

## Characterizing the Formation of Singularities in a Superlinear Indefinite Mean Curvature Problem

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In this contribution, based on the very recent paper [21], we analyze the quasilinear indefinite Neumann problem

$$\begin{cases} -\left(\frac{u'}{\sqrt{1+(u')^2}}\right)' = \lambda a(x)f(u) & \text{in } (0, 1), \\ u'(0) = u'(1) = 0. \end{cases} \tag{1}$$

Here,  $\lambda \in \mathbb{R}$  is regarded as a parameter and

(a<sub>1</sub>) the function  $a \in L^\infty(0, 1)$  satisfies, for some  $z \in (0, 1)$ ,  $a(x) > 0$  a.e. in  $(0, z)$  and  $a(x) < 0$  a.e. in  $(z, 1)$ , as well as  $\int_0^1 a(x) dx < 0$ ;

(f<sub>1</sub>) the function  $f \in \mathcal{C}(\mathbb{R}) \cap \mathcal{C}^1[0, +\infty)$  satisfies  $f(s) > 0$  and  $f'(s) \geq 0$  for all  $s > 0$ , and there exist four constants,  $h > 0$ ,  $k > 0$ ,  $q > 1$  and  $p \geq 2$ , such that

$$\lim_{s \rightarrow +\infty} \frac{f(s)}{s^{q-1}} = qh, \quad \lim_{s \rightarrow 0^+} \frac{f(s)}{s^{p-1}} = pk.$$

Condition (f<sub>1</sub>) implies that the potential  $F$  of  $f$ , defined by  $F(s) = \int_0^s f(t) dt$ , satisfies

$$\lim_{s \rightarrow +\infty} \frac{F(s)}{s^q} = h, \quad \lim_{s \rightarrow 0^+} \frac{F(s)}{s^p} = k$$

and, thus,  $F$  must be superlinear at  $+\infty$  and either quadratic or superquadratic at 0. We also introduce the following condition on the weight function  $a$  at the nodal point  $z$ , which is going to play a pivotal role in the mathematical analysis carried out in [21]

$$(a_2) \quad \left(\int_x^z a(t) dt\right)^{-\frac{1}{2}} \in L^1(0, z) \text{ and } \left(\int_x^z a(t) dt\right)^{-\frac{1}{2}} \in L^1(z, 1).$$

We use the following notions of a solution.

- A couple  $(\lambda, u)$  is said to be a *regular solution* of (1) if  $u \in W^{2,1}(0, 1)$  and it satisfies the differential equation a.e. in  $(0, 1)$ , as well as the boundary conditions.

- A couple  $(\lambda, u)$  is said to be a *bounded variation solution* of (1) if  $u \in BV(0, 1)$  and it satisfies

$$\int_0^1 \frac{D^a u D^a \phi}{\sqrt{1 + |D^a u|^2}} dx + \int_0^1 \frac{D^s u}{|D^s u|} D^s \phi = \int_0^1 \lambda a f(u) \phi dx$$

for all  $\phi \in BV(0, 1)$  such that  $|D^s \phi|$  is absolutely continuous with respect to  $|D^s u|$  (cf. [2]).

- A couple  $(\lambda, u)$  is said to be a *singular solution* of (1) whenever it is a non-regular bounded variation solution; that is,  $u \in BV(0, 1) \setminus W^{2,1}(0, 1)$ .
- When the couple  $(\lambda, u)$  solves (1) in any of the previous senses, it is said that  $(\lambda, u)$  is a *positive solution* if, in addition,

$$\lambda > 0, \quad \text{ess inf } u > 0.$$

As usual, for any function  $v \in BV(0, 1)$ ,

$$Dv = D^a v dx + D^s v$$

stands for the Lebesgue decomposition of the Radon measure  $Dv$  and  $\frac{D^s v}{|D^s v|}$  denotes the density function of the measure  $D^s v$  with respect to its total variation  $|D^s v|$  (see [1]). By [23, Prop. 3.6], any positive singular solution,  $(\lambda, u)$ , of (1) actually satisfies

$$\begin{aligned} u|_{[0,z]} &\in W_{\text{loc}}^{2,1}[0, z] \cap W^{1,1}(0, z) \text{ and is concave,} \\ u|_{(z,1]} &\in W_{\text{loc}}^{2,1}(z, 1] \cap W^{1,1}(z, 1) \text{ and is convex;} \end{aligned} \quad (2)$$

moreover,  $u'(x) < 0$  for every  $x \in (0, 1) \setminus \{z\}$ ,  $u'(0) = u'(1) = 0$  and

$$u'(z^-) = u'(z^+) = -\infty,$$

where  $u'(z^-)$  and  $u'(z^+)$  are the left and the right Dini derivatives of  $u$  at  $z$ . The same argument used in [23, Lem. 2.1] shows that  $\lambda > 0$  is necessary for the existence of positive non-constant, either regular or singular, solutions.

Problem (1) is a one-dimensional prototype model of

$$\begin{cases} -\text{div} \left( \frac{\nabla u}{\sqrt{1 + |\nabla u|^2}} \right) = g(x, u) & \text{in } \Omega, \\ -\frac{\nabla u \cdot \nu}{\sqrt{1 + |\nabla u|^2}} = \sigma & \text{on } \partial\Omega, \end{cases} \quad (3)$$

where  $\Omega$  is a bounded regular domain in  $\mathbb{R}^N$ , with outward pointing normal  $\nu$ , and  $g : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$  and  $\sigma : \partial\Omega \rightarrow \mathbb{R}$  are given functions. Problem (3) plays a central role in the mathematical analysis of a number of geometrical and physical issues, such as prescribed mean curvature problems for cartesian surfaces in the Euclidean space [3, 9, 12–15, 19, 25, 26], capillarity phenomena for incompressible fluids [6, 10, 11, 16, 17], and reaction-diffusion processes where the flux features saturation at high regimes [5, 18, 24].

The model (1) has been recently investigated by the authors in [22, 23] and [20]. In [22] the existence of bounded variation solutions was analyzed by using variational methods and in [23] the existence of regular solutions was dealt with by means of classical phase plane and bifurcation techniques. The main result of [20] established the existence of a component of bounded variation

solutions bifurcating from the trivial state  $(\lambda, 0)$  in the special, but significant, case where  $p = 2$ . According to the results of these papers, it is already known that, under conditions  $(a_1)$  and  $(f_1)$ , problem (1) cannot admit positive solutions if  $\lambda < 0$  and that it possesses at least one positive bounded variation solution for sufficiently small  $\lambda > 0$ .

Quite strikingly, whether or not these bounded variation solutions are singular depends on whether or not condition  $(a_2)$  holds true: this is the main result of [21] which can be stated as follows.

**Theorem 1.** *Assume  $(a_1)$  and  $(f_1)$ . Then, the following conclusions hold for sufficiently small  $\lambda > 0$ :*

- (i) *any positive solution of (1) is singular if  $(a_2)$  holds;*
- (ii) *any positive solution of (1) is regular if  $(a_2)$  fails.*

In other words, condition  $(a_2)$  completely characterizes, under  $(a_1)$  and  $(f_1)$ , the development of singularities by the positive solutions of (1) for sufficiently small  $\lambda > 0$ .

By having a glance at condition  $(a_2)$  it becomes apparent that it fails whenever the function  $a$  is differentiable at the nodal point  $z$ , whereas a very simple example where  $(a_2)$  holds occurs when the function  $a$  is discontinuous at  $z$ , like, for instance, in the special case when  $a$  is assumed to be a positive constant,  $A > 0$ , in  $[z - \eta_1, z)$  and a negative constant,  $-B < 0$ , in  $(z, z + \eta_2]$ , for some  $\eta_1, \eta_2 > 0$ . The huge contrast on the nature of the positive solutions of the problem with respect to the integrability properties of the function  $a$  near the node  $z$  can also be realized by considering any weight function  $a$  satisfying the requirements of  $(a_1)$  except for the fact that  $a = 0$  in  $[z - \eta, z + \eta]$  for some  $\eta > 0$ . In such case, thanks to the convexity and concavity properties of the positive bounded variation solutions of (1) guaranteed by [23, Prop. 3.6], any positive solution  $u$  must be linear in the interval  $[z - \eta, z + \eta]$  and hence, due to (2), it cannot develop singularities.

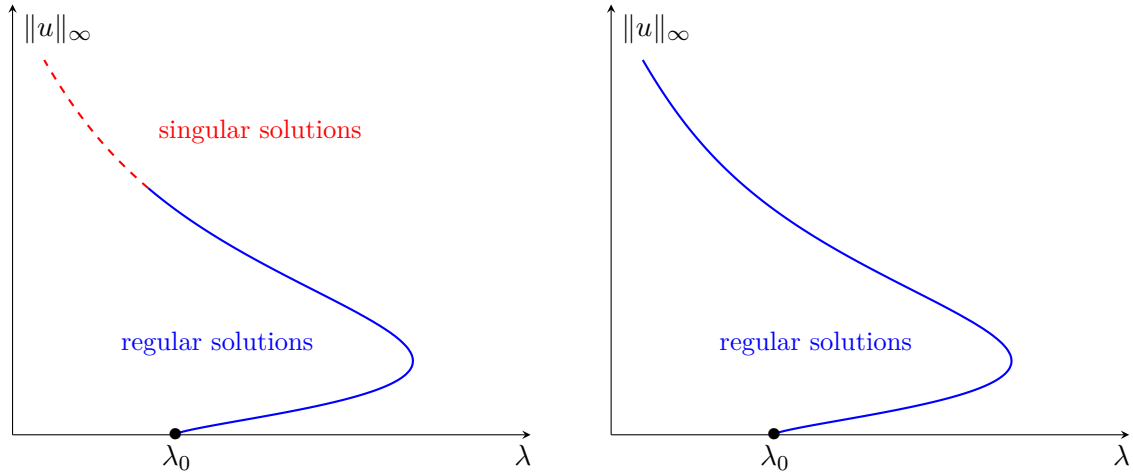
As a consequence of Theorem 1, when  $p = 2$ , the global structure of the component of the positive solutions of (1),  $\mathcal{C}_+$ , whose existence is guaranteed by the main theorem of [20], drastically changes according to whether or not the condition  $(a_2)$  holds as illustrated in Figure 1, where  $\lambda_0$  stands for the principal positive eigenvalue of the linear weighted problem

$$\begin{cases} -\varphi'' = \lambda a(x)\varphi & \text{in } (0, 1), \\ \varphi'(0) = \varphi'(1) = 0. \end{cases}$$

The non-existence of positive regular solutions of (1) in the very special cases when  $p = 2$  and the weight  $a$  is constant in  $[0, z)$  and in  $(z, 1]$  has been recently established in Section 8 of [23] by using some classical, but sophisticated, phase portrait techniques. This induced the authors to presume that an analogous non-existence result should also be valid for general weight functions  $a$ , without imposing the integrability condition  $(a_2)$ . So, they formulated [23, Th. 7.1]. Theorem 1 in particular shows that [23, Th. 7.1] has to be complemented with condition  $(a_2)$ .

Similarly as for  $p = 2$ , also in the case  $p > 2$  the global structure of the set of positive solutions of (1),  $\mathcal{C}_+$ , whose existence is now guaranteed by [22, Th.1.1] and [23, Th.10.1], changes for sufficiently small  $\lambda > 0$  according to whether or not condition  $(a_2)$  holds, as illustrated by Figure 2.

Our proof of Theorem 1 is based upon the characterization of the exact limiting profiles of the positive solutions of (1), both regular and singular, as the parameter  $\lambda$  approximates zero. These profiles are provided by the next theorem, regardless their particular nature.



**Figure 1.** Global components emanating from the positive principal eigenvalue  $\lambda_0$  in case  $p = 2$  when  $(a_2)$  holds (on the left), or  $(a_2)$  fails (on the right).

**Theorem 2.** Assume  $(a_1)$  and  $(f_1)$ , and let  $((\lambda_n, u_n))_n$  be an arbitrary sequence of positive solutions of (1) with  $\lim_{n \rightarrow \infty} \lambda_n = 0$ . Then, for sufficiently small  $\eta > 0$ , the following assertions hold:

$$\lim_{n \rightarrow +\infty} \frac{u_n(x)}{u_n(0)} = 1 \quad \text{uniformly in } x \in [0, z - \eta],$$

$$\lim_{n \rightarrow +\infty} \frac{u_n(x)}{u_n(0)} = \left( \frac{\int_0^z a(t) dt}{-\int_z^1 a(t) dt} \right)^{\frac{1}{q-1}} \quad \text{uniformly in } x \in [z + \eta, 1],$$

$$\lim_{n \rightarrow +\infty} (\lambda_n f(u_n(x))) = \frac{1}{\int_0^z a(t) dt} \quad \text{uniformly in } x \in [0, z - \eta],$$

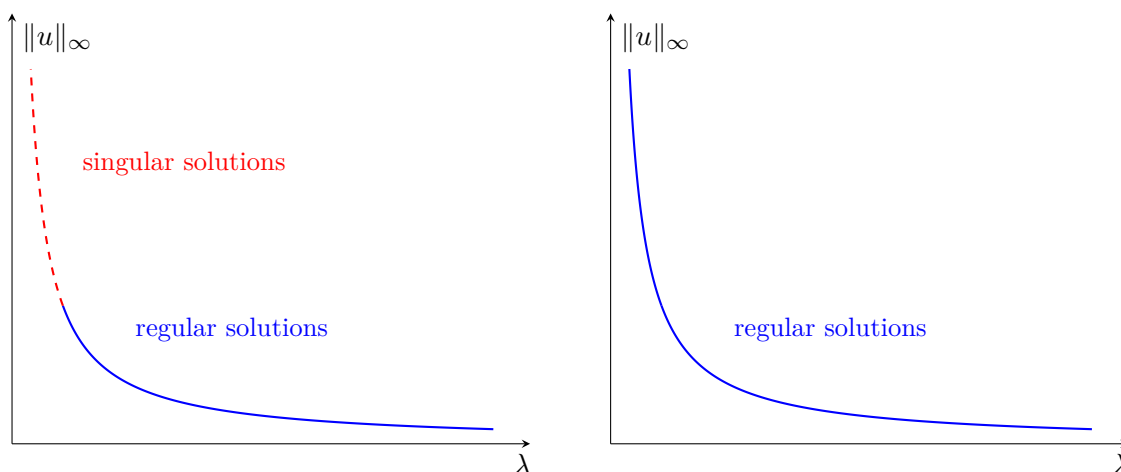
$$\lim_{n \rightarrow +\infty} (\lambda_n f(u_n(x))) = \frac{1}{-\int_z^1 a(t) dt} \quad \text{uniformly in } x \in [z + \eta, 1],$$

$$\lim_{n \rightarrow +\infty} u'_n(x) = \frac{-\int_0^x a(t) dt}{\sqrt{\left(\int_0^z a(t) dt\right)^2 - \left(\int_0^x a(t) dt\right)^2}} \quad \text{uniformly in } x \in [0, z - \eta],$$

and

$$\lim_{n \rightarrow +\infty} u'_n(x) = \frac{\int_x^1 a(t) dt}{\sqrt{\left(\int_z^1 a(t) dt\right)^2 - \left(\int_x^1 a(t) dt\right)^2}} \quad \text{uniformly in } x \in [z + \eta, 1].$$

Note that condition  $(a_2)$  is equivalent to requiring the integrability in both intervals,  $(0, z)$  and  $(z, 1)$ , of the asymptotic profile of the derivatives of the positive solutions of (1) as  $\lambda \rightarrow 0^+$ , which



**Figure 2.** Global bifurcation diagrams in case  $p > 2$  when  $(a_2)$  holds (on the left), or  $(a_2)$  fails (on the right).

is equivalent to impose that the “limiting derivative”

$$u'_\omega(x) = \begin{cases} \frac{-\int_0^x a(t) dt}{\sqrt{\left(\int_0^z a(t) dt\right)^2 - \left(\int_0^x a(t) dt\right)^2}} & \text{for } x \in [0, z), \\ \frac{\int_x^1 a(t) dt}{\sqrt{\left(\int_z^1 a(t) dt\right)^2 - \left(\int_x^1 a(t) dt\right)^2}} & \text{for } x \in (z, 1], \end{cases}$$

belongs to both  $L^1(0, z)$  and  $L^1(z, 1)$ .

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