# On the Solvability of a Boundary Value Problem for Nonlinear Wave Equations in Angular Domains 

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#### Abstract

For a one-dimensional wave equation with a weak nonlinearity, we study the Darboux boundary value problem in angular domains, for which we analyze the existence and uniqueness of a global solution and the existence of local solutions as well as the absence of global solutions.


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## 1. STATEMENT OF THE PROBLEM

In the plane of independent variables $x$ and $t$, consider the nonlinear wave equation

$$
\begin{equation*}
L u:=\square u+f(x, t, u)=F(x, t) \tag{1.1}
\end{equation*}
$$

where $f=f(x, t, s)$ is a given real function nonlinear with respect to the variable $s, F=F(x, t)$ is a given real function, and $u=u(x, t)$ is the unknown real function; moreover, we assume that $f$ and $F$ are continuous functions of their arguments and $\square:=\partial^{2} / \partial t^{2}-\partial^{2} / \partial x^{2}$.

By $D: \gamma_{2}(t)<x<\gamma_{1}(t), t>0$, we denote the angular domain lying inside the characteristic angle $\Lambda_{0}: t>|x|$ and bounded by smooth noncharacteristic curves $\gamma_{i}: x=\gamma_{i}(t), t \geq 0, i=1$, 2 (i.e., $\left.\left|\gamma_{i}^{\prime}(t)\right| \neq 1, t \geq 0, i=1,2\right)$, of the class $C^{2}$ issuing from the origin $O(0,0)$. Set $D_{T}:=D \cap\{t<T\}$ and $\gamma_{i, T}:=\gamma_{i} \cap\{t \leq T\}, T>0, i=1,2$.

For Eq. (1.1), we consider the Darboux boundary value problem, where the directional derivative of the solution of Eq. (1.1) is posed on $\gamma_{1, T}$, and a solution itself is posed on $\gamma_{2, T}$, in the following statement: find a solution $u=u(x, t)$ of that equation in the domain $D_{T}$ with the boundary conditions

$$
\begin{align*}
\left.\left(l_{1} u_{x}+l_{2} u_{t}\right)\right|_{\gamma_{1, T}} & =0  \tag{1.2}\\
\left.u\right|_{\gamma_{2, T}} & =0 \tag{1.3}
\end{align*}
$$

where $l_{1}$ and $l_{2}$ are given continuous functions; moreover, $\left.\left(\left|l_{1}\right|+\left|l_{2}\right|\right)\right|_{\gamma_{1}} \neq 0$.
Note that, in the linear case in which the function $f$ occurring in Eq. (1) is linear with respect to the variable $s$ and the conditions

$$
\begin{equation*}
\left.\left(\alpha_{i} u_{x}+\beta_{i} u_{t}\right)\right|_{\gamma_{i, T}}=0, \quad i=1,2 ; \quad u(0,0)=0 \tag{1.4}
\end{equation*}
$$

are considered instead of the boundary conditions (1.2) and (1.3), problem (1.1), (1.4) in the domain $D_{T}$ was studied in $[1-6]$. Note also that problem (1.1)-(1.3) is equivalent to problem (1.1), (1.4) in which the direction $\left(\alpha_{2}, \beta_{2}\right)$ coincides with the direction of the tangent to the curve $\gamma_{2, T}$ at each of its points. In the case of Eq. (1.1) with a power-law nonlinearity in which the homogeneous Dirichlet conditions $\left.u\right|_{\gamma_{i, T}}=0, i=1,2$, are posed on the curves $\gamma_{1}$ and $\gamma_{2}$, and moreover, one of these curves,
either $\gamma_{1}$ or $\gamma_{2}$, is characteristic, this problem was considered in the papers [7-9], and the case in which both curves are noncharacteristic was studied in [10]. The special case of the boundary conditions (1.2) and (1.3) of the form $\left.u_{x}\right|_{\gamma_{1, T}}=0$ and $\left.u\right|_{\gamma_{2, T}}=0$, where $\gamma_{1, T}: x=0,0 \leq t \leq T$, and $\gamma_{2, T}: x=-t, 0 \leq t \leq T$, is a characteristic of Eq. (1.1) with a power-law nonlinearity, was considered in $[11,12]$; moreover, the case in which $\gamma_{2, T}$ is a noncharacteristic curve was considered in $[13,14]$. As was noted in $[1,6]$, such problems arise in the mathematical modeling of small harmonic oscillations of a wedge in a supersonic flow and oscillations of a string inside a cylinder filled with a viscous fluid.

In the present paper, we consider the more general case of a nonlinear function $f(x, t, s)$, smooth noncharacteristic curves $\gamma_{1}$ and $\gamma_{2}$, and the behavior of the vector field $\left(l_{1}, l_{2}\right)$ in the boundary condition (1.2) as compared with the cases studied in the above-mentioned papers. Note that, in the case under consideration, the analysis of the solvability of problem (1.1)-(1.3) encounters additional difficulties of nontechnical character.

Set $\dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right):=\left\{v \in C^{2}\left(\bar{D}_{T}\right):\left.\left(l_{1} v_{x}+l_{2} v_{t}\right)\right|_{\gamma_{1, T}}=0,\left.v\right|_{\gamma_{2, T}}=0\right\}$ and $\gamma_{T}:=\gamma_{1, T} \cup \gamma_{2, T}$.
Definition 1.1. Let the conditions $f \in C\left(\bar{D}_{T} \times \mathbb{R}\right), F \in C\left(\bar{D}_{T}\right)$, and $l_{1}, l_{2} \in C\left(\gamma_{1, T}\right)$ be satisfied. A function $u$ is called a strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$ if $u$ belongs to $C\left(\bar{D}_{T}\right)$ and there exists a function sequence $u_{n} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ such that $u_{n} \rightarrow u$ and $L u_{n} \rightarrow F$ in the space $C\left(\bar{D}_{T}\right)$ as $n \rightarrow \infty$.

Remark 1.1. Obviously, a classical solution of problem (1.1)-(1.3) in the space $\dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ is a strong generalized solution of that problem in the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Definition 1.2. Let $f \in C\left(\bar{D}_{\infty} \times \mathbb{R}\right), F \in C\left(\bar{D}_{\infty}\right)$, and $l_{1}, l_{2} \in C\left(\gamma_{1, \infty}\right)$. We say that problem (1.1)-(1.3) is globally solvable in the class $C$ if, for any finite $T>0$, it has at least one strong generalized solution of the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Definition 1.3. Under the assumptions of Definition 1.2, a function $u \in C\left(\bar{D}_{\infty}\right)$ is called a global strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{\infty}$ if, for any finite $T>0$, the function $\left.u\right|_{D_{T}}$ is a strong generalized solution of that problem in the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Definition 1.4. Under the assumptions of Definition 1.2 , we say that problem (1.1)-(1.3) is locally solvable in the class $C$ if there exists a positive number $T_{0}=T_{0}(F)$ such that, for $T \leq T_{0}$, it has at least one strong generalized solution of the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Remark 1.2. Under the above-mentioned assumptions, one can readily note that

$$
\begin{equation*}
-t<\gamma_{2}(t)<\gamma_{1}(t)<t, \quad t>0 ; \quad\left|\gamma_{i}^{\prime}(t)\right|<1, \quad t \geq 0, \quad \gamma_{i}(0)=0, \quad i=1,2 \tag{1.5}
\end{equation*}
$$

Remark 1.3. Below, without loss of generality, one can assume that $\gamma_{1}(t) \leq 0,0 \leq t \leq T$, since otherwise, by virtue of condition (1.5), this could be achieved with the use of the Lorentz transformation

$$
x^{\prime}=\frac{x-k_{0} t}{\sqrt{1-k_{0}^{2}}}, \quad t^{\prime}=\frac{t-k_{0} x}{\sqrt{1-k_{0}^{2}}}, \quad k_{0}:=\max _{0 \leq t \leq T}\left|\gamma_{1}^{\prime}(t)\right|<1
$$

which does not change the form of Eq. (1.1) and maps the characteristic angle $\Lambda_{0}: t>|x|$ into the characteristic angle $\Lambda_{0}^{\prime}: t^{\prime}>\left|x^{\prime}\right|$.

Below, by virtue of Remark 1.3, in addition to condition (1.5), we require that

$$
\begin{equation*}
\gamma_{2}(t)<\gamma_{1}(t) \leq 0, \quad \gamma_{1}^{\prime}(t) \leq 0, \quad \gamma_{2}^{\prime}(t)<0, \quad t>0 \tag{1.6}
\end{equation*}
$$

In Section 2, we present conditions on the data of problem (1.1)-(1.3) under which we prove an a priori estimate for a strong generalized solution of that problem in the class $C$ in the domain $D_{T}$. In Section 3, we study the global solvability of the problem in the class $C$ in the
domain $D_{T}$. We analyze the smoothness of the solution in Section 4 and study the uniqueness and existence of a global solution of problem (1.1)-(1.3) in the domain $D_{\infty}$ in Section 5. Finally, in Section 6, we study the case of absence of a global solution as well as the local solvability of that problem.

## 2. A PRIORI ESTIMATE FOR THE SOLUTION OF PROBLEM (1.1)-(1.3)

Set

$$
\begin{equation*}
g(x, t, s):=\int_{0}^{s} f\left(x, t, s_{1}\right) d s_{1}, \quad(x, t, s) \in \bar{D}_{T} \times \mathbb{R} \tag{2.1}
\end{equation*}
$$

In view of notation (2.1), consider the following conditions imposed on the nonlinear function $f$ :

$$
\begin{align*}
g(x, t, s) & \geq-M_{1}-M_{2} s^{2}, \quad(x, t, s) \in \bar{D}_{T} \times \mathbb{R}  \tag{2.2}\\
g_{t}(x, t, s) & \leq M_{3}+M_{4} s^{2}, \quad(x, t, s) \in \bar{D}_{T} \times \mathbb{R} \tag{2.3}
\end{align*}
$$

where $M_{i}:=M_{i}(T)=$ const $\geq 0,1 \leq i \leq 4$.
Remark 2.1. Let $f_{0}, f_{0 t} \in C\left(\bar{D}_{\infty}\right), f_{0} \geq 0$, and $f_{0 t} \leq 0$. We present some classes of functions $f=f(x, t, s)$ that are often used in applications and satisfy conditions (2.2) and (2.3).

1. $f(x, t, s)=f_{0}(x, t)|s|^{\alpha} \operatorname{sgn} s$, where $\alpha>0, \alpha \neq 1$. In this case, we have

$$
g(x, t, s)=f_{0}(x, t) \frac{|s|^{\alpha+1}}{\alpha+1}
$$

2. $f(x, t, s)=f_{0}(x, t) \psi(s)$, where $\psi$ belongs to $C(\mathbb{R}), \psi(s) \operatorname{sgn} s \geq 0$, and $s \in \mathbb{R}$. Here

$$
g(x, t, s)=f_{0}(x, t) \int_{0}^{s} \psi(\tau) d \tau
$$

3. $f(x, t, s)=f_{0}(x, t) e^{s}$. In this case, we have $g(x, t, s)=f_{0}(x, t)\left(e^{s}-1\right)$.

Now we subject the curve $\gamma_{1, T}$ to an additional constraint of the geometric nature, which depends on the direction of the vector $\left(l_{1}, l_{2}\right)$ of the directional derivative occurring in the boundary condition (1.2),

$$
\begin{equation*}
\left[\left(l_{1}^{2}+l_{2}^{2}\right) \nu_{t}+2 l_{1} l_{2} \nu_{x}\right](P) \geq 0, \quad P \in \gamma_{1, T} \tag{2.4}
\end{equation*}
$$

where $\nu:=\left(\nu_{x}, \nu_{t}\right)$ is the unit outward normal to $\partial D_{T}$ at the point $P$.
Remark 2.2. By virtue of conditions (1.5) and (1.6), one can readily see that, on $\gamma_{1, T} \subset \partial D_{T}$, the unit vector $\nu:=\left(\nu_{x}, \nu_{t}\right)$ of the outward normal to $\partial D_{T}$ is defined by the relations

$$
\begin{equation*}
\nu_{x}=\frac{1}{\sqrt{1+\left|\gamma_{1}^{\prime}(t)\right|^{2}}}>0, \quad \nu_{t}=-\frac{\gamma_{1}^{\prime}(t)}{\sqrt{1+\left|\gamma_{1}^{\prime}(t)\right|^{2}}} \geq 0, \quad 0 \leq t \leq T \tag{2.5}
\end{equation*}
$$

It follows from relations (2.5) that condition (2.4) is satisfied for the case in which $\left.l_{1} l_{2}\right|_{\gamma_{1, T}} \geq 0$. In particular, if condition (1.2) is a homogeneous Neumann boundary condition, i.e., $\left.u_{\nu}\right|_{\gamma_{1, T}}=$ $\left.\left(\nu_{x} u_{x}+\nu_{t} u_{t}\right)\right|_{\gamma_{1, T}}=0$, then condition (2.4) is satisfied. One can also readily show that, for the case in which $l_{1}, l_{2}=$ const, $l_{2}=-k_{0} l_{1}$, where $k_{0}>0, k_{0} \neq 1$, and $\gamma_{1, T}: x=-k t, 0 \leq k=$ const $<1$, condition (2.4) is equivalent to the condition $k \geq 2 k_{0} /\left(1+k_{0}^{2}\right)$.

Lemma 2.1. Let $f \in C\left(\bar{D}_{T} \times \mathbb{R}\right), F \in C\left(\bar{D}_{T}\right)$, and $l_{1}, l_{2} \in C\left(\gamma_{1, T}\right)$, and let conditions (1.5), (1.6), (2.2)-(2.4) be satisfied. Then any strong generalized solution $u=u(x, t)$ of problem (1.1)-(1.3) of the class $C$ in the domain $D_{T}$ satisfies the a priori estimate

$$
\begin{equation*}
\|u\|_{C\left(\bar{D}_{T}\right)} \leq c_{1}\|F\|_{C\left(\bar{D}_{T}\right)}+c_{2} \tag{2.6}
\end{equation*}
$$

with nonnegative constants $c_{i}:=c_{i}\left(f, l_{1}, l_{2}, T\right), i=1,2$, independent of $u$ and $F$; moreover,$c_{1}>0$.

Proof. Let $u$ be a strong generalized solution of problem (1.1)-(1.3) of the class $C$ in the domain $D_{T}$. By Definition 1.1, there exists a sequence of functions $u_{n} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|u_{n}-u\right\|_{C\left(\bar{D}_{T}\right)}=0, \quad \lim _{n \rightarrow \infty}\left\|L u_{n}-F\right\|_{C\left(\bar{D}_{T}\right)}=0 \tag{2.7}
\end{equation*}
$$

Consider the function $u_{n} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ treated as a solution of the problem

$$
\begin{align*}
L u_{n} & =F_{n},  \tag{2.8}\\
\left.\left(l_{1} u_{n x}+l_{2} u_{n t}\right)\right|_{\gamma_{1, T}} & =0  \tag{2.9}\\
\left.u_{n}\right|_{\gamma_{2, T}} & =0 \tag{2.10}
\end{align*}
$$

where

$$
\begin{equation*}
F_{n}:=L u_{n} \tag{2.11}
\end{equation*}
$$

By multiplying both sides of Eq. (2.8) by $u_{n t}$, by integrating the resulting relation over the domain $D_{\tau}:=\left\{(x, t) \in D_{T}: t<\tau\right\}, 0<\tau \leq T$, and by using relation (2.1), we obtain

$$
\begin{align*}
\frac{1}{2} \int_{D_{\tau}}\left(u_{n t}^{2}\right)_{t} d x d t & -\int_{D_{\tau}} u_{n x x} u_{n t} d x d t+\int_{D_{\tau}} \frac{d}{d t}\left(g\left(x, t, u_{n}(x, t)\right) d x d t-\int_{D_{\tau}} g_{t}\left(x, t, u_{n}(x, t)\right) d x d t\right. \\
& =\int_{D_{\tau}} F_{n} u_{n t} d x d t \tag{2.12}
\end{align*}
$$

Set $\Omega_{\tau}:=\bar{D}_{\infty} \cap\{t=\tau\}, 0<\tau \leq T$. By integrating by parts on the left-hand side in relation (2.12) and by taking into account relations (2.1) and (2.10), we obtain

$$
\begin{align*}
& 2 \int_{D_{\tau}} F_{n} u_{n t} d x d t=\int_{\gamma_{1, \tau}}\left(u_{n t}^{2} \nu_{t}-2 u_{n x} u_{n t} \nu_{x}+u_{n x}^{2} \nu_{t}\right) d s \\
& \quad+\int_{\gamma_{2, \tau}} \frac{1}{\nu_{t}}\left[\left(u_{n x} \nu_{t}-u_{n t} \nu_{x}\right)^{2}+u_{n t}^{2}\left(\nu_{t}^{2}-\nu_{x}^{2}\right)\right] d s+\int_{\Omega_{\tau}}\left(u_{n t}^{2}+u_{n x}^{2}\right) d x+2 \int_{\Omega_{\tau}} g\left(x, \tau, u_{n}(x, \tau)\right) d x \\
& \quad+2 \int_{\gamma_{1, \tau}} g\left(x, t, u_{n}(x, t)\right) \nu_{t} d s-2 \int_{D_{\tau}} g_{t}\left(x, t, u_{n}(x, t)\right) d x d t \tag{2.13}
\end{align*}
$$

It follows from Eq. (2.9) that the relations

$$
\begin{equation*}
u_{n x}=-\lambda l_{2}, \quad u_{n t}=\lambda l_{1} \tag{2.14}
\end{equation*}
$$

hold on $\gamma_{1, T}$, where $\lambda$ is a proportionality coefficient.
By virtue of relations (2.4) and (2.14), we have

$$
\begin{equation*}
\int_{\gamma_{1, \tau}}\left(u_{n t}^{2} \nu_{t}-2 u_{n x} u_{n t} \nu_{x}+u_{n x}^{2} \nu_{t}\right) d s=\int_{\gamma_{1, \tau}} \lambda^{2}\left[\left(l_{1}^{2}+l_{2}^{2}\right) \nu_{t}+2 l_{1} l_{2} \nu_{x}\right] d s \geq 0 \tag{2.15}
\end{equation*}
$$

Since $\nu_{t} \frac{\partial}{\partial x}-\nu_{x} \frac{\partial}{\partial t}$ is the operator of differentiation along the direction of the tangent to $\gamma_{2, T}$, i.e., is an internal differential operator on $\gamma_{2, T}$, it follows from conditions (2.10) that

$$
\begin{equation*}
\left.\left(u_{n x} \nu_{t}-u_{n t} \nu_{x}\right)\right|_{\gamma_{2, \tau}}=0 \tag{2.16}
\end{equation*}
$$

By taking into account conditions (1.5) and (1.6), one can readily see that on $\gamma_{2, T} \subset \partial D_{T}$ the unit vector $\nu:=\left(\nu_{x}, \nu_{t}\right)$ of the outward normal to $\partial D_{T}$ is defined by the relations

$$
\begin{equation*}
\nu_{x}=-\frac{1}{\sqrt{1+\left|\gamma_{2}^{\prime}(t)\right|^{2}}}<0, \quad \nu_{t}=\frac{\gamma_{2}^{\prime}(t)}{\sqrt{1+\left|\gamma_{2}^{\prime}(t)\right|^{2}}}<0, \quad 0 \leq t \leq T \tag{2.17}
\end{equation*}
$$

moreover,

$$
\begin{equation*}
\left.\left(\nu_{t}^{2}-\nu_{x}^{2}\right)\right|_{\gamma_{2, T}}<0 \tag{2.18}
\end{equation*}
$$

It follows from relations (2.16)-(2.18) that

$$
\begin{equation*}
\int_{\gamma_{2, \tau}} \frac{1}{\nu_{t}}\left[\left(u_{n x} \nu_{t}-u_{n t} \nu_{x}\right)^{2}+u_{n t}^{2}\left(\nu_{t}^{2}-\nu_{x}^{2}\right)\right] d s \geq 0 . \tag{2.19}
\end{equation*}
$$

By virtue of inequalities (2.2), (2.3), and (2.5),

$$
\begin{align*}
\int_{\Omega_{\tau}} g\left(x, \tau, u_{n}(x, \tau)\right) d x+ & \int_{\gamma_{1, \tau}} g\left(x, t, u_{n}(x, t)\right) \nu_{t} d s-\int_{D_{\tau}} g_{t}\left(x, t, u_{n}(x, t)\right) d x d t \\
\geq & -\int_{\Omega_{\tau}}\left[M_{1}+M_{2}\left|u_{n}(x, \tau)\right|^{2}\right] d x-\int_{\gamma_{1, \tau}}\left[M_{1}+M_{2}\left|u_{n}(x, t)\right|^{2}\right] d s \\
& -\int_{D_{\tau}}\left[M_{3}+M_{4}\left|u_{n}(x, t)\right|^{2}\right] d x d t . \tag{2.20}
\end{align*}
$$

Now, by taking into account inequalities (2.15), (2.19), and (2.20), from relation (2.13) we obtain

$$
\begin{align*}
w_{n}(\tau):= & \int_{\Omega_{\tau}}\left(u_{n t}^{2}+u_{n x}^{2}\right) d x \leq 2 M_{1}\left(\operatorname{mes} \Omega_{T}+\operatorname{mes} \gamma_{1, T}\right)+2 M_{3} \operatorname{mes} D_{T} \\
& +2 M_{2} \int_{\Omega_{\tau}}\left|u_{n}(x, \tau)\right|^{2} d x+2 M_{2} \int_{\gamma_{1, \tau}}\left|u_{n}(x, t)\right|^{2} d s \\
& +2 M_{4} \int_{D_{\tau}}\left|u_{n}(x, t)\right|^{2} d x d t+2 \int_{D_{\tau}} F_{n} u_{n t} d x d t . \tag{2.21}
\end{align*}
$$

By virtue of conditions (1.5) and (1.6), one can readily see that

$$
\begin{equation*}
\operatorname{mes} \Omega_{T} \leq T, \quad \operatorname{mes} \gamma_{1, T} \leq \sqrt{2 T}, \quad \operatorname{mes} D_{T} \leq T^{2} \tag{2.22}
\end{equation*}
$$

Since $\Omega_{\tau}: \gamma_{2}(\tau) \leq x \leq \gamma_{1}(\tau), t=\tau$, and $\gamma_{2, T}: t=\gamma_{2}^{-1}(x), \gamma_{2}(T) \leq x \leq 0$, where $\gamma_{2}^{-1}$ is the function inverse to $\gamma_{2}$, which is uniquely determined by virtue of condition (1.6), we can use relations (2.10) and the Newton-Leibniz formulas to obtain

$$
\begin{gather*}
u_{n}(x, \tau)=\int_{\gamma_{2}^{-1}(x)}^{\tau} u_{n t}(x, t) d t, \quad \gamma_{2}(\tau) \leq x \leq \gamma_{1}(\tau), \quad(x, \tau) \in \Omega_{\tau},  \tag{2.23}\\
u_{n}\left(\gamma_{1}(t), t\right)=\int_{\gamma_{2}(t)}^{\gamma_{1}(t)} u_{n x}(x, t) d x, \quad 0 \leq t \leq \tau, \quad\left(\gamma_{1}(t), t\right) \in \gamma_{1, \tau} . \tag{2.24}
\end{gather*}
$$

This, together with the Schwarz inequality, implies that

$$
\begin{align*}
\left|u_{n}(x, \tau)\right|^{2} & \leq \int_{\gamma_{2}^{-1}(x)}^{\tau} 1^{2} d t \int_{\gamma_{2}^{-1}(x)}^{\tau}\left|u_{n t}(x, t)\right|^{2} d t \leq T \int_{\gamma_{2}^{-1}(x)}^{\tau}\left|u_{n t}(x, t)\right|^{2} d t, \quad(x, \tau) \in \Omega_{\tau},  \tag{2.25}\\
\left|u_{n}\left(\gamma_{1}(t), t\right)\right|^{2} & \leq \int_{\gamma_{2}(t)}^{\gamma_{1}(t)} 1^{2} d x \int_{\gamma_{2}(t)}^{\gamma_{1}(t)}\left|u_{n x}(x, t)\right|^{2} d x \\
& \leq T \int_{\gamma_{2}(t)}^{\gamma_{1}(t)}\left|u_{n x}(x, t)\right|^{2} d x, \quad 0 \leq t \leq \tau, \quad\left(\gamma_{1}(t), t\right) \in \gamma_{1, \tau} . \tag{2.26}
\end{align*}
$$

By integrating both sides of inequality (2.25) with respect to $x$ on the closed interval $\left[\gamma_{2}(\tau), \gamma_{1}(\tau)\right]$, we obtain

$$
\begin{align*}
\int_{\Omega_{\tau}}\left|u_{n}(x, \tau)\right|^{2} d x & \leq T \int_{\gamma_{2}(\tau)}^{\gamma_{1}(\tau)}\left[\int_{\gamma_{2}^{-1}(x)}^{\tau}\left|u_{n t}(x, t)\right|^{2} d t\right] d x=T \int_{D_{\tau} \cap\left\{x<\gamma_{1}(\tau)\right\}}\left|u_{n t}(x, t)\right|^{2} d x d t \\
& \leq T \int_{D_{\tau}}\left|u_{n t}(x, t)\right|^{2} d x d t . \tag{2.27}
\end{align*}
$$

In a similar way, since $\left|\gamma_{1}^{\prime}(t)\right|<1, t \geq 0$, it follows that, by integrating both sides of inequality (2.26) with respect to $t$ over the interval $[0, \tau]$, we obtain

$$
\begin{align*}
\int_{\gamma_{1}, \tau}\left|u_{n}(x, t)\right|^{2} d s & =\int_{0}^{\tau}\left|u_{n}\left(\gamma_{1}(t), t\right)\right|^{2} \sqrt{1+\left|\gamma_{1}^{\prime}(t)\right|^{2}} d t \leq \sqrt{2} \int_{0}^{\tau}\left|u_{n}\left(\gamma_{1}(t), t\right)\right|^{2} d t \\
& \leq \sqrt{2} T \int_{0}^{\tau}\left[\int_{\gamma_{2}(t)}^{\gamma_{1}(t)}\left|u_{n x}(x, t)\right|^{2} d x\right] d t=\sqrt{2} T \int_{D_{\tau}}\left|u_{n x}(x, t)\right|^{2} d x d t \tag{2.28}
\end{align*}
$$

From inequality (2.27), we have

$$
\begin{equation*}
\int_{D_{\tau}} u_{n}^{2} d x d t=\int_{0}^{\tau}\left[\int_{\Omega_{\sigma}} u_{n}^{2} d x\right] d \sigma \leq T \int_{0}^{\tau}\left[\int_{D_{\sigma}} u_{n t}^{2} d x d t\right] d \sigma \leq T^{2} \int_{D_{\tau}} u_{n t}^{2} d x d t \tag{2.29}
\end{equation*}
$$

The relations $2 F_{n} u_{n t} \leq F_{n}^{2}+u_{n t}^{2},(2.22)$, and (2.27)-(2.29), together with inequality (2.21), imply the estimate

$$
\begin{equation*}
w_{n}(\tau) \leq M_{5}+M_{6} \int_{D_{\tau}}\left(u_{n t}^{2}+u_{n x}^{2}\right) d x d t+\int_{D_{\tau}} F_{n}^{2} d x d t \tag{2.30}
\end{equation*}
$$

Here

$$
\begin{equation*}
M_{5}:=2(1+\sqrt{2}) T M_{1}+2 T^{2} M_{3}, \quad M_{6}:=2(1+\sqrt{2}) M_{2} T+2 M_{4} T^{2}+1 \tag{2.31}
\end{equation*}
$$

Since

$$
\int_{D_{\tau}}\left(u_{n t}^{2}+u_{n x}^{2}\right) d x d t=\int_{0}^{\tau} w_{n}(\sigma) d \sigma
$$

it follows from inequality (2.30) that

$$
\begin{equation*}
w_{n}(\tau) \leq M_{6} \int_{0}^{\tau} w_{n}(\sigma) d \sigma+M_{5}+T^{2}\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}^{2}, \quad 0<\tau \leq T . \tag{2.32}
\end{equation*}
$$

This, together with the Gronwall lemma, implies that

$$
\begin{equation*}
w_{n}(\tau) \leq\left(M_{5}+T^{2}\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}^{2}\right) \exp \left(M_{6} \tau\right), \quad 0<\tau \leq T . \tag{2.33}
\end{equation*}
$$

Next, by virtue of condition (2.10), for any $(x, t) \in \bar{D}_{T} \backslash O$ we have

$$
u_{n}(x, t)=u_{n}(x, t)-u_{n}\left(\gamma_{2}(t), t\right)=\int_{\gamma_{2}(t)}^{x} u_{n x}(\xi, t) d \xi,
$$

whence, by analogy with the derivation of inequality (2.26), one can show that

$$
\begin{equation*}
\left|u_{n}(x, t)\right|^{2} \leq T \int_{\gamma_{2}(t)}^{x}\left|u_{n x}(\xi, t)\right|^{2} d \xi, \quad(x, t) \in \bar{D}_{T} \backslash O \tag{2.34}
\end{equation*}
$$

Inequality (2.34), together with the estimate (2.33) and the definition of the quantity $w_{n}$ as the left-hand side in relation (2.21), implies that

$$
\begin{equation*}
\left|u_{n}(x, t)\right|^{2} \leq T \int_{\Omega_{t}} u_{n x}^{2} d x \leq T w_{n}(t) \leq T\left(M_{5}+T^{2}\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}^{2}\right) \exp \left(M_{6} t\right), \quad(x, t) \in \bar{D}_{T} \backslash O . \tag{2.35}
\end{equation*}
$$

By taking into account the estimate (2.35) and by using the obvious inequality

$$
\left(\sum_{i=1}^{m} a_{i}^{2}\right)^{1 / 2} \leq \sum_{i=1}^{m}\left|a_{i}\right|
$$

we obtain

$$
\begin{equation*}
\left\|u_{n}\right\|_{C\left(\bar{D}_{T}\right)} \leq c_{1}\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}+c_{2} \tag{2.36}
\end{equation*}
$$

where

$$
\begin{equation*}
c_{1}=T^{3 / 2} \exp \left(2^{-1} M_{6} T\right), \quad c_{2}=\left(T M_{5}\right)^{1 / 2} \exp \left(2^{-1} M_{6} T\right), \tag{2.37}
\end{equation*}
$$

and $M_{5}$ and $M_{6}$ are the constants defined in (2.31). By taking into account relations (2.7) and (2.11) and by passing in inequality (2.36) to the limit as $n \rightarrow \infty$, we obtain the a priori estimate (2.6). The proof of Lemma 2.1 is complete.

Remark 2.3. In the linear case in which the function $f$ occurring in Eq. (1.1) vanishes, one can introduce the notion of a strong generalized solution of problem (1.1)-(1.3) in a similar way. In this case, by virtue of relation (2.1), the function $g$ vanishes and satisfies conditions (2.2) and (2.3) for $M_{i}=0,1 \leq i \leq 4$; moreover, under conditions (1.5), (1.6), and (2.4), the a priori estimate (2.6) is valid and, by virtue of relations (2.31) and (2.37), acquires the form

$$
\begin{equation*}
\|u\|_{C\left(\bar{D}_{T}\right)} \leq T^{3 / 2} \exp \left(2^{-1} T\right)\|F\|_{C\left(\bar{D}_{T}\right)} . \tag{2.38}
\end{equation*}
$$

## 3. CASES OF GLOBAL SOLVABILITY OF PROBLEM (1.1)-(1.3) <br> IN THE CLASS $C$

In the new independent variables $\xi=t+x, \eta=t-x$, the domain $D_{T}$ becomes a curvilinear triangular domain $G_{T}$ with vertices at the points $O(0,0), Q_{1}\left(T+\gamma_{1}(T), T-\gamma_{1}(T)\right)$, and $Q_{2}\left(T+\gamma_{2}(T), T-\gamma_{2}(T)\right)$ of the plane of the variables $\xi$ and $\eta$, and problem (1.1)-(1.3) becomes the problem

$$
\begin{align*}
\widetilde{L} \widetilde{u}:=\widetilde{u}_{\xi_{\eta}}+\widetilde{f}(\xi, \eta, \widetilde{u}) & =\widetilde{F}(\xi, \eta), \quad(\xi, \eta) \in G_{T},  \tag{3.1}\\
\left(m_{1} \widetilde{u}_{\xi}+m_{2} \widetilde{u}_{\eta}\right) \mid \widetilde{\gamma}_{1, T} & =0,  \tag{3.2}\\
\left.\widetilde{u}\right|_{\tilde{\gamma}_{2}, T} & =0 \tag{3.3}
\end{align*}
$$

for the unknown function

$$
\widetilde{u}(\xi, \eta):=u\left(\frac{\xi-\eta}{2}, \frac{\xi+\eta}{2}\right) .
$$

Here

$$
\begin{align*}
\widetilde{f}(\xi, \eta, \widetilde{u}) & :=\frac{1}{4} f\left(\frac{\xi-\eta}{2}, \frac{\xi+\eta}{2}, \widetilde{u}\right), \quad \widetilde{F}(\xi, \eta):=\frac{1}{4} F\left(\frac{\xi-\eta}{2}, \frac{\xi+\eta}{2}\right), \\
m_{1} & :=l_{2}+l_{1}, \quad m_{2}:=l_{2}-l_{1} \quad \text { on } \quad \widetilde{\gamma}_{1, T}, \tag{3.4}
\end{align*}
$$

and $\widetilde{\gamma}_{1, T}$ and $\widetilde{\gamma}_{2, T}$ are the images of the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ under that transformation issuing from the common point $O(0,0)$ with terminal points $Q_{1}$ and $Q_{2}$.

By analogy with Definition 1.1, one can introduce the notion of strong generalized solution $\widetilde{u}$ of problem (3.1)-(3.3) in the class $C$ in the domain $G_{T}$.

By virtue of conditions (1.5) and (1.6), the smooth noncharacteristic curves $\widetilde{\gamma}_{1, T}$ and $\widetilde{\gamma}_{2, T}$ can be represented in the form

$$
\begin{equation*}
\widetilde{\gamma}_{1, T}: \eta=\lambda_{1}(\xi), \quad 0 \leq \xi \leq \xi_{0} ; \quad \widetilde{\gamma}_{2, T}: \xi=\lambda_{2}(\eta), \quad 0 \leq \eta \leq \eta_{0} \tag{3.5}
\end{equation*}
$$

where $\xi_{0}:=T+\gamma_{1}(T)<\eta_{0}:=T-\gamma_{2}(T)$ and

$$
\begin{align*}
& \lambda_{1}^{\prime}(\xi)>0, \quad 0 \leq \xi \leq \xi_{0} ; \quad \lambda_{2}^{\prime}(\eta)>0, \quad 0 \leq \lambda_{2}(\eta) \leq \eta, \quad 0 \leq \eta \leq \eta_{0} ;  \tag{3.6}\\
& \lambda_{2}\left(\lambda_{1}(\xi)\right)<\xi, \quad 0<\xi \leq \xi_{0} ; \quad \lambda_{1}\left(\lambda_{2}(\eta)\right)<\eta, \quad 0<\eta \leq \eta_{0} ;  \tag{3.7}\\
& G_{T}:=\left\{(\xi, \eta) \in\left(0, \xi_{0}\right) \times\left(0, \eta_{0}\right): \lambda_{1}(\xi)<\eta, \quad \lambda_{2}(\eta)<\xi, \xi+\eta<2 T\right\} . \tag{3.8}
\end{align*}
$$

Remark 3.1. Obviously, $u=u(x, t)$ is a strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$ if and only if $\widetilde{u}$ is a strong generalized solution of problem (3.1)-(3.3) in the class $C$ in the domain $G_{T}$; moreover, under the assumptions of Lemma 2.1 this solution $\widetilde{u}$ satisfies an a priori estimate of the type (2.6),

$$
\begin{equation*}
\|\widetilde{u}\|_{C\left(\bar{G}_{T}\right)}=\|u\|_{C\left(\bar{D}_{T}\right)} \leq c_{1}\|F\|_{C\left(\bar{D}_{T}\right)}+c_{2} \leq 4 c_{1}\|\widetilde{F}\|_{C\left(\bar{G}_{T}\right)}+c_{2} \tag{3.9}
\end{equation*}
$$

with the same constants $c_{1}$ and $c_{2}$.
Further, we first consider the linear case of problem (3.1)-(3.3) for which the function $\tilde{f}$ occurring in Eq. (3.1) vanishes,

$$
\begin{align*}
\widetilde{\square} \widetilde{w}:=\widetilde{w}_{\xi \eta} & =\widetilde{F}(\xi, \eta), \quad(\xi, \eta) \in G_{T},  \tag{3.10}\\
\left.\left(m_{1} \widetilde{w}_{\xi}+m_{2} \widetilde{w}_{\eta}\right)\right|_{\tilde{\gamma}_{1, T}} & =0,  \tag{3.11}\\
\widetilde{w}_{\tilde{\gamma}_{2, T}} & =0 . \tag{3.12}
\end{align*}
$$

Remark 3.2. By Remarks 2.3 and 3.1, a strong generalized solution $\widetilde{w}$ of the linear problem (3.10)-(3.12) in the class $C$ in the domain $G_{T}$ satisfies the estimate

$$
\begin{equation*}
\|\widetilde{w}\|_{C\left(\bar{G}_{T}\right)} \leq 4 T^{3 / 2} \exp \left(2^{-1} T\right)\|\widetilde{F}\|_{C\left(\bar{G}_{T}\right)} . \tag{3.13}
\end{equation*}
$$

In particular, the estimate (3.13) holds for a classical solution $\widetilde{w} \in C^{2}\left(\bar{G}_{T}\right)$ of that problem. The estimate (3.13) implies the uniqueness of both generalized and classical solutions of that problem.

Remark 3.3. It follows from the condition $\left.\left(\left|l_{1}\right|+\left|l_{2}\right|\right)\right|_{\gamma_{1}} \neq 0$ that, by virtue of relations (3.4), at each point $P \in \widetilde{\gamma}_{1, T}$ at least one of the numbers $m_{1}(P)$ and $m_{2}(P)$ is nonzero. In what follows, we assume that $\left.m_{1}\right|_{\gamma_{1}} \neq 0$; i.e.,

$$
\begin{equation*}
\left(l_{2}+l_{1}\right)(P) \neq 0, \quad P \in \gamma_{1, T} . \tag{3.14}
\end{equation*}
$$

Condition (3.14) implies that the direction $\left(l_{1}, l_{2}\right)$ is not a characteristic direction corresponding to the family of characteristics $x+t=$ const of Eq. (1.1).

Set

$$
\begin{equation*}
a(\xi):=\frac{m_{2}(\xi)}{m_{1}(\xi)} \lambda_{2}^{\prime}\left(\lambda_{1}(\xi)\right), \quad 0 \leq \xi \leq \xi_{0} \tag{3.15}
\end{equation*}
$$

and consider the equation

$$
\begin{equation*}
|a(0)|=\left|\frac{m_{2}(0)}{m_{1}(0)} \lambda_{2}^{\prime}(0)\right|<1 . \tag{3.16}
\end{equation*}
$$

Lemma 3.1. Let conditions (2.4) and (3.14) be satisfied at the point $P=O(0,0)$. If either $\left(l_{1} l_{2}\right)(O) \neq 0$ or $\left(l_{1} l_{2}\right)(O)=0$ but the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are not tangent to each other at the point $O$ or are tangent but $\gamma_{2}^{\prime}(0)<0$, then condition (3.16) is also satisfied.

Proof. By virtue of condition (3.6), we have

$$
\begin{equation*}
0<\lambda_{2}^{\prime}(0) \leq 1 . \tag{3.17}
\end{equation*}
$$

If $\left(l_{1} l_{2}\right)(O)>0$, then, obviously, $\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right|<1$; therefore, by virtue of relations (3.4) and (3.17), inequality (3.16) is satisfied.

It follows from inequalities (2.4) and (2.5) at the point $P=O(0,0)$ that

$$
\begin{equation*}
\gamma_{1}^{\prime}(0) \leq \frac{2 l_{1} l_{2}}{l_{1}^{2}+l_{2}^{2}}(O) \tag{3.18}
\end{equation*}
$$

By virtue of the representations (3.5), one can readily show that

$$
\begin{equation*}
\lambda_{1}^{\prime}(0)=\frac{1-\gamma_{1}^{\prime}(0)}{1+\gamma_{1}^{\prime}(0)}, \quad \lambda_{2}^{\prime}(0)=\frac{1+\gamma_{2}^{\prime}(0)}{1-\gamma_{2}^{\prime}(0)} . \tag{3.19}
\end{equation*}
$$

Next, by virtue of conditions (3.6) and (3.7), we have $0<\lambda_{1}^{\prime}(0) \lambda_{2}^{\prime}(0) \leq 1$, because $\left[\lambda_{2}\left(\lambda_{1}(\xi)\right)\right]^{\prime}(0)=$ $\left[\lambda_{1}\left(\lambda_{2}(\eta)\right)\right]^{\prime}(0)=\lambda_{1}^{\prime}(0) \lambda_{2}^{\prime}(0)$. Therefore,

$$
\begin{equation*}
\lambda_{2}^{\prime}(0) \leq \frac{1}{\lambda_{1}^{\prime}(0)} \tag{3.20}
\end{equation*}
$$

For $\left(l_{1} l_{2}\right)(O)<0$, one can readily see that

$$
\begin{equation*}
\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right|>1 \tag{3.21}
\end{equation*}
$$

but nevertheless, as is shown below, inequality (3.16) remains valid.

Now, in view of the fact that $\mu(s):=(1+s) /(1-s), s \in \mathbb{R}$, is an increasing function and by taking into account relations (3.4) and (3.18)-(3.21), we obtain the estimate

$$
\begin{aligned}
|a(0)| & =\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right| \lambda_{2}^{\prime}(0) \leq\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right| \frac{1+\gamma_{1}^{\prime}(0)}{1-\gamma_{1}^{\prime}(0)} \\
& \leq\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right| \frac{1+\frac{2 l_{1} l_{2}}{l_{1}^{\prime}+l_{2}^{\prime}}(O)}{1-\frac{l_{1} 2}{l_{1}^{\prime}+l_{2}^{\prime}}(O)}=\left|\left(\frac{l_{2}+l_{1}}{l_{2}-l_{1}}\right)(O)\right|<1 .
\end{aligned}
$$

It remains to consider the case in which $\left(l_{1} l_{2}\right)(O)=0$. By virtue of inequalities (1.6), we have $\gamma_{2}^{\prime}(0) \leq \gamma_{1}^{\prime}(0) \leq 0$. Therefore, if the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are not tangent at the point $O(0,0)$, then $\gamma_{2}^{\prime}(0) \neq \gamma_{1}^{\prime}(0) \leq 0$ and hence $\gamma_{2}^{\prime}(0)<0$. This, together with relations (3.19), implies that

$$
\begin{equation*}
\lambda_{2}^{\prime}(0)<1 . \tag{3.22}
\end{equation*}
$$

Since the relation $\left(l_{1} l_{2}\right)(O)=0$ implies that $\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right|=1$, it follows from the estimate (3.22) that

$$
|a(0)|=\left|\left(\frac{l_{2}-l_{1}}{l_{2}+l_{1}}\right)(O)\right| \lambda_{2}^{\prime}(0)=\lambda_{2}^{\prime}(0)<1 .
$$

In a similar way, one can consider the case in which $\left(l_{1} l_{2}\right)(O)=0$, the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are tangent at the point $O$, but $\gamma_{2}^{\prime}(0)<0$. The proof of Lemma 3.1 is complete.

Remark 3.4. One can readily see that if $\left(l_{1} l_{2}\right)(O)=0$, then the relation $|a(0)|=1$ holds if and only if $\gamma_{2}^{\prime}(0)=0$; in this case, by virtue of conditions (1.6), the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are tangent at the common point $O$.

Let $G_{0, T}:=\left\{(\xi, \eta) \in \mathbb{R}^{2}: 0<\xi<\xi_{0}, 0<\eta<\eta_{0}\right\}$ be the characteristic rectangle in the plane of the variables $\xi$ and $\eta$ corresponding to Eq. (3.10). By virtue of (3.8), we have $G_{T} \subset G_{0, T}$. If $\widetilde{F}$ belongs to $C\left(\bar{G}_{T}\right)$, then we extend that function as a continuous function into the closed domain $\bar{G}_{0, T}$ and keep the previous notation for it by setting, for example, $\widetilde{F}(\xi, \eta)=\widetilde{F}\left(\xi, \lambda_{1}(\xi)\right)$ for $0 \leq \eta \leq \lambda_{1}(\xi), 0 \leq \xi \leq \xi_{0}, \widetilde{F}(\xi, \eta)=\widetilde{F}\left(\lambda_{2}(\eta), \eta\right)$ for $0 \leq \xi \leq \lambda_{2}(\eta), 0 \leq \eta \leq \eta_{0}$, and $\widetilde{F}(\xi, \eta)=\widetilde{F}(2 T-\eta, \eta)$ for $(\xi, \eta) \in G_{0, T} \cap\{\xi+\eta \geq 2 T\}$. Since the space $C^{1}\left(\bar{G}_{0, T}\right)$ is dense in the space $C\left(\bar{G}_{0, T}\right)[15$, p. 37 of the Russian translation], it follows that there exists a function sequence $\widetilde{F}_{n}$ such that

$$
\begin{equation*}
\widetilde{F}_{n} \in C^{1}\left(\bar{G}_{0, T}\right), \quad \lim _{n \rightarrow \infty}\left\|\widetilde{F}_{n}-\widetilde{F}\right\|_{C\left(\bar{G}_{0, T}\right)}=0 . \tag{3.23}
\end{equation*}
$$

We introduce the function $\widetilde{u}_{n} \in C^{2}\left(\bar{G}_{0, T}\right)$ that is the solution of the Goursat problem

$$
\begin{array}{rlrl}
\widetilde{\square} \widetilde{u}_{n} & =\widetilde{F}_{n}(\xi, \eta), & (\xi, \eta) \in G_{0, T}, & \\
\widetilde{u}_{n}(\xi, 0) & =\varphi_{n}(\xi), \quad 0 \leq \xi \leq \xi_{0} ; \quad \widetilde{u}_{n}(0, \eta)=\psi_{n}(\eta), \quad 0 \leq \eta \leq \eta_{0},
\end{array}
$$

where $\varphi_{n} \in C^{2}\left(\left[0, \xi_{0}\right]\right)$ and $\psi_{n} \in C^{2}\left(\left[0, \eta_{0}\right]\right)$ are some functions satisfying the matching condition

$$
\begin{equation*}
\varphi_{n}(0)=\psi_{n}(0)=0 . \tag{3.24}
\end{equation*}
$$

It is well known that the unique solution of this problem can be represented in the form [16, p. 246]

$$
\begin{equation*}
\widetilde{u}_{n}(\xi, \eta)=\varphi_{n}(\xi)+\psi_{n}(\eta)+\int_{0}^{\xi} d \xi^{\prime} \int_{0}^{\eta} \widetilde{F}_{n}\left(\xi^{\prime}, \eta^{\prime}\right) d \eta^{\prime}, \quad(\xi, \eta) \in \bar{G}_{0, T} \tag{3.25}
\end{equation*}
$$

By assuming that

$$
\begin{equation*}
\gamma_{i} \in C^{i}([0, T]), \quad l_{i} \in C^{1}\left(\gamma_{1, T}\right), \quad i=1,2, \tag{3.26}
\end{equation*}
$$

we readily obtain

$$
\begin{equation*}
\lambda_{1} \in C^{1}\left(\left[0, \xi_{0}\right]\right), \quad \lambda_{2} \in C^{2}\left(\left[0, \eta_{0}\right]\right), \quad m_{i} \in C^{1}\left(\widetilde{\gamma}_{1, T}\right), \quad i=1,2 . \tag{3.27}
\end{equation*}
$$

Now we construct functions $\varphi_{n} \in C^{2}\left(\left[0, \xi_{0}\right]\right)$ and $\psi_{n} \in C^{2}\left(\left[0, \eta_{0}\right]\right)$ such that the function $\widetilde{w}=\widetilde{u}_{n}$ defined by relation (3.25) satisfies the boundary conditions (3.11) and (3.12). By differentiating relation (3.12) in the direction of the tangent to $\widetilde{\gamma}_{2, T}$ with regard of (3.5), we obtain

$$
\begin{equation*}
\lambda_{2}^{\prime}(\eta) \widetilde{u}_{n \xi}\left(\lambda_{2}(\eta), \eta\right)+\widetilde{u}_{n \eta}\left(\lambda_{2}(\eta), \eta\right)=0, \quad 0 \leq \eta \leq \eta_{0} . \tag{3.28}
\end{equation*}
$$

Obviously, relation (3.28), together with the condition $\widetilde{u}_{n}(0,0)=0$, is equivalent to condition (3.12). By substituting the expression for $\widetilde{u}_{n}$ in (3.25) into relations (3.11) and (3.28) and by using the representations (3.5), for the functions $\varphi_{n}^{\prime}$ and $\psi_{n}^{\prime}$ we obtain the system of functional equations

$$
\begin{align*}
m_{1}(\xi) \varphi_{n}^{\prime}(\xi)+m_{2}(\xi) \psi_{n}^{\prime}\left(\lambda_{1}(\xi)\right) & =\omega_{1 n}(\xi), & & 0 \leq \xi \leq \xi_{0}  \tag{3.29}\\
\lambda_{2}^{\prime}(\eta) \varphi_{n}^{\prime}\left(\lambda_{2}(\eta)\right)+\psi_{n}^{\prime}(\eta) & =\omega_{2 n}(\eta), & & 0 \leq \eta \leq \eta_{0} . \tag{3.30}
\end{align*}
$$

Here

$$
\begin{align*}
& \omega_{1 n}(\xi):=-m_{1}(\xi) \int_{0}^{\lambda_{1}(\xi)} \widetilde{F}_{n}\left(\xi, \eta^{\prime}\right) d \eta^{\prime}-m_{2}(\xi) \int_{0}^{\xi} \widetilde{F}_{n}\left(\xi^{\prime}, \lambda_{1}(\xi)\right) d \xi^{\prime}, \quad 0 \leq \xi \leq \xi_{0},  \tag{3.31}\\
& \omega_{2 n}(\eta):=-\lambda_{2}^{\prime}(\eta) \int_{0}^{\eta} \widetilde{F}_{n}\left(\lambda_{2}(\eta), \eta^{\prime}\right) d \eta^{\prime}-\int_{0}^{\lambda_{2}(\eta)} \widetilde{F}_{n}\left(\xi^{\prime}, \eta\right) d \xi^{\prime}, \quad 0 \leq \eta \leq \eta_{0} . \tag{3.32}
\end{align*}
$$

If condition (3.14) is satisfied, which is equivalent to the condition $\left.m_{1}\right|_{\tilde{\gamma}_{1, T}} \neq 0$, then, by eliminating the function $\psi_{n}^{\prime}$ from the system of equations (3.29) and (3.30), for $\varphi_{0 n}:=\varphi_{n}^{\prime}$, we obtain the functional equation

$$
\begin{equation*}
\varphi_{0 n}(\xi)-a(\xi) \varphi_{0 n}\left(\lambda_{2}\left(\lambda_{1}(\xi)\right)\right)=\omega_{n}(\xi), \quad 0 \leq \xi \leq \xi_{0} . \tag{3.33}
\end{equation*}
$$

Here $a(\xi), 0 \leq \xi \leq \xi_{0}$, is the function defined by relation (3.15), and

$$
\begin{equation*}
\omega_{n}(\xi):=\frac{1}{m_{1}(\xi)}\left[\omega_{1 n}(\xi)-m_{2}(\xi) \omega_{2 n}\left(\lambda_{1}(\xi)\right)\right], \quad 0 \leq \xi \leq \xi_{0} \tag{3.34}
\end{equation*}
$$

By setting

$$
\begin{equation*}
\tau(\xi):=\lambda_{2}\left(\lambda_{1}(\xi)\right), \quad 0 \leq \xi \leq \xi_{0} \tag{3.35}
\end{equation*}
$$

and by taking into account relations (3.7) and (3.27), we obtain

$$
\begin{equation*}
\tau \in C^{1}\left(\left[0, \xi_{0}\right]\right), \quad \tau(0)=0, \quad \tau(\xi)<\xi \quad \text { if } \quad 0<\xi \leq \xi_{0} \tag{3.36}
\end{equation*}
$$

Since $a \in C\left(\left[0, \xi_{0}\right]\right)$, it follows that, under condition (3.16), there exists a positive number $\varepsilon$ such that

$$
\begin{equation*}
|a(\xi)| \leq q:=\text { const }<1 \quad \text { if } \quad 0 \leq \xi \leq \varepsilon . \tag{3.37}
\end{equation*}
$$

From relations (3.36), we find that if $\tau_{k}(\xi):=\tau\left(\tau_{k-1}(\xi)\right)$ and $\tau_{0}(\xi):=\xi, 0 \leq \xi \leq \xi_{0}$, then the function sequence $\left\{\tau_{k}(\xi)\right\}_{k=1}^{\infty}$ converges uniformly to zero on the interval $\left[0, \xi_{0}\right]$; i.e., there exists a positive integer $n_{0}=n_{0}(\varepsilon)$ such that

$$
\begin{equation*}
\tau_{k}(\xi) \leq \varepsilon, \quad 0 \leq \xi \leq \xi_{0}, \quad k \geq n_{0} \tag{3.38}
\end{equation*}
$$

By $\Lambda: C\left(\left[0, \xi_{0}\right]\right) \rightarrow C\left(\left[0, \xi_{0}\right]\right)$ we denote the linear continuous operator acting by the rule

$$
\begin{equation*}
\left(\Lambda \omega_{n}\right)(\xi):=a(\xi) \omega_{n}(\tau(\xi)), \quad 0 \leq \xi \leq \xi_{0} \tag{3.39}
\end{equation*}
$$

Obviously,

$$
\begin{equation*}
\left(\Lambda^{k} \omega_{n}\right)(\xi)=a(\xi) a(\tau(\xi)) \cdots a\left(\tau_{k-1}(\xi)\right) \omega_{n}\left(\tau_{k}(\xi)\right), \quad k \geq 2, \tag{3.40}
\end{equation*}
$$

and for $k=1$ and $k=0$, we set

$$
\begin{equation*}
\Lambda^{1}=\Lambda \quad \text { and } \quad \Lambda^{0}=I \tag{3.41}
\end{equation*}
$$

where $I$ is the identity operator.
By virtue of relations (3.36)-(3.41), we have the estimate

$$
\begin{aligned}
\left|\left(\Lambda^{k} \omega_{n}\right)(\xi)\right| & \leq\left[a(\xi) a(\tau(\xi)) \cdots a\left(\tau_{n_{0}-1}(\xi)\right)\right]\left[a\left(\tau_{n_{0}}(\xi)\right) \cdots a\left(\tau_{k-1}(\xi)\right)\right] \omega_{n}\left(\tau_{k}(\xi)\right) \\
& \leq\|a\|_{C([0, \xi])}^{n_{0}} q^{k-n_{0}}\left\|\omega_{n}\right\|_{C(0, \xi])}, \quad 0 \leq \xi \leq \xi_{0}, \quad k>n_{0},
\end{aligned}
$$

whence we obtain

$$
\begin{equation*}
\left\|\Lambda^{k}\right\|_{C\left(\left[0, \xi_{0}\right]\right) \rightarrow C\left(\left[0, \xi_{0}\right]\right)} \leq M_{0} q^{k}, \quad k>n_{0} \tag{3.42}
\end{equation*}
$$

where

$$
M_{0}:=\left(q^{-1}\|a\|_{C\left(\left[0, \xi_{0}\right]\right)}\right)^{n_{0}} .
$$

It follows from inequality (3.42), where $q<1$, that if condition (3.16) is satisfied, then the Neumann series

$$
(I-\Lambda)^{-1}=\sum_{k=0}^{\infty} \Lambda^{k}
$$

of the operator $\Lambda$ is convergent in the space $C\left(\left[0, \xi_{0}\right]\right)$, and by (3.35), the unique solution $\varphi_{0 n} \in$ $C\left(\left[0, \xi_{0}\right]\right)$ of Eq. (33) can be represented in the form

$$
\begin{equation*}
\varphi_{0 n}(\xi)=\left[\sum_{k=0}^{\infty} \Lambda^{k} \omega_{n}\right](\xi), \quad 0 \leq \xi \leq \xi_{0} \tag{3.43}
\end{equation*}
$$

Remark 3.5. Note that, by virtue of Remark 3.4, if, in the case $\left(l_{1} l_{2}\right)(O)=0$, we have $\gamma_{2}^{\prime}(0)=0$, which is equivalent to the condition $\lambda_{2}^{\prime}(0)=1$, then the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are tangent at the point $O$; moreover, $|a(0)|=1, \tau^{\prime}(0)=\lambda_{2}^{\prime}(0) \lambda_{1}^{\prime}(0)=1$, and Eq. (3.33) is not solvable in the class $C\left(\left[0, \xi_{0}\right]\right)$ for any right-hand side $\omega_{n} \in C\left(\left[0, \xi_{0}\right]\right)$. In this case, a necessary and sufficient condition for the solvability of Eq. (3.33) in the class $C\left(\left[0, \xi_{0}\right]\right)$ is given by the uniform convergence of the series on the right-hand side in relation (3.43) on the interval $\left[0, \xi_{0}\right]$, which is not necessarily true for any function $\omega_{n} \in C\left(\left[0, \xi_{0}\right]\right)$.

Remark 3.6. One can readily see that if we additionally require that the functions $a, \tau$, and $\omega_{n}$ belong to $C^{1}\left(\left[0, \xi_{0}\right]\right)$, then the solution $\varphi_{0 n}$ of Eq. (3.33), which can be represented as the convergent series (3.43) in $C\left(\left[0, \xi_{0}\right]\right)$, belongs to the space $C^{1}\left(\left[0, \xi_{0}\right]\right)$ as well; moreover, its derivative $\chi_{n}:=\varphi_{0 n}^{\prime}$ can be found from the functional equation

$$
\begin{equation*}
\chi_{n}(\xi)-a_{1}(\xi) \chi_{n}(\tau(\xi))=\widetilde{\omega}_{1 n}(\xi), \quad 0 \leq \xi \leq \xi_{0} \tag{3.44}
\end{equation*}
$$

where $a_{1}(\xi):=a(\xi) \tau^{\prime}(\xi)$ and $\widetilde{\omega}_{1 n}(\xi):=\omega_{n}^{\prime}(\xi)+a^{\prime}(\xi) \varphi_{0 n}(\tau(\xi)), 0 \leq \xi \leq \xi_{0}$, and since $\left|\tau^{\prime}(0)\right| \leq 1$ by virtue of relations (3.36), we have $\left|a_{1}(0)\right|<1$ under condition (3.16); consequently, by analogy with (3.43), the solution $\chi_{n}$ of Eq. (3.44) can be represented in the form

$$
\begin{equation*}
\chi_{n}=\sum_{k=0}^{\infty} \Lambda_{1}^{k} \widetilde{\omega}_{1 n}, \tag{3.45}
\end{equation*}
$$

where $\left(\Lambda_{1} \widetilde{\omega}_{1 n}\right)(\xi):=a_{1}(\xi) \widetilde{\omega}_{1 n}(\tau(\xi)), 0 \leq \xi \leq \xi_{0}$. By setting

$$
\begin{equation*}
\widetilde{\varphi}_{0 n}(\xi):=\int_{0}^{\xi} \chi_{n}\left(\xi^{\prime}\right) d \xi^{\prime}+\varphi_{0 n}(0), \quad 0 \leq \xi \leq \xi_{0}, \tag{3.46}
\end{equation*}
$$

and by integrating Eq. (3.44), we obtain

$$
\widetilde{\varphi}_{0 n}(\xi)-\varphi_{0 n}(0)-\int_{0}^{\xi} a\left(\xi^{\prime}\right) d \widetilde{\varphi}_{0 n}\left(\tau\left(\xi^{\prime}\right)\right)=\int_{0}^{\xi} a^{\prime}\left(\xi^{\prime}\right) \varphi_{0 n}\left(\tau\left(\xi^{\prime}\right)\right) d \xi^{\prime}+\omega_{n}(\xi)-\omega_{n}(0), \quad 0 \leq \xi \leq \xi_{0}
$$

By integrating the third term on the left-hand side in the last relation, we obtain

$$
\begin{aligned}
\widetilde{\varphi}_{0 n}(\xi)-\varphi_{0 n}(0)- & a(\xi) \widetilde{\varphi}_{0 n}(\tau(\xi))+a(0) \widetilde{\varphi}_{0 n}(\tau(0))+\int_{0}^{\xi} a^{\prime}\left(\xi^{\prime}\right) \widetilde{\varphi}_{0 n}\left(\tau\left(\xi^{\prime}\right)\right) d \xi^{\prime} \\
& =\int_{0}^{\xi} a^{\prime}\left(\xi^{\prime}\right) \varphi_{0 n}\left(\tau\left(\xi^{\prime}\right)\right) d \xi^{\prime}+\omega_{n}(\xi)-\omega_{n}(0), \quad 0 \leq \xi \leq \xi_{0}
\end{aligned}
$$

By using relation (3.35), by subtracting relation (3.33) from the last relation, and by taking into account the equalities $\tau(0)=0$ and $\widetilde{\varphi}_{0 n}(0)=\varphi_{0 n}(0)$ in view of relations (3.36) and (3.46), for $\psi_{0 n}:=\widetilde{\varphi}_{0 n}-\varphi_{0 n}$, we obtain the Volterra integro-functional equation

$$
\psi_{0 n}(\xi)-a(\xi) \psi_{0 n}(\tau(\xi))+\int_{0}^{\xi} a^{\prime}\left(\xi^{\prime}\right) \psi_{0 n}\left(\tau\left(\xi^{\prime}\right)\right) d \xi^{\prime}=0, \quad 0 \leq \xi \leq \xi_{0}
$$

By applying the standard successive approximation method [4] to that equation, we obtain $\psi_{0 n}=0 ;$ i.e., $\widetilde{\varphi}_{0 n}=\varphi_{0 n}$, and therefore,

$$
\varphi_{0 n}(\xi)=\int_{0}^{\xi} \chi_{n}\left(\xi^{\prime}\right) d \xi^{\prime}+\varphi_{0 n}(0), \quad 0 \leq \xi \leq \xi_{0}
$$

taking into account the representation (3.46). Hence it follows that $\varphi_{0 n}$ belongs to $C^{1}\left(\left[0, \xi_{0}\right]\right)$. Since $\varphi_{0 n}:=\varphi_{n}^{\prime}$, we have

$$
\begin{equation*}
\psi_{n}^{\prime}(\eta)=\omega_{2 n}(\eta)-\lambda_{2}^{\prime}(\eta) \varphi_{0 n}\left(\lambda_{2}(\eta)\right), \quad 0 \leq \eta \leq \eta_{0} \tag{3.47}
\end{equation*}
$$

by virtue of relation (3.30); by relations (3.24), (3.27), and (3.32), we have

$$
\begin{equation*}
\varphi_{n}(\xi)=\int_{0}^{\xi} \varphi_{0 n}\left(\xi^{\prime}\right) d \xi^{\prime} \in C^{2}\left(\left[0, \xi_{0}\right]\right), \quad \psi_{n}(\eta)=\int_{0}^{\eta} \psi_{n}^{\prime}\left(\eta^{\prime}\right) d \eta^{\prime} \in C^{2}\left(\left[0, \eta_{0}\right]\right) \tag{3.48}
\end{equation*}
$$

Remark 3.7. By keeping the same notation for the restrictions of the functions $\widetilde{u}_{n}$ and $\widetilde{F}_{n}$ to the subdomain $G_{T}$ of the domain $G_{0, T}$ and by taking into account their definition, we find that the function $\widetilde{u}_{n} \in C^{2}\left(\bar{G}_{T}\right)$ is a classical solution of the linear problem (3.10)-(3.12) for $\widetilde{F}=\widetilde{F}_{n}$; by Remark 3.2 and the estimate (3.13), the following inequality holds:

$$
\left\|\widetilde{u}_{n}-\widetilde{u}_{k}\right\|_{C\left(\bar{G}_{T}\right)} \leq 4 T^{3 / 2} \exp \left(2^{-1} T\right)\left\|\widetilde{F}_{n}-\widetilde{F}_{k}\right\|_{C\left(\bar{G}_{T}\right)}
$$

This, together with relations (3.23), implies that the function sequence $\widetilde{u}_{n} \in C^{2}\left(\bar{G}_{T}\right)$ is a Cauchy sequence in the complete space $C\left(\bar{G}_{T}\right)$; therefore, the exists a function $\widetilde{w} \in C\left(\bar{G}_{T}\right)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|\widetilde{u}_{n}-\widetilde{w}\right\|_{C\left(\bar{G}_{T}\right)}=0 \tag{3.49}
\end{equation*}
$$

By virtue of relations (3.23) and (3.49), the function $\widetilde{w}$ thus defined is a strong generalized solution of the linear problem (3.10)-(30.12) in the class $C$ in the domain $G_{T}$, whose uniqueness follows from the estimate (3.13). We denote this solution $\widetilde{w}$ by $\widetilde{\square}^{-1} \widetilde{F}$; i.e.,

$$
\begin{equation*}
\widetilde{w}=\widetilde{\square}^{-1} \widetilde{F}, \tag{3.50}
\end{equation*}
$$

where the linear operator $\widetilde{\square}^{-1}: C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)$ is continuous, and by (3.13), its norm satisfies the estimate

$$
\begin{equation*}
\left\|\widetilde{\square}^{-1}\right\|_{C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)} \leq 4 T^{3 / 2} \exp \left(2^{-1} T\right) \tag{3.51}
\end{equation*}
$$

Moreover, it follows from relations (3.31), (3.32), (3.34), (3.39)-(3.41), and (3.43)-(3.45) that the operator $\widetilde{\square}^{-1}$ occurring in relation (3.55) indeed maps any continuous function $\widetilde{F} \in C\left(\bar{G}_{T}\right)$ to a function $\widetilde{w} \in C^{1}\left(\bar{G}_{T}\right)$ and the linear operator

$$
\tilde{\square}^{-1}: C\left(\bar{G}_{T}\right) \rightarrow C^{1}\left(\bar{G}_{T}\right)
$$

is also continuous. [For details on the smoothness of $\widetilde{w}$ in (3.50), see Section 4, the representation (4.10).] The above-performed argument implies that, for the validity of the representation (3.50), i.e., for the unique solvability of the linear problem (3.10)-(3.12) in the class $C$, it suffices to require that $f \in C\left(\bar{D}_{T} \times \mathbb{R}\right), F \in C\left(\bar{D}_{T}\right)$, conditions (1.5), (1.6), and (2.4) are satisfied at the point $O$, and relations (3.14) and (3.26) and assumptions of Lemma 3.1 are valid.

Remark 3.8. Since the space $C^{1}\left(\bar{G}_{T}\right)$ is compactly embedded in $C\left(\bar{G}_{T}\right)$ [17, p. 135 of the Russian translation], it follows from Remark 3.7 that the linear operator $\widetilde{\square}^{-1}: C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)$ is compact, and its norm can be estimated as (3.51).

Remark 3.9. By virtue of Remarks 3.1 and 3.7 and relation (3.50), the function $u=u(x, t)$ is a strong generalized solution of problem (1.1)-(1.3) of the class $C$ in the domain $D_{T}$ if and only if $\widetilde{u}(\xi, \eta):=u\left(\frac{\xi-\eta}{2}, \frac{\xi+\eta}{2}\right)$ is a continuous solution of the functional equation

$$
\begin{equation*}
\widetilde{u}=K_{0} \widetilde{u}:=\widetilde{\square}^{-1}(-\widetilde{f}(\xi, \eta, \widetilde{u})+\widetilde{F}) \tag{3.52}
\end{equation*}
$$

in the class $C\left(\bar{G}_{T}\right)$, where $K_{0}: C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)$ is a continuous compact operator, because the nonlinear operator $N: C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)$ acting by the rule $N \widetilde{u}=-\widetilde{f}(\xi, \eta, \widetilde{u})+\widetilde{F}$, where $\widetilde{f} \in C\left(\bar{G}_{T} \times \mathbb{R}\right)$ and $\widetilde{F} \in C\left(\bar{G}_{T}\right)$, is bounded and continuous, and the linear operator $\widetilde{\square}^{-1}: C\left(\bar{G}_{T}\right) \rightarrow$ $C\left(\bar{G}_{T}\right)$ is compact by virtue of Remark 3.8. At the same time, by virtue of the estimate (3.9) and relations (2.37), the same a priori estimate (3.9) with the same constants $c_{1}$ and $c_{2}$ is valid for any parameter $\tau \in[0,1]$ and for any solution $\widetilde{u} \in C\left(\bar{G}_{T}\right)$ of the equation $\widetilde{u}=\tau K_{0} \widetilde{u}$. Therefore, by the Leray-Schauder theorem [18, p. 375], Eq. (3.52) has at least one solution $\widetilde{u} \in C\left(\bar{G}_{T}\right)$. Therefore, in view of Remarks 3.1 and 3.9 , we have thereby proved the following assertion.

Theorem 3.1. Let the conditions $f \in C\left(\bar{D}_{T} \times \mathbb{R}\right)$ and $F \in C\left(\bar{D}_{T}\right)$ as well as conditions (1.5), (1.6), (2.2)-(2.4), (3.14), and (3.26) be satisfied; moreover, in the case of $\left(l_{1} l_{2}\right)(O)=0$, assume that the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ either are not tangent at the point $O$ or are tangent but $\gamma_{2}^{\prime}(0)<0$. Then problem (1.1)-(1.3) has at least one strong generalized solution $u$ of the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Remark 3.10. One can readily see that if the assumptions of Theorem 3.1 are true for $T=\infty$, then problem (1.1)-(1.3) is globally solvable in the class $C$ in the sense of Definition 1.2.

## 4. SMOOTHNESS OF SOLUTION OF PROBLEM (1.1)-(1.3)

Now let us study the smoothness of the strong generalized solution of the nonlinear problem (1.1)-(1.3) depending on the smoothness of the data of that problem. To this end, under the
assumptions of Theorem 3.1 with regard of Remark 3.1 , we trace the scheme of the construction of a strong generalized solution $\widetilde{w}$ of the linear problem (3.10)-(3.12) in the class $C$ in the domain $G_{T}$ and show that such a solution actually belongs to the class $C^{1}\left(\bar{G}_{T}\right)$, and the boundary conditions (3.11) and (3.12) are satisfied pointwise. Indeed, by virtue of relations (3.31), (3.32), and (3.34), the right-hand side $\omega_{n}$ of Eq. (33) can be represented in the form

$$
\begin{align*}
\omega_{n}(\xi)= & -\frac{1}{m_{1}(\xi)}\left[m_{1}(\xi) \int_{0}^{\lambda_{1}(\xi)} \widetilde{F}_{n}\left(\xi, \eta^{\prime}\right) d \eta^{\prime}+m_{2}(\xi) \int_{0}^{\xi} \widetilde{F}_{n}\left(\xi^{\prime}, \lambda_{1}(\xi)\right) d \xi^{\prime}\right. \\
& \left.-m_{2}(\xi) \lambda_{2}^{\prime}\left(\lambda_{1}(\xi)\right) \int_{0}^{\lambda_{1}(\xi)} \widetilde{F}_{n}\left(\tau(\xi), \eta^{\prime}\right) d \eta^{\prime}-m_{2}(\xi) \int_{0}^{\tau(\xi)} \widetilde{F}_{n}\left(\xi^{\prime}, \lambda_{1}(\xi)\right) d \xi^{\prime}\right], \quad 0 \leq \xi \leq \xi_{0} \tag{4.1}
\end{align*}
$$

This, together with conditions (3.23), implies that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|\omega_{n}-\omega\right\|_{C\left(\bar{G}_{T}\right)}=0 \tag{4.2}
\end{equation*}
$$

where

$$
\begin{align*}
\omega(\xi):= & -\frac{1}{m_{1}(\xi)}\left[m_{1}(\xi) \int_{0}^{\lambda_{1}(\xi)} \widetilde{F}\left(\xi, \eta^{\prime}\right) d \eta^{\prime}+m_{2}(\xi) \int_{0}^{\xi} \widetilde{F}\left(\xi^{\prime}, \lambda_{1}(\xi)\right) d \xi^{\prime}\right. \\
& \left.-m_{2}(\xi) \lambda_{2}^{\prime}\left(\lambda_{1}(\xi)\right) \int_{0}^{\lambda_{1}(\xi)} \widetilde{F}\left(\tau(\xi), \eta^{\prime}\right) d \eta^{\prime}-m_{2}(\xi) \int_{0}^{\tau(\xi)} \widetilde{F}\left(\xi^{\prime}, \lambda_{1}(\xi)\right) d \xi^{\prime}\right], \quad 0 \leq \xi \leq \xi_{0} \tag{4.3}
\end{align*}
$$

In turn, it follows from relations (3.39)-(3.43), (4.1)-(4.3) that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|\varphi_{0 n}-\varphi_{0}\right\|_{C\left(\left[0, \xi_{0}\right]\right)}=0 \tag{4.4}
\end{equation*}
$$

where $\varphi_{0 n}:=\varphi_{n}^{\prime}$ and

$$
\begin{equation*}
\varphi_{0}:=\left[\sum_{k=0}^{\infty} \Lambda^{k} \omega\right] \in C\left(\left[0, \xi_{0}\right]\right) \tag{4.5}
\end{equation*}
$$

Since the derivative $\psi_{n}^{\prime}$ of the function $\psi_{n}$ occurring in the representation $(3.25)$ is defined by relation (3.47), it follows from (3.23), (3.32), and (4.4) that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|\psi_{n}^{\prime}-\psi_{0}\right\|_{C\left(\left[0, \eta_{0}\right]\right)}=0 \tag{4.6}
\end{equation*}
$$

where

$$
\begin{align*}
\psi_{0} \in C\left(\left[0, \eta_{0}\right]\right), \quad \psi_{0}(\eta):=\omega_{2}(\eta)-\lambda_{2}^{\prime}(\eta) \varphi_{0}\left(\lambda_{2}(\eta)\right), & 0 \leq \eta \leq \eta_{0}  \tag{4.7}\\
\omega_{2}(\eta):=-\lambda_{2}^{\prime}(\eta) \int_{0}^{\eta} \widetilde{F}\left(\lambda_{2}(\eta), \eta^{\prime}\right) d \eta^{\prime}-\int_{0}^{\lambda_{2}(\eta)} \widetilde{F}\left(\xi^{\prime}, \eta\right) d \xi^{\prime}, & 0 \leq \eta \leq \eta_{0} \tag{4.8}
\end{align*}
$$

Finally, by using Remark 3.7 and the limit relations (3.23), (3.49), (4.4), (4.6), and (3.48) in the notation

$$
\begin{equation*}
\varphi(\xi):=\int_{0}^{\xi} \varphi_{0}\left(\xi^{\prime}\right) d \xi^{\prime}, \quad 0 \leq \xi \leq \xi_{0}, \quad \psi(\eta):=\int_{0}^{\eta} \psi_{0}\left(\eta^{\prime}\right) d \eta^{\prime}, \quad 0 \leq \eta \leq \eta_{0} \tag{4.9}
\end{equation*}
$$

and by passing to the limit in relation (3.25), for the strong generalized solution $\widetilde{w}$ of the linear problem (3.10)-(3.12) in the class $C$ in the domain $G_{T}$ we obtain the representation

$$
\begin{equation*}
\widetilde{w}(\xi, \eta)=\varphi(\xi)+\psi(\eta)+\int_{0}^{\xi} d \xi^{\prime} \int_{0}^{\eta} \widetilde{F}\left(\xi^{\prime}, \eta^{\prime}\right) d \eta^{\prime}, \quad(\xi, \eta) \in \bar{G}_{T} \tag{4.10}
\end{equation*}
$$

If $\widetilde{F}$ belongs to $C\left(\bar{G}_{T}\right)$, then, by virtue of relations (4.5) and (4.7), it follows from the representation (4.10) that

$$
w \in C^{1}\left(\bar{G}_{T}\right)
$$

Next, by virtue of relations (4.2), (4.4) and (3.33), (3.35), the function $\varphi_{0}$ satisfies the functional equation

$$
\begin{equation*}
\varphi_{0}(\xi)-a(\xi) \varphi_{0}(\tau(\xi))=\omega(\xi), \quad 0 \leq \xi \leq \xi_{0} \tag{4.11}
\end{equation*}
$$

Remark 4.1. If the function $\widetilde{F}$ belongs to $C^{1}\left(\bar{G}_{T}\right)$ and the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are not tangent at the point $O$, then, by [19, p. 595], one can extend that function in the rectangle $\bar{G}_{0, T}$ (keeping the same notation for it) so as to ensure that the function $\widetilde{F}$ belongs to $C^{1}\left(\bar{G}_{0, T}\right)$. In the case of tangency of the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ at the point $O$, throughout the following we assume that such an extension is possible.

It follows from relation (4.2) that if condition (3.26) is satisfied and one additionally requires that the function $\widetilde{F}$ belongs to $C^{1}\left(\bar{G}_{T}\right)$, then the right-hand side $\omega$ of Eq. (4.11) belongs to the class $C^{1}\left(\left[0, \xi_{0}\right]\right)$. This, together with the argument carried out in Remark 3.6, implies that the solution of Eq. (4.11) belongs to the space $C^{1}\left(\left[0, \xi_{0}\right]\right)$; consequently, by (4.7) and (4.8), the function $\psi_{0}$ belongs to the space $C^{1}\left(\left[0, \eta_{0}\right]\right)$ as well. Therefore, under the above-stipulated assumptions with regard of notation (4.9), we find that the function $\widetilde{w}$ occurring in (4.10) belongs to the space $C^{2}\left(\bar{G}_{T}\right)$. Thus, in view of Remark 3.7, we have proved the following assertion.

Theorem 4.1. If conditions (1.5), (1.6), (2.4), (3.14), and (3.26) are satisfied, $\widetilde{F} \in C\left(\bar{G}_{T}\right)$, and moreover, for $\left(l_{1} l_{2}\right)(O)=0$ the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ either are not tangent at the point $O$ or are tangent but $\gamma_{2}^{\prime}(0)<0$, then the strong generalized solution $\widetilde{w}$ of the linear problem (3.10)-(3.12) in the class $C$ in the domain $G_{T}$ belongs to the space $C^{1}\left(\bar{G}_{T}\right)$; i.e., by relation (3.50), $\widetilde{w}=\widetilde{\square}^{-1} \widetilde{F}$ in the class $C^{1}\left(\bar{G}_{T}\right)$; and if it is additionally required that the function $\widetilde{F}$ belongs to $C^{1}\left(\bar{G}_{T}\right)$, then $\widetilde{w}$ belongs to $C^{2}\left(\bar{G}_{T}\right)$; in addition, the boundary conditions (3.11) and (3.12) are valid pointwise in both cases.

The following assertion is a consequence of Remarks 3.1 and 3.9 , relation (3.52), and Theorem 4.1.

Theorem 4.2. If the assumptions of Theorem 3.1 are satisfied, then a strong generalized solution $u$ of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$ belongs to the space $C^{1}\left(\bar{D}_{T}\right)$; under the additional requirements $f \in C^{1}\left(\bar{D}_{T} \times \mathbb{R}\right)$ and $F \in C^{1}\left(\bar{D}_{T}\right)$, this solution belongs to the space $C^{2}\left(\bar{D}_{T}\right)$, i.e., is classical; moreover, in both cases the boundary conditions (1.1) and (1.3) are satisfied pointwise.

## 5. UNIQUENESS THEOREM. EXISTENCE OF GLOBAL SOLUTION OF PROBLEM (1.1)-(1.3) IN THE DOMAIN $D_{\infty}$

By definition, a function $f=f(x, t, s)$ satisfies the local Lipschitz condition with respect to the variable $s$ on the set $\bar{D}_{T} \times \mathbb{R}$ if

$$
\begin{equation*}
\left|f\left(x, t, s_{2}\right)-f\left(x, t, s_{1}\right)\right| \leq M(T, r)\left|s_{2}-s_{1}\right|, \quad(x, t) \in \bar{D}_{T}, \quad\left|s_{i}\right| \leq r, \quad i=1,2 \tag{5.1}
\end{equation*}
$$

where $M(T, r):=$ const $\geq 0$.

Theorem 5.1. Let condition (2.4) be satisfied, let the function $f \in C\left(\bar{D}_{T} \times \mathbb{R}\right)$ satisfy condition (5.1), let $F$ belong to $C\left(\bar{D}_{T}\right)$, and let $l_{1}$ and $l_{2}$ belong to the class $C\left(\gamma_{1, T}\right)$. Then problem (1.1)-(1.3) has at most one strong generalized solution in the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Proof. Indeed, assume that problem (1.1)-(1.3) has two possible distinct strong generalized solutions $u^{1}$ and $u^{2}$ in the class $C$ in the domain $D_{T}$. Then, by Definition 1.1, there exist sequences of functions $u_{n}^{i} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right), i=1,2$, such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|u_{n}^{i}-u^{i}\right\|_{C\left(\bar{D}_{T}\right)}=0, \quad \lim _{n \rightarrow \infty}\left\|L u_{n}^{i}-F\right\|_{C\left(\bar{D}_{T}\right)}=0, \quad i=1,2 \tag{5.2}
\end{equation*}
$$

Set $\omega_{n}:=u_{n}^{2}-u_{n}^{1}$. One can readily see that the function $\omega_{n} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ is a classical solution of the problem

$$
\begin{equation*}
\square \omega_{n}+g_{n}=F_{n},\left.\quad\left(l_{1} \omega_{n x}+l_{2} \omega_{n t}\right)\right|_{\gamma_{1, T}}=0,\left.\quad \omega_{n}\right|_{\gamma_{2, T}}=0 \tag{5.3}
\end{equation*}
$$

Here

$$
\begin{equation*}
g_{n}:=f\left(x, t, u_{n}^{2}\right)-f\left(x, t, u_{n}^{1}\right), \quad F_{n}:=L u_{n}^{2}-L u_{n}^{1} \tag{5.4}
\end{equation*}
$$

By virtue of relations (5.2), there exists a number $m:=$ const $>0$ independent of the indices $i$ and $n$ such that $\left\|u_{n}^{i}\right\|_{C\left(\bar{D}_{T}\right)} \leq m$, which, together with relations (5.1) and (5.2), implies that

$$
\begin{equation*}
\left|g_{n}\right| \leq M(T, m)\left|\omega_{n}\right| \tag{5.5}
\end{equation*}
$$

By virtue of relations (5.2) and the second relation in (5.4), we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty}\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}=0 \tag{5.6}
\end{equation*}
$$

By multiplying both sides of the first relation in (5.3) by $\omega_{n t}$, by integrating the resulting relation over the domain

$$
D_{\tau}:=\left\{(x, t) \in D_{T}: t<\tau\right\}, \quad 0<\tau \leq T,
$$

and by following the derivation of relation (2.13) in (2.8)-(2.10), we obtain

$$
\begin{align*}
w_{n}(\tau):= & \int_{\Omega_{\tau}}\left(\omega_{n t}^{2}+\omega_{n x}^{2}\right) d x=-\int_{\gamma_{1, \tau}}\left(\omega_{n t}^{2} \nu_{t}-2 \omega_{n x} \omega_{n t} \nu_{x}+\omega_{n x}^{2} \nu_{t}\right) d s \\
& -\int_{\gamma_{2, \tau}} \frac{1}{\nu_{t}}\left[\left(\omega_{n x} \nu_{t}-\omega_{n t} \nu_{x}\right)^{2}+\omega_{n t}^{2}\left(\nu_{t}^{2}-\nu_{x}^{2}\right)\right] d s+2 \int_{D_{\tau}}\left(F_{n}-g_{n}\right) \omega_{n t} d x d t . \tag{5.7}
\end{align*}
$$

By virtue of inequality (5.5) and the Cauchy inequality, we have the estimate

$$
\begin{align*}
\left|2 \int_{D_{\tau}}\left(F_{n}-g_{n}\right) \omega_{n t} d x d t\right| & \leq \int_{D_{\tau}}\left(F_{n}-g_{n}\right)^{2} d x d t+\int_{D_{\tau}} \omega_{n t}^{2} d x d t \\
& \leq 2 \int_{D_{\tau}} F_{n}^{2} d x d t+2 M^{2}(T, m) \int_{D_{\tau}} \omega_{n}^{2} d x d t+\int_{D_{\tau}} \omega_{n t}^{2} d x d t . \tag{5.8}
\end{align*}
$$

Since inequalities (2.15) and (2.19), true for $u_{n}$, also hold for $\omega_{n}$, it follows from relations (5.7) and (5.8) that

$$
\begin{equation*}
w_{n}(\tau) \leq 2 M^{2}(T, m) \int_{D_{\tau}} \omega_{n}^{2} d x d t+\int_{D_{\tau}} \omega_{n t}^{2} d x d t+2 \int_{D_{\tau}} F_{n}^{2} d x d t \tag{5.9}
\end{equation*}
$$

Since inequality (2.29), true for $u_{n}$, also holds for $\omega_{n}$, from the estimate (5.9), we have

$$
\begin{align*}
w_{n}(\tau) & \leq\left(2 M^{2}(T, m) T^{2}+1\right) \int_{D_{\tau}} \omega_{n t}^{2} d x d t+2 \int_{D_{T}} F_{n}^{2} d x d t \\
& \leq M_{0} \int_{D_{\tau}}\left(\omega_{n t}^{2}+\omega_{n x}^{2}\right) d x d t+2 \int_{D_{T}} F_{n}^{2} d x d t \tag{5.10}
\end{align*}
$$

where $M_{0}:=2 M^{2}(T, m) T^{2}+1$.
By taking into account the relation

$$
\int_{D_{\tau}}\left(\omega_{n t}^{2}+\omega_{n x}^{2}\right) d x d t=\int_{0}^{\tau} w_{n}(\sigma) d \sigma,
$$

from inequality (5.10) we obtain

$$
w_{n}(\tau) \leq M_{0} \int_{0}^{\tau} w_{n}(\sigma) d \sigma+2\left\|F_{n}\right\|_{C\left(\bar{D}_{T)}\right.}^{2} \operatorname{mes} D_{T}, \quad 0<\tau \leq T .
$$

This, together with the Gronwall lemma, implies that

$$
\begin{equation*}
w_{n}(\tau) \leq 2\left\|F_{n}\right\|_{C\left(\bar{D}_{T}\right)}^{2}\left(\operatorname{mes} D_{T}\right) \exp \left(M_{0} T\right), \quad 0<\tau \leq T \tag{5.11}
\end{equation*}
$$

Since inequality (2.34), true for $u_{n}$, also holds for $\omega_{n}$, it follows from the estimate (5.11) and inequality (2.35) that

$$
\begin{equation*}
\left|\omega_{n}(x, t)\right|^{2} \leq T w_{n}(t) \leq 2 T\left\|F_{n}\right\|_{C\left(\bar{D}_{T)}\right.}^{2}\left(\operatorname{mes} D_{T}\right) \exp \left(M_{0} T\right), \quad(x, t) \in \bar{D}_{T} \backslash O . \tag{5.12}
\end{equation*}
$$

By using relations (5.2) and (5.6) and the relation $\omega_{n}:=u_{n}^{2}-u_{n}^{1}$ and by passing in inequality (5.12) to the limit as $n \rightarrow \infty$, we obtain $\left|\left(u^{2}-u^{1}\right)(x, t)\right|^{2} \leq 0,(x, t) \in \bar{D}_{T} \backslash O$; i.e., $u^{2}=u^{1}$, which contradicts the above assumption. The proof of Theorem 5.1 is complete.

Remark 5.1. Obviously, condition (5.1) is satisfied if $f \in C^{1}\left(\bar{D}_{T} \times \mathbb{R}\right)$.
Theorems 3.1, 4.2, and 5.1 and Remark 5.1 imply the following assertion.
Theorem 5.2. Let $f \in C^{1}\left(\bar{D}_{T} \times \mathbb{R}\right)$ and $F \in C^{1}\left(\bar{D}_{T}\right)$, and let conditions (1.5), (1.6), (2.2)-(2.4), (3.14), and (3.26) be satisfied. Moreover, assume that in the case of $\left(l_{1} l_{2}\right)(O)=0$ either the curves $\gamma_{1, T}$ and $\gamma_{2, T}$ are not tangent at the point $O$ or $\gamma_{2}^{\prime}(0)<0$. Then problem (1.1)-(1.3) has a unique classical solution $u \in C^{2}\left(\bar{D}_{T}\right)$ in the domain $D_{T}$.

Corollary 5.1. If the assumptions of Theorem 5.2 hold for $T=\infty$, then problem (1.1)-(1.3) has a unique global classical solution $u \in C^{2}\left(\bar{D}_{\infty}\right)$.

Indeed, by Theorem 5.2, problem (1.1)-(1.3) for $T=n$ has a unique classical solution $u_{n}$ in the domain $D_{n}$. Since $u_{n+1}$ is a classical solution of that problem in the domain $D_{n}$ as well, we have $\left.u_{n+1}\right|_{D_{n}}=u_{n}$ by virtue of the uniqueness Theorem 5.1. Therefore, the function $u$ constructed in the domain $D_{\infty}$ by the rule $u(x, t)=u_{n}(x, t)$ for $n=[t]+1$, where $[t]$ is the integer part of the number $t$ and $(x, t) \in \bar{D}_{\infty}$, is the unique global classical solution of problem (1.1)-(1.3) in the domain $D_{\infty}$.

## 6. CASES OF THE ABSENCE OF GLOBAL SOLVABILITY OF PROBLEM (1.1)-(1.3) AND ITS LOCAL SOLVABILITY

In what follows, we show that if condition (2.2) fails, then problem (1.1)-(1.3) is not necessarily globally solvable in the class $C$ in the sense of Definition 1.2. To this end, we use the method of test functions described in [20, pp. 10-14].

Lemma 6.1. Let $u$ be a strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1. Then the integral relation

$$
\begin{equation*}
\int_{D_{T}} u \square \varphi d x d t+\int_{D_{T}} f(x, t, u) \varphi d x d t=\int_{D_{T}} F \varphi d x d t \tag{6.1}
\end{equation*}
$$

holds for any test function $\varphi$ such that

$$
\begin{equation*}
\varphi \in C^{2}\left(\bar{D}_{T}\right),\left.\quad \varphi\right|_{\partial D_{T}}=0,\left.\quad \nabla \varphi\right|_{\partial D_{T}}=0 \tag{6.2}
\end{equation*}
$$

where $\nabla:=\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial t}\right)$.
Proof. By the definition of a strong generalized solution $u$ of problem (1.1)-(1.3) in the class $C$, in the domain $D_{T}$, we have $u \in C\left(\bar{D}_{T}\right)$, and there exists a function sequence $u_{n} \in \dot{C}^{2}\left(\bar{D}_{T}, \gamma_{T}\right)$ such that the limit relations (2.7) hold.

Set $F_{n}:=L u_{n}$. We multiply both sides of the relation $L u_{n}=F_{n}$ by the function $\varphi$ and integrate the resulting relation over the domain $D_{T}$. By virtue of condition (6.2), after the integration by parts in the resulting integral relation, we obtain

$$
\begin{equation*}
\int_{D_{T}} u_{n} \square \varphi d x d t+\int_{D_{T}} f\left(x, t, u_{n}\right) \varphi d x d t=\int_{D_{T}} F_{n} \varphi d x d t \tag{6.3}
\end{equation*}
$$

By taking into account the limit relations (2.7) and by passing in relation (6.3) to the limit as $n \rightarrow \infty$, we obtain the desired relation (6.1). The proof of Lemma 6.1 is complete.

Consider the following condition imposed on the function $f$ :

$$
\begin{equation*}
f(x, t, s) \leq-\lambda|s|^{\alpha+1}, \quad(x, t, s) \in \bar{D}_{\infty} \times \mathbb{R} ; \quad \lambda, \alpha:=\text { const }>0 \tag{6.4}
\end{equation*}
$$

One can readily see that condition (2.2) fails in case (6.4).
We introduce the function $\varphi^{0}:=\varphi^{0}(x, t)$ satisfying the conditions

$$
\begin{equation*}
\varphi^{0} \in C^{2}\left(\bar{D}_{\infty}\right),\left.\quad \varphi^{0}\right|_{D_{T=1}}>0,\left.\quad \varphi^{0}\right|_{\partial D_{T=1}}=0,\left.\quad \nabla \varphi^{0}\right|_{\partial D_{T=1}}=0,\left.\quad \varphi^{0}\right|_{t \geq 1}=0 \tag{6.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\kappa_{0}:=\int_{D_{T=1}} \frac{\left|\square \varphi^{0}\right|^{p^{\prime}}}{\left|\varphi^{0}\right| p^{p^{\prime}-1}} d x d t<\infty, \quad p^{\prime}=1+\frac{1}{\alpha} \tag{6.6}
\end{equation*}
$$

To simplify the exposition, we consider the case in which the curves $\gamma_{1}$ and $\gamma_{2}$ are rays; i.e.,

$$
\begin{equation*}
\gamma_{i}: x=-k_{i} t, \quad k_{i}:=\text { const }, \quad i=1,2 ; \quad 0<k_{1}<k_{2}<1 . \tag{6.7}
\end{equation*}
$$

One can readily see that, in the case of (6.7), for the function $\varphi^{0}$ satisfying conditions (6.5) and (6.6) one can take the function

$$
\varphi^{0}(x, t)= \begin{cases}{\left[\left(x+k_{1} t\right)\left(x+k_{2} t\right)(1-t)\right]^{m}} & \text { for } \quad(x, t) \in D_{T=1} \\ 0 & \text { for } \quad t \geq 1\end{cases}
$$

for a sufficiently large $m:=$ const $>0$.
By setting $\varphi_{T}(x, t):=\varphi^{0}\left(\frac{x}{T}, \frac{t}{T}\right), T>0$, and by using conditions (6.5), one can readily see that

$$
\begin{equation*}
\varphi_{T} \in C^{2}\left(\bar{D}_{\infty}\right),\left.\quad \varphi_{T}\right|_{D_{T}}>0,\left.\quad \varphi_{T}\right|_{\partial D_{T}}=0,\left.\quad \nabla \varphi_{T}\right|_{\partial D_{T}}=0,\left.\quad \varphi_{T}\right|_{t \geq T}=0 \tag{6.8}
\end{equation*}
$$

By assuming that $F \in C\left(\bar{D}_{\infty}\right)$ is a fixed function, we introduce the following function of one variable $T$ :

$$
\begin{equation*}
\zeta(T):=\int_{D_{T}} F \varphi_{T} d x d t, \quad T>0 \tag{6.9}
\end{equation*}
$$

We have the following assertion on the absence of the global solvability of problem (1.1)-(1.3).

Theorem 6.1. Let the function $f \in C\left(\bar{D}_{\infty} \times \mathbb{R}\right)$ satisfy condition (6.4), let $F$ belong to $C\left(\bar{D}_{\infty}\right)$, $F \geq 0$, in the domain $D_{\infty}$, and let

$$
\begin{equation*}
\liminf _{T \rightarrow \infty} \zeta(T)>0 \tag{6.10}
\end{equation*}
$$

Then there exists a positive number $T_{0}=T_{0}(F)$ such that problem (1.1)-(1.3) with $T>T_{0}$ cannot have a strong generalized solution in the class $C$ in the domain $D_{T}$ in the sense of Definition 1.1.

Proof. Suppose that, under assumptions of the theorem, there exists a strong generalized solution $u$ of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$. Then, by Lemma 6.1, relation (6.1) holds, where, by virtue of condition (6.8), for the test function $\varphi$ one can take the function $\varphi_{T}$; i.e.,

$$
\begin{equation*}
-\int_{D_{T}} f(x, t, u) \varphi_{T} d x d t+\int_{D_{T}} F \varphi_{T} d x d t=\int_{D_{T}} u \square \varphi_{T} d x d t \tag{6.11}
\end{equation*}
$$

Since the function $\varphi_{T}$ is positive in the domain $D_{T}$, it follows from condition (6.4), notation (6.9), and relation (6.11) that

$$
\begin{equation*}
\lambda \int_{D_{T}}|u|^{p} \varphi_{T} d x d t \leq \int_{D_{T}}|u|\left|\square \varphi_{T}\right| d x d t-\zeta(T), \quad p:=\alpha+1 . \tag{6.12}
\end{equation*}
$$

If in the Young inequality with parameter $\varepsilon>0$,

$$
a b \leq \frac{\varepsilon}{p} a^{p}+\frac{1}{p^{\prime} \varepsilon^{p^{\prime}-1}} b^{p^{\prime}} ; \quad a, b \geq 0, \quad \frac{1}{p}+\frac{1}{p^{\prime}}=1, \quad p>1,
$$

we take $a=|u| \varphi_{T}^{1 / p}$ and $b=\left|\square \varphi_{T}\right| / \varphi_{T}^{1 / p}$, then, by virtue of the relation $p^{\prime} / p=p^{\prime}-1$, we obtain

$$
\begin{equation*}
\left|u \square \varphi_{T}\right|=|u| \varphi_{T}^{1 / p} \frac{\left|\square \varphi_{T}\right|}{\varphi_{T}^{1 / p}} \leq \frac{\varepsilon}{p}|u|^{p} \varphi_{T}+\frac{1}{p^{\prime} \varepsilon^{p^{\prime}-1}} \frac{\left|\square \varphi_{T}\right|^{p^{\prime}}}{\varphi_{T}^{p^{\prime}-1}} . \tag{6.13}
\end{equation*}
$$

It follows from inequalities (6.12) and (6.13) that

$$
\left(\lambda-\frac{\varepsilon}{p}\right) \int_{D_{T}}|u|^{p} \varphi_{T} d x d t \leq \frac{1}{p^{\prime} \varepsilon^{p^{\prime}-1}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{p^{\prime}}}{\varphi_{T}^{p^{\prime}-1}} d x d t-\zeta(T),
$$

whence for $\varepsilon<\lambda p$ we obtain

$$
\begin{equation*}
\int_{D_{T}}|u|^{p} \varphi_{T} d x d t \leq \frac{p}{(\lambda p-\varepsilon) p^{\prime} \varepsilon^{p^{\prime}-1}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{p^{\prime}}}{\varphi_{T}^{p^{\prime}-1}} d x d t-\frac{p}{\lambda p-\varepsilon} \zeta(T) . \tag{6.14}
\end{equation*}
$$

By taking into account the relations $p^{\prime}=\frac{p}{p-1}, p=\frac{p^{\prime}}{p^{\prime}-1}$, and the minimum

$$
\min _{0<\varepsilon<\lambda p} \frac{p}{(\lambda p-\varepsilon) p^{\prime} \varepsilon^{p^{\prime}-1}}=\frac{1}{\lambda^{p^{\prime}}},
$$

which is attained for $\varepsilon=\lambda$, we rewrite inequality (6.14) in the form

$$
\begin{equation*}
\int_{D_{T}}|u|^{p} \varphi_{T} d x d t \leq \frac{1}{\lambda^{p^{\prime}}} \int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{p^{\prime}}}{\varphi_{T}^{p^{\prime}-1}} d x d t-\frac{p^{\prime}}{\lambda} \zeta(T) . \tag{6.15}
\end{equation*}
$$

Since $\varphi_{T}(x, t):=\varphi^{0}\left(\frac{x}{T}, \frac{t}{T}\right)$, it follows that, by taking into account relation (6.6) and by performing the change of variables $x=T x_{1}$ and $t=T t_{1}$, one can represent the integral on the right-hand side in inequality (6.15) in the form

$$
\begin{equation*}
\int_{D_{T}} \frac{\left|\square \varphi_{T}\right|^{p^{\prime}}}{\varphi_{T}^{p^{\prime}-1}} d x d t=T^{-2\left(p^{\prime}-1\right)} \int_{D_{T=1}} \frac{\left|\square \varphi^{0}\right| p^{p^{\prime}}}{\left.\left|\varphi^{0}\right|\right|^{p^{\prime}-1}} d x_{1} d t_{1}=T^{-2\left(p^{\prime}-1\right)} \kappa_{0}<\infty . \tag{6.16}
\end{equation*}
$$

By virtue of relations (6.8) and (6.16), it follows from inequality (6.15) that

$$
\begin{equation*}
0 \leq \int_{D_{T}}|u|^{p} \varphi_{T} d x d t \leq \frac{1}{\lambda^{p^{\prime}}} T^{-2\left(p^{\prime}-1\right)} \kappa_{0}-\frac{p^{\prime}}{\lambda} \zeta(T) \tag{6.17}
\end{equation*}
$$

Since $p^{\prime}>1$, we have $-2\left(p^{\prime}-1\right)<0$, and by virtue of (6.6),

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \frac{1}{\lambda^{p^{\prime}}} T^{-2\left(p^{\prime}-1\right)} \kappa_{0}=0 \tag{6.18}
\end{equation*}
$$

By virtue of relations (6.10) and (6.18), there exists a positive number $T_{0}=T_{0}(F)$ such that for $T>T_{0}$ the right-hand side of inequality (6.17) is negative, while the left-hand side of this inequality is nonnegative. Hence it follows that if there exists a strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$, then the inequality $T \leq T_{0}$ is necessarily true. The proof of Theorem 6.1 is complete.

Remark 6.1. One can readily see that if the conditions $F \in C\left(\bar{D}_{\infty}\right), F \geq 0$, and $F(x, t) \geq c t^{-m}$ are satisfied for $t \geq 1$, where $c:=$ const $>0$ and $0 \leq m:=$ const $\leq 2$, then inequality (6.10) holds, and by Theorem 6.1 problem (1.1)-(1.3) does not have a strong generalized solution in the class $C$ in the domain $D_{T}$ for sufficiently large $T$ in this case.

Corollary 5.2. Under the assumptions of Theorem 6.1, problem (1.1)-(1.3) is not globally solvable in the class $C$ in the sense of Definition 1.2; i.e., it cannot have a global strong generalized solution in the class $C$ in the domain $D_{\infty}$ in the sense of Definition 1.3.

In what follows, we show that, although the global solvability of problem (1.1)-(1.3) has been proved under condition (2.2), the local solvability of this problem remains valid if that condition fails.

Theorem 6.2. Let $f \in C\left(\bar{D}_{\infty} \times \mathbb{R}\right)$ and $F \in C\left(\bar{D}_{\infty}\right)$, and let conditions (1.5), (1.6), (3.14), and (3.26) be satisfied; moreover, suppose that in the case $\left(l_{1} l_{2}\right)(O)=0$ the curves $\gamma_{1}$ and $\gamma_{2}$ either are not tangent at the point $O$ or are tangent but $\gamma_{2}^{\prime}(0)<0$. Then problem (1.1)-(1.3) is locally solvable in the class $C$ in the sense of Definition 1.4; i.e., there exists a positive number $T_{0}=T_{0}(F)$ such that this problem with $T \leq T_{0}$ has at least one strong generalized solution $u$ in the class $C$ in the domain $D_{T}$.

Proof. By Remarks 3.7 and 3.9, a function $u \in C\left(\bar{D}_{T}\right)$ is a strong generalized solution of problem (1.1)-(1.3) in the class $C$ in the domain $D_{T}$ if and only if

$$
\widetilde{u}(\xi, \eta):=u\left(\frac{\xi-\eta}{2}, \frac{\xi+\eta}{2}\right)
$$

is a solution of the functional equation (3.52) in the class $C\left(\bar{G}_{T}\right)$, where $K_{0}: C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)$ is a continuous compact operator. Therefore, by virtue of the Schauder theorem [18, p. 370], for the solvability of Eq. (3.52) in the space $C\left(\bar{G}_{T}\right)$, it suffices to show that the operator $K_{0}$ maps some ball

$$
B(O, r):=\left\{w \in C\left(\bar{G}_{T}\right):\|w\|_{C\left(\bar{G}_{T}\right)} \leq r\right\}
$$

of radius $r>0$ (which is a closed convex set in the Banach space $C\left(\bar{G}_{T}\right)$ ) into itself for sufficiently small $T$.

Take an arbitrary positive number $T_{*}$ and assume that $T \leq T_{*}$. By virtue of relations (3.15) and (3.52) for

$$
\begin{equation*}
\|w\|_{C\left(\bar{G}_{T}\right)} \leq r, \quad f^{*}:=\sup _{\substack{(\xi, \eta) \in \bar{G}_{T_{*}} \\|s| \leq r}}|\widetilde{f}(\xi, \eta, s)|, \quad F^{*}:=\|\widetilde{F}\|_{C\left(\bar{G}_{T_{*}}\right)} \tag{6.19}
\end{equation*}
$$

with regard of the embedding $D_{T} \subset D_{T_{*}}$, we obtain

$$
\begin{align*}
\left\|K_{0} w\right\|_{C\left(\bar{G}_{T}\right)} & \leq\left\|\widetilde{\square}^{-1}\right\|_{C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)} \sup _{\substack{(\xi, \eta) \in \bar{G}_{T_{*}},|s| \leq r}}|\widetilde{f}(\xi, \eta, s)|+\left\|\widetilde{\square}^{-1}\right\|_{C\left(\bar{G}_{T}\right) \rightarrow C\left(\bar{G}_{T}\right)}\|\widetilde{F}\|_{C\left(\bar{G}_{T}\right)} \\
& \leq 4 T^{3 / 2} \exp \left(2^{-1} T\right)\left(f^{*}+F^{*}\right) \tag{6.20}
\end{align*}
$$

It follows from relations (6.19) and (6.20) that if

$$
T \leq T_{0}:=\min \left\{T_{*}, h^{-1}\left[4^{-1} r\left(f^{*}+F^{*}\right)^{-1}\right]\right\}
$$

where $h^{-1}$ is the function inverse to $h(s):=s^{3 / 2} \exp \left(2^{-1} s\right), s>0$, then $\left\|K_{0} w\right\|_{C\left(\bar{G}_{T}\right)} \leq r$ for $\|w\|_{C\left(\bar{G}_{T}\right)} \leq r$. The proof of Theorem 6.2 is complete.

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