$ :*

$$\|u\|_{C^1(\overline{D}_T)} \leq c_1 \tag{2.16}$$

with positive constant  $c_1 = c_1(T, c_0, \|f\|_{C(\overline{D}_T)})$ , where

$$\|u\|_{C^1(\overline{D}_T)} := \max\{\|u\|_{C(\overline{D}_T)}, \|u_x\|_{C(\overline{D}_T)}, \|u_t\|_{C(\overline{D}_T)}\}.$$

**Proof.** Let  $u$  be a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ . Then we have the limit equalities (2.3), (2.4), where  $u_n$  can be regarded as the solution of problem (2.5), (2.6) with the right-hand side of  $f_n$  from (2.7). For a fixed natural number  $n$ , we introduce the following functions:

$$u_{n1} := u_{nt} - u_{nx}, \quad u_{n2} := u_{nt} + u_{nx}, \quad u_{n3} := u_n, \tag{2.17}$$

which, in view of (2.2), for  $0 \leq t \leq T$ , satisfy the boundary conditions

$$u_{n1}(0, t) = u_{n2}(0, t), \quad u_{n2}(t, 0) = 0, \quad u_{n3}(t, 0) = 0. \tag{2.18}$$

In view of (1.1) and (2.17), the unknown functions  $u_{n1}, u_{n2}, u_{n3}$  satisfy the following system of first-order partial differential equations:

$$\begin{cases} \frac{\partial u_{n1}}{\partial t} + \frac{\partial u_{n1}}{\partial x} = f_n(x, t) - \frac{1}{2}g(x, t, u_{n3})(u_{n1} + u_{n2}), \\ \frac{\partial u_{n2}}{\partial t} - \frac{\partial u_{n2}}{\partial x} = f_n(x, t) - \frac{1}{2}g(x, t, u_{n3})(u_{n1} + u_{n2}), \\ \frac{\partial u_{n3}}{\partial t} - \frac{\partial u_{n3}}{\partial x} = u_{n1}. \end{cases} \tag{2.19}$$

Integrating the equations of system (2.19) along the corresponding characteristic curves and taking into account the boundary conditions (2.18), we obtain

$$\begin{cases} u_{n1}(x, t) - u_{n1}(0, t) = \int_{t-x}^t \left[ f_n(P_\tau) - \frac{1}{2}g(P_\tau, u_{n3}(P_\tau))(u_{n1}(P_\tau) + u_{n2}(P_\tau)) \right] d\tau, \\ u_{n2}(x, t) = \int_{(x+t)/2}^t \left[ f_n(Q_\tau) - \frac{1}{2}g(Q_\tau, u_{n3}(Q_\tau))(u_{n1}(Q_\tau) + u_{n2}(Q_\tau)) \right] d\tau, \\ u_{n3}(x, t) = \int_{(x+t)/2}^t u_{n1}(Q_\tau) d\tau, \end{cases}$$

where  $P_\tau := (x - t + \tau, \tau)$  and  $Q_\tau := (x + t - \tau, \tau)$ . Using the second equation of this system and the

first equality from (2.18), as well as the notation  $P_{\tau_0} := (t - \tau, \tau)$ , we can rewrite this system as

$$\begin{cases} u_{n1}(x, t) = -\frac{1}{2} \int_{t-x}^t [g(P_{\tau}, u_{n3}(P_{\tau}))(u_{n1}(P_{\tau}) + u_{n2}(P_{\tau}))] d\tau \\ \quad - \frac{1}{2} \int_{t/2}^t [g(P_{\tau_0}, u_{n3}(P_{\tau_0}))(u_{n1}(P_{\tau_0}) + u_{n2}(P_{\tau_0}))] d\tau + F_{n1}(x, t), \\ u_{n2}(x, t) = -\frac{1}{2} \int_{(x+t)/2}^t g(Q_{\tau}, u_{n3}(Q_{\tau}))(u_{n1}(Q_{\tau}) + u_{n2}(Q_{\tau})) d\tau + F_{n2}(x, t), \\ u_{n3}(x, t) = \int_{(x+t)/2}^t u_{n1}(Q_{\tau}) d\tau. \end{cases} \tag{2.20}$$

Here

$$F_{n1}(x, t) := \int_{t-x}^t f_n(P_{\tau}) d\tau + \int_{t/2}^t f_n(P_{\tau_0}) d\tau, \quad F_{n2}(x, t) := \int_{(x+t)/2}^t f_n(Q_{\tau}) d\tau. \tag{2.21}$$

Passing in relations (2.20), (2.21) to the limit as  $n \rightarrow \infty$  in the space  $C(\overline{D}_T)$  and taking into account relations (2.3), (2.4), (2.7), and (2.17), we obtain

$$\begin{cases} u_1(x, t) = -\frac{1}{2} \int_{t-x}^t [g(P_{\tau}, u_3(P_{\tau}))(u_1(P_{\tau}) + u_2(P_{\tau}))] d\tau \\ \quad - \frac{1}{2} \int_{t/2}^t [g(P_{\tau_0}, u_3(P_{\tau_0}))(u_1(P_{\tau_0}) + u_2(P_{\tau_0}))] d\tau + F_1(x, t), \\ u_2(x, t) = -\frac{1}{2} \int_{(x+t)/2}^t g(Q_{\tau}, u_3(Q_{\tau}))(u_1(Q_{\tau}) + u_2(Q_{\tau})) d\tau + F_2(x, t), \\ u_3(x, t) = \int_{(x+t)/2}^t u_1(Q_{\tau}) d\tau, \end{cases} \tag{2.22}$$

where  $u_i := \lim_{n \rightarrow \infty} u_{ni}$  (in the norm of the space  $C(\overline{D}_T)$ )  $i = 1, 2, 3$ , and

$$F_1(x, t) := \int_{t-x}^t f(P_{\tau}) d\tau + \int_{t/2}^t f(P_{\tau_0}) d\tau, \quad F_2(x, t) := \int_{(x+t)/2}^t f(Q_{\tau}) d\tau. \tag{2.23}$$

Obviously,  $u_3 = u$  is a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ . Further,

$$u_1 := u_t - u_x, \quad u_2 := u_t + u_x. \tag{2.24}$$

Let  $G_T := \{(x, t, s) \in \mathbb{R}^3 : (x, t) \in \overline{D}_T, |s| \leq c_0 \|f\|_{C(\overline{D}_T)}\}$ , and let

$$K := \sup_{(x,t,s) \in G_T} |g(x, t, s)| < +\infty, \tag{2.25}$$

where  $K = K(T, c_0, \|f\|_{C(\overline{D}_T)})$ . Then in view of the a priori estimate (2.2), the strong generalized solution  $u_3 = u$  of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$  satisfies the estimate

$$|g(x, t, u_3(x, t))| \leq K, \quad (x, t) \in \overline{D}_T. \tag{2.26}$$

Let

$$v_i(t) := \sup_{(\xi, \tau) \in \overline{D}_t} |u_i(\xi, \tau)|, \quad i = 1, 2, 3, \quad F(t) := \sup_{(\xi, \tau) \in \overline{D}_t} |f(\xi, \tau)|. \tag{2.27}$$

In view of (2.23), (2.26), and (2.27), relations (2.22) imply

$$|u_1(x, t)| \leq K \int_0^t (v_1(\tau) + v_2(\tau)) d\tau + 2tF(t),$$

$$|u_2(x, t)| \leq \frac{K}{2} \int_0^t (v_1(\tau) + v_2(\tau)) d\tau + tF(t),$$

$$|u_3(x, t)| \leq \int_0^t v_1(\tau) d\tau.$$

Hence, for  $(\xi, \tau) \in \overline{D}_t$ , we have

$$|u_1(\xi, \tau)| \leq K \int_0^\tau (v_1(\tau_1) + v_2(\tau_1)) d\tau_1 + 2\tau F(\tau),$$

$$|u_2(\xi, \tau)| \leq \frac{K}{2} \int_0^\tau (v_1(\tau_1) + v_2(\tau_1)) d\tau_1 + \tau F(\tau),$$

$$|u_3(\xi, \tau)| \leq \int_0^\tau v_1(\tau_1) d\tau_1$$

and, therefore, in view of (2.27) and the fact that  $tF(t)$  is a nondecreasing function, we can write

$$v_1(t) \leq K \int_0^t (v_1(\tau) + v_2(\tau)) d\tau + 2tF(t),$$

$$v_2(t) \leq \frac{K}{2} \int_0^t (v_1(\tau) + v_2(\tau)) d\tau + tF(t),$$

$$v_3(t) \leq \int_0^t v_1(\tau) d\tau.$$

Setting  $v(t) := \max_{1 \leq i \leq 3} v_i(t)$ , from the above inequalities we obtain

$$v(t) \leq 2K \int_0^t v(\tau) d\tau + 2tF(t),$$

whence, applying Gronwall's lemma, we have

$$v(t) \leq 2tF(t) \exp(2tK) \leq 2T \exp(2TK) \|f\|_{C(\overline{D}_T)}, \quad 0 \leq t \leq T.$$

Now it readily follows from (2.24) that

$$\|u\|_{C^1(\overline{D}_T)} \leq \|v\|_{C[0,T]} \leq 2T \exp(2TK) \|f\|_{C(\overline{D}_T)}.$$

Lemma 2.2 is proved and

$$c_1 := 2T \exp(2TK) \|f\|_{C(\overline{D}_T)}, \tag{2.28}$$

where  $K$  is given by (2.25). □

### 3. EQUIVALENCE OF PROBLEM (1.1), (1.2) TO A SYSTEM OF VOLTERRA-TYPE NONLINEAR INTEGRAL EQUATIONS AND ITS LOCAL SOLVABILITY

First, let us show that problem (2.5), (2.6) is equivalent to problem (2.19), (2.18) in the classical sense. Indeed, if  $u_n \in C^2$  is a solution of problem (2.5), (2.6), then the system of functions  $u_{n1}$ ,  $u_{n2}$ , and  $u_{n3}$  will, obviously, be a solution of problem (2.19), (2.18). Conversely, let  $u_{n1}, u_{n2}, u_{n3} \in C^1$  be solutions of problem (2.19), (2.18). Let us show that  $u_n := u_{n3} \in C^2$  is a solution of problem (2.5), (2.6) and satisfies relations (2.17). If we show that  $u_{n2} = u_{nt} + u_{nx}$ , then, obviously, we have the equalities

$$u_{nt} = \frac{u_{n2} + u_{n1}}{2} \quad \text{and} \quad u_{nx} = \frac{u_{n2} - u_{n1}}{2},$$

whence it immediately follows that  $u_n \in C^2$  is a solution of problem (2.5), (2.6) in the classical sense.

Indeed, it follows from the first and second equations of system (2.19) that

$$\frac{\partial u_{n1}}{\partial t} + \frac{\partial u_{n1}}{\partial x} = \frac{\partial u_{n2}}{\partial t} - \frac{\partial u_{n2}}{\partial x}. \tag{3.1}$$

Further, because  $u_{n1} \in C^1$ , it follows from the third equation of system (2.19) that

$$\frac{\partial}{\partial t} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_{n3} \in C \quad \text{and} \quad \frac{\partial}{\partial x} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_{n3} \in C.$$

Hence, taking into account the fact that the first-order differential operators with constant coefficients are interchangeable, we obtain

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_{n3} &= \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) \frac{\partial}{\partial t} u_{n3} \in C, \\ \frac{\partial}{\partial x} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_{n3} &= \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) \frac{\partial}{\partial x} u_{n3} \in C. \end{aligned}$$

In view of these equalities, (3.1), and the third equality of system (2.19), we can write

$$\begin{aligned} &\left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) (u_{n2} - u_{nt} - u_{nx}) \\ &= \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_{n2} - \frac{\partial}{\partial t} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_n - \frac{\partial}{\partial x} \left( \frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u_n \\ &= \frac{\partial u_{n2}}{\partial t} - \frac{\partial u_{n2}}{\partial x} - \frac{\partial u_{n1}}{\partial t} - \frac{\partial u_{n1}}{\partial x} = 0. \end{aligned}$$

Therefore, in view of the second and the third equality from (2.18), we find that  $u_{n2} = u_{nt} + u_{nx}$ . This establishes the equivalence of problems (2.5), (2.6) and (2.19), (2.18) in the classical sense.

Above we have reduced problem (1.1), (1.2) to the system of Volterra-type nonlinear integral equations (2.22). Before considering the local solvability of problem (1.1), (1.2), let us make the following remark immediately following from the arguments given in Sec. 2.

**Remark 3.1.** Let  $u$  be a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ ; then  $u_1 := u_t - u_x$ ,  $u_2 := u_t + u_x$ ,  $u_3 := u$  is a continuous solution of the system of Volterra-type nonlinear integral equations (2.22). Conversely, if  $u_1, u_2, u_3$  is a continuous solution of system (2.22), then  $u := u_3$  is a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ , and the relations  $u_1 := u_t - u_x$ ,  $u_2 := u_t + u_x$  hold.

Now let us turn our attention to the proof of the local solvability of the system of Volterra-type nonlinear integral equations (2.22).

Let

$$f \in C(\overline{D}_\infty), \quad f_\infty := \sup_{(x,t) \in \overline{D}_\infty} |f(x,t)| < +\infty, \quad g \in C(\overline{D}_\infty \times \mathbb{R}) \tag{3.2}$$

and, for  $(x,t) \in \overline{D}_\infty$  and  $s, s_1, s_2$  such that  $|s|, |s_1|, |s_2| \leq R$ , we have

$$|g(x,t,s)| \leq M(R), \quad |g(x,t,s_2) - g(x,t,s_1)| \leq c(R)|s_2 - s_1|, \tag{3.3}$$

where  $M(R)$  and  $c(R)$  are nonnegative continuous functions of the argument  $R \geq 0$ .

**Theorem 3.1.** *Let the functions  $f$  and  $g$  satisfy conditions (3.2), (3.3). Then there exists a positive number  $T_* := T_*(f, g)$  such that, for  $T \leq T_*$ , problem (1.1), (1.2) has at least one strong generalized solution  $u$  of class  $C^1$  in the domain  $D_T$ .*

**Proof.** By Remark 3.1, problem (1.1), (1.2) in the space  $C^1(\overline{D}_T)$  is equivalent to the system of Volterra-type nonlinear integral equations (2.22) for the class  $C(\overline{D}_T)$ . To prove the unique solvability of system (2.22), we use the contraction mapping principle [26, p. 390].

Set  $U := (u_1, u_2, u_3)$ . Let us introduce the vector operator  $\Phi := (\Phi_1, \Phi_2, \Phi_3)$ , by the formula

$$\left\{ \begin{aligned} (\Phi_1 U)(x, t) &= -\frac{1}{2} \int_{t-x}^t g(P_\tau, u_3(P_\tau))(u_1(P_\tau) + u_2(P_\tau)) d\tau \\ &\quad + \frac{1}{2} \int_{t/2}^t g(P_{\tau_0}, u_3(P_{\tau_0}))(u_1(P_{\tau_0}) + u_2(P_{\tau_0})) d\tau + F_1(x, t), \\ (\Phi_2 U)(x, t) &= -\frac{1}{2} \int_{(x+t)/2}^t g(Q_\tau, u_3(Q_\tau))(u_1(Q_\tau) + u_2(Q_\tau)) d\tau + F_2(x, t), \\ (\Phi_3 U)(x, t) &= \int_{(x+t)/2}^t u_1(Q_\tau) d\tau. \end{aligned} \right. \tag{3.4}$$

Then system (2.22) can be rewritten in vector form:

$$U = \Phi U. \tag{3.5}$$

Let

$$\|U\|_{X_T} := \max_{1 \leq i \leq 3} \{ \|u_i\|_{C(\overline{D}_T)} \}, \quad U \in X_T := C(\overline{D}_T; \mathbb{R}^3),$$

where  $C(\overline{D}_T; \mathbb{R}^3)$  is the set of continuous vector functions  $U: \overline{D}_T \rightarrow \mathbb{R}^3$ . Let

$$B_R := \{U \in X_T : \|U\|_{X_T} \leq R\}$$

denote the closed ball of radius  $R > 0$  in the Banach space  $X_T$  centered at the zero element.

Below we shall prove that

- (1)  $\Phi$  maps the ball  $B_R$  into itself;
- (2)  $\Phi$  is a contraction mapping on  $B_R$ .

Indeed, in view of the first inequality (3.3) from (3.4), for  $U$  such that  $\|U\|_{X_T} \leq R$ , we have

$$\begin{aligned} |(\Phi_1 U)(x, t)| &\leq 2T(RM(R) + \|f\|_{C(\overline{D}_T)}), \\ |(\Phi_2 U)(x, t)| &\leq T(RM(R) + \|f\|_{C(\overline{D}_T)}), \quad |(\Phi_3 U)(x, t)| \leq TR. \end{aligned}$$

It follows from these estimates that

$$\|\Phi U\|_{X_T} \leq 2T(RM(R) + R + \|f\|_{C(\overline{D}_T)}) \leq 2T(RM(R) + R + f_\infty),$$

where  $f_\infty$  is defined in (3.2).

For a fixed  $R > 0$ , let the number  $T$  be small so that

$$2T(RM(R) + R + f_\infty) \leq R, \tag{3.6}$$

i.e.,  $\Phi U \in B_R$ , and thus condition (1) holds.

Further, in view of (3.3) from (3.4), for  $U^i$  such that  $\|U^i\|_{X_T} \leq R$ ,  $i = 1, 2$ , we have

$$\begin{aligned} &|(\Phi_1 U^2 - \Phi_1 U^1)(x, t)| \\ &\leq \frac{1}{2} \int_{t-x}^t (|g(P_\tau, u_3^2(P_\tau)) - g(P_\tau, u_3^1(P_\tau))| |u_1^2(P_\tau) + u_2^2(P_\tau)| \\ &\quad + |g(P_\tau, u_3^1(P_\tau))| |u_1^2(P_\tau) - u_1^1(P_\tau) + u_2^2(P_\tau) - u_2^1(P_\tau)|) d\tau \\ &\quad + \frac{1}{2} \int_{t/2}^t (|g(P_{\tau_0}, u_3^2(P_{\tau_0})) - g(P_{\tau_0}, u_3^1(P_{\tau_0}))| |u_1^2(P_{\tau_0}) + u_2^2(P_{\tau_0})| \\ &\quad + |g(P_{\tau_0}, u_3^1(P_{\tau_0}))| |u_1^2(P_{\tau_0}) - u_1^1(P_{\tau_0}) + u_2^2(P_{\tau_0}) - u_2^1(P_{\tau_0})|) d\tau \\ &\leq 2T[Rc(R) + M(R)] \|U^2 - U^1\|_{X_T}. \end{aligned}$$



Similarly,

$$\begin{aligned} & |(\Phi_2 U^2 - \Phi_2 U^1)(x, t)| \\ & \leq \frac{1}{2} \int_{(x+t)/2}^t (|g(Q_\tau, u_3^2(Q_\tau)) - g(Q_\tau, u_3^1(Q_\tau))| |u_1^2(Q_\tau) + u_2^2(Q_\tau)| \\ & \quad + |g(Q_\tau, u_3^1(Q_\tau))| |u_1^2(Q_\tau) - u_1^1(Q_\tau) + u_2^2(Q_\tau) - u_2^1(Q_\tau)|) d\tau \\ & \leq T[Rc(R) + M(R)] \|U^2 - U^1\|_{X_T}, \\ & |(\Phi_3 U^2 - \Phi_3 U^1)(x, t)| \leq \int_{(x+t)/2}^t |u_1^2(Q_\tau) - u_1^1(Q_\tau)| d\tau \leq T \|U^2 - U^1\|_{X_T}. \end{aligned}$$

For a fixed  $R > 0$ , let the number  $T$  be small so that

$$\max\{T, 2T(Rc(R) + M(R))\} \leq \frac{1}{2} < 1, \tag{3.7}$$

and thus

$$\|\Phi U^2 - \Phi U^1\|_{X_T} \leq \frac{1}{2} \|U^2 - U^1\|_{X_T}.$$

Thus, the operator  $\Phi$  is a contraction mapping on the set  $B_R$ , i.e., condition (2) holds.

In turn, it follows from (3.6) and (3.7) that if  $0 < T \leq T_*$ , where

$$T_* := \min\left\{ \frac{R}{2(RM(R) + R + f_\infty)}, \frac{1}{2}, \frac{1}{4(Rc(R) + M(R))} \right\}, \tag{3.8}$$

then

$$\|\Phi U\|_{X_T} \leq R \quad \text{and} \quad \|\Phi U^2 - \Phi U^1\|_{X_T} \leq \frac{1}{2} \|U^2 - U^1\|_{X_T} \quad \text{for } U, U^1, U^2 \in B_R.$$

Therefore, in view of the contraction mapping principle, there exists a solution  $U$  of Eq. (3.5) in the space  $C(\overline{D}_T; \mathbb{R}^3)$ . Theorem 3.1 is proved.  $\square$

#### 4. THE CASE OF THE GLOBAL SOLVABILITY OF PROBLEM (1.1), (1.2)

The following statement is valid.

**Theorem 4.1.** *Let conditions (2.1), (3.2), and (3.3) hold. Then, for any  $T > 0$ , problem (1.1), (1.2) has a strong generalized solution of class  $C^1$  in the domain  $D_T$ .*

**Proof.** As was noted in Remark 3.1, problem (1.1), (1.2) for the class  $C^1(\overline{D}_T)$  is equivalent to the system of nonlinear integral equations (2.22) for the class  $C(\overline{D}_T)$ . In view of (3.2), (3.3), the validity of this theorem for sufficiently small  $T$ , namely, for  $T \leq T_*$ , where  $T_*$  is given by equality (3.8), follows from Theorem 3.1. Now suppose that  $T > T_*$  and  $U^{T_*} := (u_1^{T_*}, u_2^{T_*}, u_3^{T_*})$  is a solution of the system of nonlinear integral equations (2.22) or, equivalently, of the vector equation (3.5) in the domain  $D_{T_*}$  of class  $C(\overline{D}_{T_*})$  by Theorem 3.1. For  $t > \Delta t_1 := T_*$ , system (2.22) can be rewritten as

$$\begin{cases} u_1(x, t) = -\frac{1}{2} \int_{\alpha_1(x,t,\Delta t_1)}^t g(P_\tau, u_3(P_\tau))(u_1(P_\tau) + u_2(P_\tau)) d\tau \\ \quad + \frac{1}{2} \int_{\alpha_2(x,t,\Delta t_1)}^t g(P_\tau, u_3(P_\tau))(u_1(P_\tau) + u_2(P_\tau)) d\tau + F_{1,\Delta t_1}(x, t), \\ u_2(x, t) = -\frac{1}{2} \int_{\alpha_3(x,t,\Delta t_1)}^t g(Q_\tau, u_3(Q_\tau))(u_1(Q_\tau) + u_2(Q_\tau)) d\tau + F_{2,\Delta t_1}(x, t), \\ u_3(x, t) = \int_{\alpha_3(x,t,\Delta t_1)}^t u_1(Q_\tau) d\tau + F_{3,\Delta t_1}(x, t), \end{cases} \tag{4.1}$$

where

$$\alpha_1(x, t, \Delta t_1) := \max(\Delta t_1, t - x), \quad \alpha_2(x, t, \Delta t_1) := \max\left(\Delta t_1, \frac{t}{2}\right),$$

$$\alpha_3(x, t, \Delta t_1) := \max\left(\Delta t_1, \frac{x+t}{2}\right);$$

$$\left\{ \begin{aligned} F_{1,\Delta t_1}(x, t) &:= -\frac{1}{2} \int_{t-x}^{\alpha_1(x,t,\Delta t_1)} g(P_\tau, u_3^{T_*}(P_\tau))(u_1^{T_*}(P_\tau) + u_2^{T_*}(P_\tau)) d\tau \\ &\quad + \frac{1}{2} \int_{t/2}^{\alpha_2(x,t,\Delta t_1)} g(P_{\tau_0}, u_3^{T_*}(P_{\tau_0}))(u_1^{T_*}(P_{\tau_0}) + u_2^{T_*}(P_{\tau_0})) d\tau + F_1(x, t), \\ F_{2,\Delta t_1}(x, t) &:= -\frac{1}{2} \int_{(x+t)/2}^{\alpha_3(x,t,\Delta t_1)} g(Q_\tau, u_3^{T_*}(Q_\tau))(u_1^{T_*}(Q_\tau) + u_2^{T_*}(Q_\tau)) d\tau + F_2(x, t), \\ F_{3,\Delta t_1}(x, t) &:= \int_{(x+t)/2}^{\alpha_3(x,t,\Delta t_1)} u_1^{T_*}(Q_\tau) d\tau. \end{aligned} \right. \tag{4.2}$$

Since the assumptions of Lemma 2.2 hold, it follows that, for any positive  $\tau \leq T$ , the solution of the vector equation (3.5) in the domain  $D_\tau$  of class  $C(\overline{D}_\tau)$  satisfies the a priori estimate

$$\|U\|_{C(\overline{D}_\tau)} \leq R^T(\|f\|_{C(\overline{D}_\tau)}), \tag{4.3}$$

where  $R^T = R^T(s)$  is a nondecreasing continuous function of the argument  $s \geq 0$ . Set

$$R_* := R^T(\|f\|_{C(\overline{D}_T)}).$$

In the second step with respect to  $t$ , for  $\Delta t_2$  we take

$$\Delta t_2 := \min\left\{\frac{1}{4M(R_1)R_*}, \frac{1}{4c(R_1)R_*}\right\}, \tag{4.4}$$

where

$$R_1 := 1 + 2TM(R_*)R_* + \|F\|_{C(\overline{D}_T)}, \quad F := (F_1, F_2, F_3). \tag{4.5}$$

For  $t \in [T_*, T_* + \Delta t_2]$ , the system of equations (4.1) can be rewritten as the single vector equation

$$U = \Psi U, \tag{4.6}$$

where the operator  $\Psi := (\Psi_1, \Psi_2, \Psi_3)$  acts by the formula

$$\left\{ \begin{aligned} (\Psi_1 U)(x, t) &= -\frac{1}{2} \int_{\alpha_1(x,t,\Delta t_1)}^t g(P_\tau, u_3(P_\tau))(u_1(P_\tau) + u_2(P_\tau)) d\tau \\ &\quad + \frac{1}{2} \int_{\alpha_2(x,t,\Delta t_1)}^t g(P_{\tau_0}, u_3(P_{\tau_0}))(u_1(P_{\tau_0}) + u_2(P_{\tau_0})) d\tau + F_{1,\Delta t_1}(x, t), \\ (\Psi_2 U)(x, t) &= -\frac{1}{2} \int_{\alpha_3(x,t,\Delta t_1)}^t g(Q_\tau, u_3(Q_\tau))(u_1(Q_\tau) + u_2(Q_\tau)) d\tau + F_{2,\Delta t_1}(x, t), \\ (\Psi_3 U)(x, t) &= \int_{\alpha_3(x,t,\Delta t_1)}^t u_1(Q_\tau) d\tau + F_{3,\Delta t_1}(x, t). \end{aligned} \right. \tag{4.7}$$

First, let us show that the operator  $\Psi$  takes the ball

$$B([T_1, T_2]; R_1) := \{U \in C(\overline{D}_{T_1, T_2}) : \|U\|_{C(\overline{D}_{T_1, T_2})} \leq R_1\}$$

into itself, where

$$T_1 = T_*, \quad T_2 = T_* + \Delta t_2, \quad \overline{D}_{T_1, T_2} := \overline{D} \cap \{T_1 \leq t \leq T_2\}.$$

Indeed, in view of (3.3), (4.2)–(4.5), and (4.7) we can write

$$\begin{aligned} \|\Psi_1 U\|_{C(\overline{D}_{T_1, T_2})} &\leq 2M(R_1)R_1\Delta t_2 + 2M(R_*)R_*\Delta t_1 + \|F_1\|_{C(\overline{D}_T)} \\ &\leq 2^{-1} + 2TM(R_*)R_* + \|F\|_{C(\overline{D}_T)} \leq R_1. \end{aligned}$$

Similarly,

$$\|\Psi_i U\|_{C(\overline{D}_{T_1, T_2})} \leq R_1, \quad i = 2, 3,$$

and thus, finally, we have

$$\|\Psi U\|_{C(\overline{D}_{T_1, T_2})} \leq R_1.$$

Let us now show that the operator  $\Psi$  is a contraction mapping in this ball. Indeed, for  $(x, t) \in \overline{D}_{T_1, T_2}$ , using (3.3), (4.4), and (4.7), we obtain

$$\begin{aligned} &|(\Psi_1 U^2 - \Psi_1 U^1)(x, t)| \\ &\leq \frac{1}{2} \int_{\alpha_1(x, t, \Delta t_1)}^t (|g(P_\tau, u_3^2(P_\tau)) - g(P_\tau, u_3^1(P_\tau))| |u_1^2(P_\tau) + u_2^2(P_\tau)| \\ &\quad + |g(P_\tau, u_3^1(P_\tau))| |u_1^2(P_\tau) - u_1^1(P_\tau) + u_2^2(P_\tau) - u_2^1(P_\tau)|) d\tau \\ &\quad + \frac{1}{2} \int_{\alpha_1(x, t, \Delta t_1)}^t (|g(P_{\tau_0}, u_3^2(P_{\tau_0})) - g(P_{\tau_0}, u_3^1(P_{\tau_0}))| |u_1^2(P_{\tau_0}) + u_2^2(P_{\tau_0})| \\ &\quad + |g(P_{\tau_0}, u_3^1(P_{\tau_0}))| |u_1^2(P_{\tau_0}) - u_1^1(P_{\tau_0}) + u_2^2(P_{\tau_0}) - u_2^1(P_{\tau_0})|) d\tau \\ &\leq 2c(R_1)R_1\Delta t_2 \|u_3^2 - u_3^1\|_{C(\overline{D}_{T_1, T_2})} + 2M(R_1)\Delta t_2 \|U^2 - U^1\|_{C(\overline{D}_{T_1, T_2})} \\ &\leq \frac{1}{2} \|u_3^2 - u_3^1\|_{C(\overline{D}_{T_1, T_2})} + \frac{1}{2R_1} \|U^2 - U^1\|_{C(\overline{D}_{T_1, T_2})} \\ &\leq \left(\frac{1}{2} + \frac{1}{2R_1}\right) \|U^2 - U^1\|_{C(\overline{D}_{T_1, T_2})} = q_1 \|U^2 - U^1\|_{C(\overline{D}_{T_1, T_2})}, \end{aligned}$$

where  $q_1 := (1/2)(1 + 1/R_1) < 1$ , because  $R_1 > 1$  in view of (4.5). Similarly, we find that

$$|(\Psi_i U^2 - \Psi_i U^1)(x, t)| \leq q_i \|U^2 - U^1\|_{C(\overline{D}_{T_1, T_2})}, \quad 0 < q_i := \text{const} < 1, \quad i = 2, 3.$$

Hence, in view of the contraction mapping theorem, we establish the solvability of system (4.6) for the class  $C(\overline{D}_{T_1, T_2})$ . Continuing this process step by step and taking into account the fact that, in view of the global a priori estimate (4.3), the length of each step  $\Delta t_i$  is independent of its number  $i$ , we establish the global solvability of system (3.5), and hence also that of problem (1.1), (1.2) itself in the domain  $D_T$  for any  $T > 0$ . □

### 5. UNIQUENESS OF THE SOLUTION OF PROBLEM (1.1), (1.2)

**Lemma 5.1.** *Let conditions (3.2), (3.3) hold. Then, for any  $T > 0$ , problem (1.1), (1.2) cannot have more than one strong generalized solution of class  $C^1$  in the domain  $D_T$ .*

**Proof.** Indeed, suppose that problem (1.1), (1.2) has two possible different strong generalized solutions  $u^1$  and  $u^2$  of class  $C^1$  in the domain  $D_T$ . By definition 1.1, there exists a sequence of functions  $u_n^i \in \dot{C}^2(\overline{D}_T, \Gamma_T)$  such that

$$\lim_{n \rightarrow \infty} \|u_n^i - u^i\|_{C^1(\overline{D}_T)} = 0, \quad \lim_{n \rightarrow \infty} \|Lu_n^i - f\|_{C(\overline{D}_T)} = 0, \quad i = 1, 2, \tag{5.1}$$

$$\lim_{n \rightarrow \infty} \|g(x, t, u_n^i)u_{nt}^i - g(x, t, u^i)u_t^i\|_{C(\overline{D}_T)} = 0, \quad i = 1, 2. \tag{5.2}$$

Let us use the well-known notation  $\square := \partial^2/\partial t^2 - \partial^2/\partial x^2$  and set  $\omega_{nm} := u_n^2 - u_m^1$ . It is easy to see that the function  $\omega_{nm} \in \dot{C}^2(\overline{D}_T, \Gamma_T)$  satisfies the following equalities:

$$\square \omega_{nm} + g_{nm} = f_{nm}, \tag{5.3}$$

$$\left. \frac{\partial \omega_{nm}}{\partial x} \right|_{\gamma_{2,T}} = 0, \quad \omega_{nm}|_{\gamma_{1,T}} = 0, \tag{5.4}$$

where

$$g_{nm} := g(x, t, u_n^2)u_{nt}^2 - g(x, t, u_m^1)u_{mt}^1, \quad f_{nm} := Lu_n^2 - Lu_m^1. \tag{5.5}$$

By the first equality from (5.1), there exists a number  $A := \text{const} > 0$  independent of the indices  $i$  and  $n$  such that

$$\|u_n^i\|_{C^1(\overline{D}_T)} \leq A. \tag{5.6}$$

By the second equalities from (5.1) and (5.5), we have

$$\lim_{n,m \rightarrow \infty} \|f_{nm}\|_{C(\overline{D}_T)} = 0. \tag{5.7}$$

In view of (3.2), (3.3), (5.6) and the first equality from (5.5), it is easy to see that

$$\begin{aligned} g_{nm}^2 &= \left( g(x, t, u_n^2) \frac{\partial \omega_{nm}}{\partial t} + (g(x, t, u_n^2) - g(x, t, u_m^1)) u_{mt}^1 \right)^2 \\ &\leq 2M^2(A) \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 + 2A^2 c^2(A) \omega_{nm}^2. \end{aligned} \tag{5.8}$$

Multiplying both sides of relation (5.3) by  $\partial \omega_{nm}/\partial t$ , integrating the resulting equality over the domain  $D_\tau$ , and using (5.4), just as in the derivation of inequality (2.11), from (2.5), (2.6), we obtain

$$w_{nm}(\tau) := \int_{\Omega_\tau} \left[ \left( \frac{\partial \omega_{nm}}{\partial x} \right)^2 + \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 \right] dx = 2 \int_{D_\tau} (f_{nm} - g_{nm}) \frac{\partial \omega_{nm}}{\partial t} dx dt. \tag{5.9}$$

By estimate (5.8) and Cauchy's inequality, we have

$$\begin{aligned} &2 \int_{D_\tau} (f_{nm} - g_{nm}) \frac{\partial \omega_{nm}}{\partial t} dx dt \\ &\leq \int_{D_\tau} (f_{nm} - g_{nm})^2 dx dt + \int_{D_\tau} \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 dx dt \\ &\leq 2 \int_{D_\tau} f_{nm}^2 dx dt + 2 \int_{D_\tau} g_{nm}^2 dx dt + \int_{D_\tau} \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 dx dt \\ &\leq 2 \int_{D_\tau} f_{nm}^2 dx dt + 4A^2 c^2(A) \int_{D_\tau} \omega_{nm}^2 dx dt + (1 + 4M^2(A)) \int_{D_\tau} \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 dx dt. \end{aligned} \tag{5.10}$$

Further, in view of the equality

$$\omega_{nm}(x, t) = \int_x^t \frac{\partial \omega_{nm}(x, \tau)}{\partial t} d\tau, \quad (x, t) \in \overline{D}_T,$$

which follows from the second equality in (5.4), using standard arguments, we obtain the inequality [25, p. 63]

$$\int_{D_\tau} \omega_{nm}^2 dx dt \leq \tau^2 \int_{D_\tau} \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 dx dt. \tag{5.11}$$

It follows from (5.9)–(5.11) that

$$w_{nm}(\tau) \leq (1 + 4M^2(A) + 4\tau^2 A^2 c^2(A)) \int_{D_\tau} \left( \frac{\partial \omega_{nm}}{\partial t} \right)^2 dx dt + 2 \int_{D_\tau} f_{nm}^2 dx dt$$

$$\leq (1 + 4M^2(A) + 4T^2c^2(A)) \int_0^\tau w_{nm}(\sigma) d\sigma + 2 \int_{D_T} f_{nm}^2 dx dt.$$

Therefore, by Gronwall’s lemma [24, p. 13 (Russian transl.)], we can write

$$w_{nm}(\tau) \leq c_2 \|f_{nm}\|_{L_2(D_T)}^2, \quad 0 < \tau \leq T, \tag{5.12}$$

where

$$c_2 := 2 \exp(1 + 4M^2(A) + 4T^2 A^2 c^2(A))T.$$

Arguing in the same way as in the derivation of estimate (2.13), taking into account the obvious inequality

$$\|f_{nm}\|_{L_2(D_T)}^2 \leq \|f_{nm}\|_{C(\overline{D_T})}^2 \text{mes } D_T,$$

and using (5.12), for  $(x, t) \in \overline{D_T}$ , we can write

$$|\omega_{nm}(x, t)|^2 \leq t w_{nm}(t) \leq T c_2 \text{mes } D_T \|f_{nm}\|_{C(\overline{D_T})}^2 = \frac{c_2 T^3}{2} \|f_{nm}\|_{C(\overline{D_T})}^2.$$

This yields

$$\|\omega_{nm}\|_{C(\overline{D_T})} \leq T \sqrt{\frac{c_2 T}{2}} \|f_{nm}\|_{C(\overline{D_T})}. \tag{5.13}$$

By the definition of the function  $\omega_{nm}$  and, in view of the first equality, it is easy to see that

$$\lim_{n,m \rightarrow \infty} \|\omega_{nm}\|_{C^1(\overline{D_T})} = \|u^2 - u^1\|_{C^1(\overline{D_T})}$$

and, particularly,

$$\lim_{n,m \rightarrow \infty} \|\omega_{nm}\|_{C(\overline{D_T})} = \|u^2 - u^1\|_{C(\overline{D_T})}.$$

Therefore, in inequality (5.13), passing to the limit as  $n, m \rightarrow \infty$  and using (5.7), we obtain

$$\|u^2 - u^1\|_{C(\overline{D_T})} = 0, \quad \text{i.e., } u^1 = u^2.$$

Lemma 5.1 is proved. □

### 6. THE CASE OF THE BLOW-UP OF THE GLOBAL SOLUTION OF PROBLEM (1.1), (1.2)

In this section, we show that the violation of condition (2.1) can lead to the blow-up of the global solution of problem (1.1), (1.2) in the sense of Definition 1.2. Indeed, let  $g(x, t, s) = -|s|^\alpha s$ ,  $s \in \mathbb{R}$ , and let  $\alpha > -1$  (nonlinearity exponent).

**Lemma 6.1.** *Let  $u$  be a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$  in the sense of Definition 1.1. Then the following integral equality holds:*

$$\int_{D_T} u \square \varphi dx dt = \int_{D_T} |u|^\alpha u u_t \varphi dx dt + \int_{D_T} f \varphi dx dt \tag{6.1}$$

for any function  $\varphi$  such that

$$\varphi \in C^2(\overline{D_T}), \quad \varphi|_{t=T} = 0, \quad \varphi_t|_{t=T} = 0, \quad \varphi_x|_{\gamma_{2,T}} = 0. \tag{6.2}$$

**Proof.** By the definition of a strong generalized solution  $u$  of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ , the function  $u$  belongs to  $C^1(\overline{D}_T)$  and there exists a sequence of functions  $u_n \in \mathring{C}^2(\overline{D}_T, \Gamma_T)$  such that relations (2.5) and (2.6) hold for  $g = -|s|^\alpha s$ . Set  $f_n := Lu_n$ . Let us multiply both sides of the equality  $Lu_n = f_n$  by the function  $\varphi$ , and integrate the resulting equality over the domain  $D_T$ . Integrating by parts the left-hand side of this equality, using (6.2) and conditions (1.2), we see that

$$\int_{D_T} u_n \square \varphi \, dx \, dt = \int_{D_T} |u_n|^\alpha u_n u_{nt} \varphi \, dx \, dt + \int_{D_T} f_n \varphi \, dx \, dt.$$

In this equality, passing to the limit as  $n \rightarrow \infty$  and using (2.5) and (2.6), we obtain (6.1). This lemma 6.1 is proved.  $\square$

Consider a function  $\varphi^0 := \varphi^0(x, t)$  such that

$$\begin{aligned} \varphi^0 \in C^2(\overline{D}_\infty), \quad \varphi^0 + \varphi_t^0 \leq 0, \quad \varphi^0|_{D_{T=1}} > 0, \\ \varphi_x^0|_{\gamma_{2,\infty}} = 0, \quad \varphi^0|_{t \geq 1} = 0 \end{aligned} \tag{6.3}$$

and also the number

$$\kappa_0 := \int_{D_{T=1}} \frac{|\square \varphi^0|^{p'}}{|\varphi^0|^{p'-1}} \, dx \, dt < +\infty, \quad p' = \frac{\alpha + 2}{\alpha + 1}. \tag{6.4}$$

It is easy to verify that, for sufficiently large positive constants  $n$  and  $m$ , the function  $\varphi^0$  satisfying conditions (6.3) and (6.4) can be taken as the function

$$\varphi^0(x, t) = \begin{cases} x^n(1-t)^m, & (x, t) \in D_{T=1}, \\ 0, & t \geq 1. \end{cases}$$

Set  $\varphi_T(x, t) := \varphi^0(x/T, t/T)$ ,  $T > 0$ . In view of (6.3), it is easy to see that

$$\begin{aligned} \varphi_T \in C^2(\overline{D}_T), \quad \varphi_T + T \frac{\partial \varphi_T}{\partial t} \leq 0, \quad \varphi_T|_{D_T} > 0, \\ \frac{\partial \varphi_T}{\partial x} \Big|_{\gamma_{2,T}} = 0, \quad \varphi_T|_{t=T} = 0, \quad \frac{\partial \varphi_T}{\partial t} \Big|_{t=T} = 0. \end{aligned} \tag{6.5}$$

For a given  $f$ , consider the function

$$\zeta(T) := \int_{D_T} f \varphi_T \, dx \, dt, \quad T > 0. \tag{6.6}$$

The following theorem on the blow-up of the global solution of problem (1.1), (1.2) is valid.

**Theorem 6.1.** *Let*

$$g(x, t, s) = -|s|^\alpha s, \quad s \in \mathbb{R}, \quad \alpha > -1,$$

*let  $f \in C(\overline{D}_\infty)$ , and suppose that  $f \geq 0$  in the domain  $D_\infty$ . Then if*

$$\liminf_{T \rightarrow +\infty} \zeta(T) > 0, \tag{6.7}$$

*then there exists a positive number  $T^* := T^*(f)$  such that, for  $T > T^*$ , problem (1.1), (1.2) cannot have a strong generalized solution  $u$  of class  $C^1$  in the domain  $D_T$ .*

**Proof.** Suppose that, under the assumptions of this theorem, there exists a strong generalized solution  $u$  of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ . Then, by Lemma 6.1, equality (6.1) holds; in view of (6.5),  $\varphi$  in this equality can be taken as the function  $\varphi = \varphi_T$  i.e.,

$$\int_{D_T} u \square \varphi_T \, dx \, dt = \int_{D_T} |u|^\alpha u u_t \varphi_T \, dx \, dt + \int_{D_T} f \varphi_T \, dx \, dt. \tag{6.8}$$

By (1.2) and (6.5), we have

$$\begin{aligned} \int_{D_T} |u|^\alpha uu_t \varphi_T \, dx \, dt &= \frac{1}{\alpha + 2} \int_{D_T} \varphi_T \frac{\partial}{\partial t} |u|^{\alpha+2} \, dx \, dt \\ &= -\frac{1}{\alpha + 2} \int_{D_T} |u|^{\alpha+2} \frac{\partial \varphi_T}{\partial t} \, dx \, dt \geq \frac{1}{(\alpha + 2)T} \int_{D_T} |u|^{\alpha+2} \varphi_T \, dx \, dt. \end{aligned}$$

Hence, using (6.6), from (6.8) we obtain

$$\frac{1}{pT} \int_{D_T} |u|^p \varphi_T \, dx \, dt \leq \int_{D_T} u \square \varphi_T \, dx \, dt - \zeta(T), \quad p := \alpha + 2 > 1. \tag{6.9}$$

If, in Young’s inequality with parameter  $\varepsilon > 0$ ,

$$ab \leq \frac{\varepsilon}{p} a^p + \frac{1}{p' \varepsilon^{p'-1}} b^{p'}, \quad a, b \geq 0, \quad \frac{1}{p} + \frac{1}{p'} = 1, \quad p > 1,$$

we put

$$a = |u| \varphi_T^{1/p}, \quad b = \frac{|\square \varphi_T|}{\varphi_T^{1/p}},$$

then, using the equality  $p'/p = p' - 1$ , we obtain

$$|u \square \varphi_T| = |u| \varphi_T^{1/p} \frac{|\square \varphi_T|}{\varphi_T^{1/p}} \leq \frac{1}{pT} |u|^p \varphi_T + \frac{T^{p'-1}}{p'} \frac{|\square \varphi_T|^{p'}}{\varphi_T^{p'-1}}.$$

In view of (6.9) and the last inequality, we can write

$$\frac{1 - \varepsilon}{p} \int_{D_T} \varphi_T \, dx \, dt \leq \frac{1}{p' \varepsilon^{p'-1}} \int_{D_T} \frac{|\square \varphi_T|^{p'}}{\varphi_T^{p'-1}} \, dx \, dt - \zeta(T),$$

whence we have

$$\int_{D_T} |u|^p \varphi_T \, dx \, dt \leq \frac{p}{(1 - \varepsilon) p' \varepsilon^{p'-1}} \int_{D_T} \frac{|\square \varphi_T|^{p'}}{\varphi_T^{p'-1}} \, dx \, dt - \frac{p}{1 - \varepsilon} \zeta(T).$$

Taking into account the equalities  $p' = p/(p - 1)$ ,  $p = p'/(p' - 1)$ , and

$$\min_{0 < \varepsilon < 1} \frac{p}{(1 - \varepsilon) p' \varepsilon^{p'-1}} = p^{p'},$$

where the minimum is realized at  $\varepsilon = 1/p$ , from the last inequality we obtain

$$\int_{D_T} |u|^p \varphi_T \, dx \, dt \leq p^{p'} \int_{D_T} \frac{|\square \varphi_T|^{p'}}{\varphi_T^{p'-1}} \, dx \, dt - \frac{p}{1 - \varepsilon} \zeta(T). \tag{6.10}$$

Since  $\varphi_T(x, t) := \varphi^0(x/T, t/T)$ , using (6.3), (6.4) and making the change of variables  $x = Tx_1$ ,  $t = Tt_1$ , we can easily verify the relations

$$\int_{D_T} \frac{|\square \varphi_T|^{p'}}{\varphi_T^{p'-1}} \, dx \, dt = T^{-2(p'-1)} \int_{D_{T=1}} \frac{|\square \varphi^0|^{p'}}{|\varphi^0|^{p'-1}} \, dx_1 \, dt_1 = T^{-2(p'-1)} \kappa_0.$$

Hence, using (6.5), from (6.10) we obtain

$$0 \leq \frac{\kappa_0}{p' T^{p'-1}} - \zeta(T). \tag{6.11}$$

Since  $p' = p/(p - 1) > 1$ , it follows that  $-2(p' - 1) < 0$  and, in view of (6.4), we will have

$$\lim_{T \rightarrow +\infty} \frac{\kappa_0}{p' T^{p'-1}} = 0.$$

Therefore, in view of (6.7), there exists a positive number  $T^* := T^*(f)$  such that, for  $T > T^*$ , the right-hand side of inequality (6.11) is negative, while the left-hand side of this inequality is nonnegative. This implies that if  $u$  is a strong generalized solution of problem (1.1), (1.2) of class  $C^1$  in the domain  $D_T$ , then necessarily  $T \leq T^*$ , which proves Theorem 6.1.  $\square$

**Remark 6.1.** It is easy to verify that if

$$f \in C(\overline{D_\infty}), \quad f \geq 0, \quad f(x, t) \geq ct^{-m} \quad \text{for } t \geq 1,$$

where  $c = \text{const} > 0$  and  $0 \leq m = \text{const} \leq 2$ , then condition (6.7) will hold, and thus, for  $g = -|s|^\alpha s$ ,  $s \in \mathbb{R}$ ,  $\alpha > -1$ , problem (1.1), (1.2) does not have a strong generalized solution  $u$  of class  $C^1$  in the domain  $D_T$  for sufficiently large  $T$ . Indeed, in (6.6), introducing the transformation of the independent variables  $x$  and  $t$  by the formulas  $x = Tx_1$ ,  $t = Tt_1$  and assuming that  $T > 1$ , after a few manipulations, we obtain

$$\begin{aligned} \zeta(T) &= T^2 \int_{D_{T=1}} f(Tx_1, Tt_1) \varphi^0(x_1, t_1) dx_1 dt_1 \\ &\geq cT^{2-m} \int_{D_{T=1} \cap \{t_1 \geq T^{-1}\}} t_1^{-m} \varphi^0(x_1, t_1) dx_1 dt_1 \\ &\quad + T^2 \int_{D_{T=1} \cap \{t_1 < T^{-1}\}} f(Tx_1, Tt_1) \varphi^0(x_1, t_1) dx_1 dt_1. \end{aligned}$$

Further, let  $T_1 > 1$  be an arbitrary fixed number. Then, using the last inequality for the function  $\zeta$ , we obtain

$$\zeta(T) \geq cT^{2-m} \int_{D_{T=1} \cap \{t_1 \geq T^{-1}\}} t_1^{-m} \varphi^0(x_1, t_1) dx_1 dt_1 \geq c \int_{D_{T=1} \cap \{t_1 \geq T_1^{-1}\}} t_1^{-m} \varphi^0(x_1, t_1) dx_1 dt_1,$$

if  $T \geq T_1 > 1$ . The last inequality immediately yields condition (6.7).

#### ACKNOWLEDGMENTS

This work was supported by INTAS (grant no. 05-1000008-7921) and by the Georgian National Foundation (grant no. GNSF/ST06/3-005).

#### REFERENCES

1. J.-L. Lions, *Quelques méthodes de résolution des problèmes aux limites nonlinéaires* (Dunod, Paris, 1969; Mir, Moscow, 1972).
2. J.-L. Lions and W. A. Strauss, "Some non linear evolution equations," *Bull. Soc. Math. France* **93**, 43–96 (1965).
3. A. V. Bitsadze, *Some Classes of Partial Differential Equations* (Nauka, Moscow, 1981) [in Russian].
4. F. John, "Blow-up of solutions of nonlinear wave equation in three space dimensions," *Manuscripta Math.* **28** (1-3), 235–268 (1979).
5. T. Kato, "Blow-up of solutions of some nonlinear hyperbolic equations," *Comm. Pure Appl. Math.* **33** (4), 501–505 (1980).
6. T. C. Sideris, "Nonexistence of global solutions to semilinear wave equations in high dimensions," *J. Differential Equations* **52** (3), 378–406 (1984).
7. V. Georgiev, H. Lindblad, and C. D. Sogge, "Weighted Strichartz estimates and global existence for semilinear wave equations," *Amer. J. Math.* **119** (6), 1291–1319 (1997).
8. L. Hörmander, *Lectures on Nonlinear Hyperbolic Differential Equations*, in *Math. Appl. (Berlin)* (Springer-Verlag, Berlin, 1997), Vol. 26.
9. S. I. Pohozaev and L. Véron, "Blow-up results for nonlinear hyperbolic inequaties," *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)* **29** (2), 393–420 (2001).
10. É. Mitidieri and S. I. Pokhozhaev, "A priori estimates and blow-up of solutions to nonlinear partial differential equations and inequalities," in *Trudy Mat. Inst. Steklov* (Nauka, Moscow, 2001), Vol. 234, pp. 3–383 [Proc. Steklov Inst. Math. **234**, 1–362 (2001)].
11. G. Todorova and E. Vitillaro, "Blow-up for nonlinear dissipative wave equations in  $\mathbb{R}^n$ ," *J. Math. Anal. Appl.* **303** (1), 242–257 (2005).



12. L. Liu and M. Wang, “Global existence and blow-up of solutions for some hyperbolic systems with damping and source terms,” *Nonlinear Anal.* **64** (1), 69–91 (2006).
13. J. Zhu, “Blow-up of solutions of a semilinear hyperbolic equation and a parabolic equation with general forcing term and boundary condition,” *Nonlinear Anal.* **67** (1), 33–38 (2007).
14. S. Kharibegashvili, “On the solvability of one multidimensional version of the first Darboux problem for some nonlinear wave equations,” *Nonlinear Anal.* **68** (4), 912–924 (2008).
15. S. S. Kharibegashvili, “On the solvability of the characteristic Cauchy problem for some nonlinear wave equations in the future light cone,” *Differ. Uravn.* **44** (1), 129–139 (2008) [*Differ. Equations* **44** (1), 135–146 (2008)].
16. O. Jokhadze, “On existence and nonexistence of global solutions of Cauchy–Goursat problem for nonlinear wave equations,” *J. Math. Anal. Appl.* **340** (2), 1033–1045 (2008).
17. G. K. Berikelashvili, O. M. Jokhadze, B. G. Midodashvili, and S. S. Kharibegashvili, “On the existence and nonexistence of global solutions of the first Darboux problem for nonlinear wave equations,” *Differ. Uravn.* **44** (3), 359–372 (2008) [*Differ. Equations* **44** (3), 374–389 (2008)].
18. O. M. Jokhadze and S. S. Kharibegashvili, “On the first Darboux problem for second-order nonlinear hyperbolic equations,” *Mat. Zametki* **84** (5), 693–712 (2008) [*Math. Notes* **84** (5), 646–663 (2008)].
19. E. Goursat, *Course of Mathematical Analysis*, Vol. 3, Part 1: *Infinitely Near Integrals. Partial Differential Equations* (Gostekhizdat, Moscow–Leningrad, 1933) [Russian transl.].
20. E. I. Moiseev, “Approximation of the classical solution of the Darboux problem by smooth solutions,” *Differ. Uravn.* **20** (1), 73–87 (1984).
21. E. I. Moiseev, *Equations of Mixed Type with Spectral Parameter* (Izd. Moskov. Univ., Moscow, 1988) [in Russian].
22. A. M. Nakhushhev, *Equations of Mathematical Biology* (Vysshaya Shkola, Moscow, 1995) [in Russian].
23. S. Kharibegashvili, “Goursat and Darboux type problems for linear hyperbolic partial differential equations and systems,” *Mem. Differential Equations Math. Phys.* **4**, 1–127 (1995).
24. D. Henry *Geometric Theory of Semilinear Parabolic Equations* (Springer-Verlag, Heidelberg, 1981; Mir, Moscow, 1985).
25. O. A. Ladyzhenskaya, *Boundary-Value Problems of Mathematical Physics* (Nauka, Moscow, 1973) [in Russian].
26. V. A. Trenogin, *Functional Analysis* (Nauka, Moscow, 1980) [in Russian].